

A class of unitary irreducible representations of the Lie superalgebra $\mathfrak{osp}(1|2n)$

S. Lievens, N.I. Stoilova* and J. Van der Jeugt

Department of Applied Mathematics and Computer Science, Ghent University,
Krijgslaan 281-S9, B-9000 Gent, Belgium

Stijn.Lievens@UGent.be, Neli.Stoilova@UGent.be, Joris.VanderJeugt@UGent.be

Abstract

Using the equivalence of the defining relations of the orthosymplectic Lie superalgebra $\mathfrak{osp}(1|2n)$ to the defining triple relations of n pairs of parabose operators b_i^\pm we construct a class of unitary irreducible (infinite-dimensional) lowest weight representations $V(p)$ of $\mathfrak{osp}(1|2n)$. We introduce an orthogonal basis of $V(p)$ in terms of Gelfand-Zetlin patterns, where the subalgebra $\mathfrak{u}(n)$ of $\mathfrak{osp}(1|2n)$ plays a crucial role and we present explicit actions of the $\mathfrak{osp}(1|2n)$ generators.

Following some physical ideas we construct a class of infinite dimensional unitary irreducible representations of the Lie superalgebra (LS) $\mathfrak{osp}(1|2n)$ [1] in an explicit form. In 1953 Green [2] introduced the so called parabose operators (PBOs) b_j^\pm ($j = 1, 2, \dots$) satisfying

$$[\{b_j^\xi, b_k^\eta\}, b_l^\epsilon] = (\epsilon - \xi)\delta_{jl}b_k^\eta + (\epsilon - \eta)\delta_{kl}b_j^\xi, \quad (1)$$

$j, k, l \in \{1, 2, \dots\}$ and $\eta, \epsilon, \xi \in \{+, -\}$ (to be interpreted as $+1$ and -1 in the algebraic expressions $\epsilon - \xi$ and $\epsilon - \eta$) as a generalization of the ordinary Bose operators. The Fock space $V(p)$ of n pairs of PBOs is a Hilbert space with a vacuum $|0\rangle$, defined by means of ($j, k = 1, 2, \dots, n$)

$$\langle 0|0\rangle = 1, \quad b_j^-|0\rangle = 0, \quad (b_j^\pm)^\dagger = b_j^\mp, \quad \{b_j^-, b_k^+\}|0\rangle = p \delta_{jk}|0\rangle, \quad (2)$$

and by irreducibility under the action of the algebra spanned by the elements b_j^+, b_j^- ($j = 1, \dots, n$), subject to (1). The parameter p is referred to as the order of the paraboson system. However the structure of the parabose Fock space is not known, also a proper basis has not been introduced. We solve these problems using the relation between n pairs of PBOs and the defining relations of the LS $\mathfrak{osp}(1|2n)$, discovered by Ganchev and Palev [3]. The orthosymplectic superalgebra $\mathfrak{osp}(1|2n)$ [1] consists of matrices of the form

$$\begin{pmatrix} 0 & a & a_1 \\ a_1^t & b & c \\ -a^t & d & -b^t \end{pmatrix}, \quad (3)$$

where a and a_1 are $(1 \times n)$ -matrices, b is any $(n \times n)$ -matrix, and c and d are symmetric $(n \times n)$ -matrices. The even elements have $a = a_1 = 0$ and the odd elements are those with

*Permanent address: Institute for Nuclear Research and Nuclear Energy, Boul. Tsarigradsko Chaussee 72, 1784 Sofia, Bulgaria

$b = c = d = 0$. Denote the row and column indices running from 0 to $2n$ and by e_{ij} the matrix with zeros everywhere except a 1 on position (i, j) . Then as a basis in the Cartan subalgebra \mathfrak{h} of $\mathfrak{osp}(1|2n)$ consider

$$h_j = e_{jj} - e_{n+j, n+j} \quad (j = 1, \dots, n). \quad (4)$$

In terms of the dual basis δ_j of \mathfrak{h}^* , the root vectors and corresponding roots of $\mathfrak{osp}(1|2n)$ are given by:

$$\begin{aligned} e_{0,k} - e_{n+k,0} &\leftrightarrow -\delta_k, & e_{0,n+k} + e_{k,0} &\leftrightarrow \delta_k, & k = 1, \dots, n, \text{ odd}, \\ e_{j,n+k} + e_{k,n+j} &\leftrightarrow \delta_j + \delta_k, & e_{n+j,k} + e_{n+k,j} &\leftrightarrow -\delta_j - \delta_k, & j \leq k = 1, \dots, n, \text{ even}, \\ e_{j,k} - e_{n+k,n+j} &\leftrightarrow \delta_j - \delta_k, & j \neq k = 1, \dots, n, &\text{ even}. \end{aligned}$$

Introduce the following multiples of the odd root vectors

$$b_k^+ = \sqrt{2}(e_{0,n+k} + e_{k,0}), \quad b_k^- = \sqrt{2}(e_{0,k} - e_{n+k,0}) \quad (k = 1, \dots, n). \quad (5)$$

Then the following holds [3]

Theorem 1 (Ganchev and Palev). *As a Lie superalgebra defined by generators and relations, $\mathfrak{osp}(1|2n)$ is generated by $2n$ odd elements b_k^\pm subject to the parabose relations (1).*

From (4) and (5) it follows that $\{b_j^-, b_j^+\} = 2h_j, j = 1, \dots, n$ and using (2) we have:

Corollary 2. *The parabose Fock space $V(p)$ is the unitary irreducible representation of $\mathfrak{osp}(1|2n)$ with lowest weight $(\frac{p}{2}, \frac{p}{2}, \dots, \frac{p}{2})$.*

We can construct the representation $V(p)$ [4] using an induced module construction with an appropriate chain of subalgebras.

Proposition 3. *A basis for the even subalgebra $\mathfrak{sp}(2n)$ of $\mathfrak{osp}(1|2n)$ is given by the $2n^2 + n$ elements $\{b_j^\pm, b_k^\pm\}, 1 \leq j \leq k \leq n, \{b_j^+, b_k^-\}, 1 \leq j, k \leq n$. The n^2 elements $\{b_j^+, b_k^-\}, j, k = 1, \dots, n$ are a basis for the $\mathfrak{sp}(2n)$ subalgebra $\mathfrak{u}(n)$.*

The subalgebra $\mathfrak{u}(n)$ can be extended to a parabolic subalgebra $\mathcal{P} = \text{span}\{\{b_j^+, b_k^-\}, b_j^-, \{b_j^-, b_k^-\}, j, k = 1, \dots, n\}$ [4] of $\mathfrak{osp}(1|2n)$. Recall that $\{b_j^-, b_k^+\}|0\rangle = p\delta_{jk}|0\rangle$, with $\{b_j^-, b_j^+\} = 2h_j$. Then the space spanned by $|0\rangle$ is a trivial one-dimensional $\mathfrak{u}(n)$ module $\mathbb{C}|0\rangle$ of weight $(\frac{p}{2}, \dots, \frac{p}{2})$. Since $b_j^-|0\rangle = 0$, the module $\mathbb{C}|0\rangle$ can be extended to a one-dimensional \mathcal{P} module. Now we define the induced $\mathfrak{osp}(1|2n)$ module $\bar{V}(p)$ with lowest weight $(\frac{p}{2}, \dots, \frac{p}{2})$: $\bar{V}(p) = \text{Ind}_{\mathcal{P}}^{\mathfrak{osp}(1|2n)} \mathbb{C}|0\rangle$.

By the Poincaré-Birkhoff-Witt theorem [1, 4], it is easy to give a basis for $\bar{V}(p)$:

$$(b_1^+)^{k_1} \dots (b_n^+)^{k_n} (\{b_1^+, b_2^+\})^{k_{12}} (\{b_1^+, b_3^+\})^{k_{13}} \dots (\{b_{n-1}^+, b_n^+\})^{k_{n-1,n}} |0\rangle,$$

where $k_1, \dots, k_n, k_{12}, k_{13}, \dots, k_{n-1,n} \in \mathbb{Z}_+$. However in general $\bar{V}(p)$ is not a simple module and let $M(p)$ be the maximal nontrivial submodule of $\bar{V}(p)$. Then the simple module (irreducible module), corresponding to the paraboson Fock space, is $V(p) = \bar{V}(p)/M(p)$. The purpose is now to determine the vectors belonging to $M(p)$ and also to find explicit matrix elements of the $\mathfrak{osp}(1|2n)$ generators b_j^\pm in an appropriate basis of $V(p)$.

From the basis in $\bar{V}(p)$, it is easy to write down the character of $\bar{V}(p)$:

$$\text{char } \bar{V}(p) = \frac{(x_1 \dots x_n)^{p/2}}{\prod_{i=1}^n (1 - x_i) \prod_{1 \leq j < k \leq n} (1 - x_j x_k)}. \quad (6)$$

Such expressions have an interesting expansion in terms of Schur functions.

Proposition 4 (Cauchy, Littlewood). *Let x_1, \dots, x_n be a set of n variables. Then [5]*

$$\frac{1}{\prod_{i=1}^n (1-x_i) \prod_{1 \leq j < k \leq n} (1-x_j x_k)} = \sum_{\lambda} s_{\lambda}(x_1, \dots, x_n) = \sum_{\lambda} s_{\lambda}(x). \quad (7)$$

Herein the sum is over all partitions λ and $s_{\lambda}(x)$ is the Schur symmetric function [6].

The characters of finite dimensional $\mathfrak{u}(n)$ representations are given by such Schur functions $s_{\lambda}(x)$. For such finite dimensional $\mathfrak{u}(n)$ representations labelled by a partition λ , there is a known basis: the Gelfand-Zetlin basis (GZ) [7]. We shall use the $\mathfrak{u}(n)$ GZ basis vectors as our new basis for $\bar{V}(p)$. Thus the new basis of $\bar{V}(p)$ consists of vectors of the form

$$|m\rangle \equiv |m\rangle^n \equiv \begin{pmatrix} m_{1n} & \cdots & \cdots & m_{n-1,n} & m_{nn} \\ m_{1,n-1} & \cdots & \cdots & m_{n-1,n-1} & \\ \vdots & & \ddots & & \\ m_{11} & & & & \end{pmatrix} = \begin{pmatrix} [m]^n \\ |m\rangle^{n-1} \end{pmatrix}, \quad (8)$$

where the top line of the pattern, also denoted by the n -tuple $[m]^n$, is any partition λ (consisting of non increasing nonnegative numbers) with $\ell(\lambda) \leq n$. The label p itself is dropped in the notation of $|m\rangle$. The remaining $n-1$ lines of the pattern will sometimes be denoted by $|m\rangle^{n-1}$. So all m_{ij} in the above GZ-pattern are nonnegative integers, satisfying the *betweenness conditions* $m_{i,j+1} \geq m_{ij} \geq m_{i+1,j+1}$, $1 \leq i \leq j \leq n-1$. Note that, since the weight of $|0\rangle$ is $(\frac{p}{2}, \dots, \frac{p}{2})$, the weight of the above vector is determined by

$$h_k |m\rangle = \left(\frac{p}{2} + \sum_{j=1}^k m_{jk} - \sum_{j=1}^{k-1} m_{j,k-1} \right) |m\rangle. \quad (9)$$

The triple relations (1), imply that $(b_1^+, b_2^+, \dots, b_n^+)$ is a standard $\mathfrak{u}(n)$ tensor of rank $(1, 0, \dots, 0)$. Therefore we can attach a unique GZ-pattern with top line $10 \cdots 0$ to every b_j^+ , corresponding to the weight $+\delta_j$. Explicitly:

$$b_j^+ \sim \begin{matrix} 10 \cdots 000 \\ 10 \cdots 00 \\ \cdots \\ 0 \cdots 0 \\ \cdots \\ 0 \end{matrix}, \quad (10)$$

where the pattern consists of $j-1$ zero rows at the bottom, and the first $n-j+1$ rows are of the form $10 \cdots 0$. The tensor product rule in $\mathfrak{u}(n)$ reads $([m]^n) \otimes (10 \cdots 0) = ([m]_{+1}^n) \oplus ([m]_{+2}^n) \oplus \cdots \oplus ([m]_{+n}^n)$ where $([m]^n) = (m_{1n}, m_{2n}, \dots, m_{nn})$ and a subscript $\pm k$ indicates an increment of the k th label by ± 1 : $([m]_{\pm k}^n) = (m_{1n}, \dots, m_{kn} \pm 1, \dots, m_{nn})$. A general matrix element of b_j^+ can now be written as follows:

$$(m' | b_j^+ | m) = \left(\begin{array}{c|c} [m]_{+k}^n & [m]^n \\ \hline (m')^{n-1} & |m\rangle^{n-1} \end{array} \right) = \left(\begin{array}{c|c} [m]^n & \begin{matrix} 10 \cdots 00 \\ 10 \cdots 0 \\ \cdots \\ 0 \end{matrix} \\ \hline |m\rangle^{n-1} & (m')^{n-1} \end{array} \right) ([m]_{+k}^n || b^+ || [m]^n).$$

The first factor in the right hand side is a $\mathfrak{u}(n)$ Clebsch-Gordan coefficient [8], the second factor is a reduced matrix element. By the tensor product rule, the first line of $|m'\rangle$ has to be $[m']^n = [m]_{+k}^n$ for some k -value.

The special $\mathfrak{u}(n)$ CGCs appearing here are well known, and have fairly simple expressions. They can be found, e.g. in [8, 9]. The actual problem is now converted into finding expressions

for the reduced matrix elements, i.e. for the functions $F_k([m]^n)$, for arbitrary n -tuples of non increasing nonnegative integers $[m]^n = (m_{1n}, m_{2n}, \dots, m_{nn})$:

$$F_k([m]^n) = F_k(m_{1n}, m_{2n}, \dots, m_{nn}) = ([m]_{+k}^n || b^+ || [m]^n). \quad (11)$$

So one can write:

$$b_j^+ |m\rangle = \sum_{k, m'} \left(\begin{array}{c|c} [m]^n & 10 \cdots 00 \\ |m\rangle^{n-1} & 10 \cdots 0 \\ & \cdots \\ & 0 \end{array} \middle| \begin{array}{c} [m]_{+k}^n \\ |m'\rangle^{n-1} \end{array} \right) F_k([m]^n) \left| \begin{array}{c} [m]_{+k}^n \\ |m'\rangle^{n-1} \end{array} \right\rangle, \quad (12)$$

$$b_j^- |m\rangle = \sum_{k, m'} \left(\begin{array}{c|c} [m]_{-k}^n & 10 \cdots 00 \\ |m'\rangle^{n-1} & 10 \cdots 0 \\ & \cdots \\ & 0 \end{array} \middle| \begin{array}{c} [m]^n \\ |m\rangle^{n-1} \end{array} \right) F_k([m]_{-k}^n) \left| \begin{array}{c} [m]_{-k}^n \\ |m'\rangle^{n-1} \end{array} \right\rangle. \quad (13)$$

For $j = n$, the CGCs in (12)-(13) take a simple form [8], and we have

$$b_n^+ |m\rangle = \sum_{i=1}^n \left(\frac{\prod_{k=1}^{n-1} (m_{k, n-1} - m_{in} - k + i - 1)}{\prod_{k \neq i=1}^n (m_{kn} - m_{in} - k + i)} \right)^{1/2} F_i(m_{1n}, m_{2n}, \dots, m_{nn}) |m\rangle_{+in}; \quad (14)$$

$$b_n^- |m\rangle = \sum_{i=1}^n \left(\frac{\prod_{k=1}^{n-1} (m_{k, n-1} - m_{in} - k + i)}{\prod_{k \neq i=1}^n (m_{kn} - m_{in} - k + i + 1)} \right)^{1/2} F_i(m_{1n}, \dots, m_{in} - 1, \dots, m_{nn}) |m\rangle_{-in}. \quad (15)$$

In order to determine the n unknown functions F_k , one can start from the following action:

$$\{b_n^-, b_n^+\} |m\rangle = 2h_n |m\rangle = (p + 2(\sum_{j=1}^n m_{jn} - \sum_{j=1}^{n-1} m_{j, n-1})) |m\rangle. \quad (16)$$

Expressing the left hand side by means of (14)-(15), one finds a system of coupled recurrence relations for the functions F_k . Taking the appropriate boundary conditions into account, we have been able to solve this system of relations [9].

Proposition 5. *The reduced matrix elements F_k appearing in the actions of b_j^\pm on vectors $|m\rangle$ of $\bar{V}(p)$ are given by:*

$$F_k(m_{1n}, m_{2n}, \dots, m_{nn}) = (-1)^{m_{k+1, n} + \dots + m_{nn}} (m_{kn} + n + 1 - k + \mathcal{E}_{m_{kn}}(p - n))^{1/2} \times \prod_{j \neq k=1}^n \left(\frac{m_{jn} - m_{kn} - j + k}{m_{jn} - m_{kn} - j + k - \mathcal{O}_{m_{jn} - m_{kn}}} \right)^{1/2}, \quad (17)$$

where \mathcal{E} and \mathcal{O} are the even and odd functions defined by $\mathcal{E}_j = 1$ if j is even and 0 otherwise, $\mathcal{O}_j = 1$ if j is odd and 0 otherwise.

The proof consists of verifying that all triple relations (1) hold when acting on any vector $|m\rangle$. Each such verification leads to an algebraic identity in n variables m_{1n}, \dots, m_{nn} . In these computations, there are some intermediate verifications: e.g. the action $\{b_j^+, b_k^-\} |m\rangle$ should leave the top row of the GZ-pattern $|m\rangle$ invariant (since $\{b_j^+, b_k^-\}$ belongs to $\mathfrak{u}(n)$). Furthermore, it must yield (up to a factor 2) the known action of the standard $\mathfrak{u}(n)$ matrix elements E_{jk} in the classical GZ-basis. Consider now the factor $(m_{kn} + n + 1 - k + \mathcal{E}_{m_{kn}}(p - n))$ in the expression of $F_k([m]^n)$. This is the only factor in the right hand side of (17) that may become zero. If this factor is zero or negative, the assigned vector $|m'\rangle$ belongs to $M(p)$. Recall that the integers m_{jn} satisfy $m_{1n} \geq m_{2n} \geq \dots \geq m_{nn} \geq 0$. If $m_{kn} = 0$ (its smallest possible value), then this factor in F_k takes the value $(p - k + 1)$. So the p -values $1, 2, \dots, n - 1$ play a special role leading to the following result [9]:

Theorem 6. *The $\mathfrak{osp}(1|2n)$ representation $V(p)$ with lowest weight $(\frac{p}{2}, \dots, \frac{p}{2})$ is a unirrep if and only if $p \in \{1, 2, \dots, n-1\}$ or $p > n-1$. For $p > n-1$, $V(p) = \bar{V}(p)$ and $\text{char } V(p) = \frac{(x_1 \cdots x_n)^{p/2}}{\prod_i (1-x_i) \prod_{j < k} (1-x_j x_k)} = (x_1 \cdots x_n)^{p/2} \sum_{\lambda} s_{\lambda}(x)$. For $p \in \{1, 2, \dots, n-1\}$, $V(p) = \bar{V}(p)/M(p)$ with $M(p) \neq 0$. The structure of $V(p)$ is determined by $\text{char } V(p) = (x_1 \cdots x_n)^{p/2} \sum_{\lambda, \ell(\lambda) \leq p} s_{\lambda}(x)$ where $\ell(\lambda)$ is the length of the partition λ .*

The explicit action of the $\mathfrak{osp}(1|2n)$ generators in $V(p)$ is given by (12)-(13), and the basis is orthogonal and normalized. For $p \in \{1, 2, \dots, n-1\}$ this action remains valid, provided one keeps in mind that all vectors with $m_{p+1, n} \neq 0$ must vanish.

Note that the first line of Theorem 6 can also be deduced from [10, 11], where all unirreps of $\mathfrak{osp}(1|2n)$ are classified.

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References

- [1] V.G. Kac. Representations of Classical Lie Superalgebras. Lecture Notes in Math. **626** (1978), 597-626 (Berlin, Springer).
- [2] H.S. Green. A Generalized Method of Field Quantization. Phys. Rev. **90** (1953), 270-273.
- [3] A.Ch. Ganchev and T.D. Palev. A Lie Superalgebraic Interpretation of the Para-Bose Statistics. J. Math. Phys. **21** (1980), 797-799.
- [4] T.D. Palev. Lie Superalgebras, Infinite-Dimensional Algebras and Quantum Statistics. Rep. Math. Phys. **31** (1992), 241-262.
- [5] D.E. Littlewood. *The theory of Group Characters and Matrix Representations of Groups*. Oxford University Press, Oxford, 1950.
- [6] I.G. Macdonald. *Symmetric Functions and Hall Polynomials*. Oxford University Press, Oxford, 2nd edition, 1995.
- [7] I.M. Gel'fand and M.L. Zetlin. Finite-Dimensional Representations of the Group of Unitary Matrices. Dokl. Akad. Nauk SSSR **71** (1950), 825-828.
- [8] N.Ja. Vilenkin and A.U. Klimyk. *Representation of Lie Groups and Special Functions, Vol. 3* Kluwer Academic Publishers, Dordrecht, 1992.
- [9] S. Lievens, N.I. Stoilova and J. Van der Jeugt. The paraboson Fock space and unitary irreducible representations of the Lie superalgebra $\mathfrak{osp}(1|2n)$. Preprint arXiv:0706.4196 [hep-th].
- [10] H.P. Jakobsen. The full set of unitarizable highest weight modules of basic classical Lie superalgebras. Mem. Amer. Math. Soc. **111** No 532, 1994.
- [11] V.K. Dobrev and R.B. Zhang. Positive Energy Unitary Irreducible Representations of the Superalgebras $\mathfrak{osp}(1|2n; \mathbb{R})$. Phys. Atom. Nuclei **68** (2005), 1660-1669.