

The East and Southeast Asia Workshop on Inverse Problems and Optimal Control (ESEAW: IPOC)

# EXISTENCE OF A WEAK SOLUTION TO A NONLINEAR INDUCTION HARDENING PROBLEM FOR A STEEL WORKPIECE

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#### Outline

Mathematical model

Variational formulation

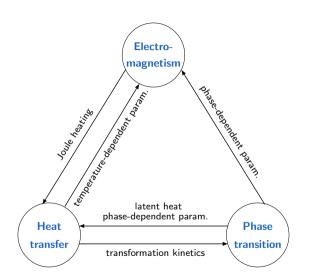
Time discretization

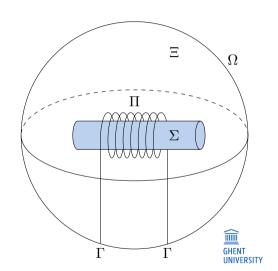
Existence of a solution





## Induction hardening





## Jumping material coefficients

The subscripts  $\Sigma, \Pi$  and  $\Xi$  are used to distinguish the value of material coefficients on subdomains. For instances,

$$\mu = \begin{cases} \mu_\Pi & \text{ in } \Pi, \\ \mu_\Sigma & \text{ in } \Sigma, \\ \mu_\Xi & \text{ in } \Xi. \end{cases} \quad \text{and} \quad \sigma = \begin{cases} \sigma_\Pi & \text{ in } \Pi, \\ \sigma_\Sigma & \text{ in } \Sigma, \\ 0 & \text{ in } \Xi. \end{cases}$$

In general, material coefficients have jumps at the interfaces of different subdomains.

$$[\![\mu]\!]_{\partial\Sigma} := \mu_{\Xi} - \mu_{\Sigma} \neq 0,$$
$$[\![\mu]\!]_{\partial\Pi\backslash\Gamma} := \mu_{\Xi} - \mu_{\Pi} \neq 0.$$





## Electromagnetic subproblem

#### Maxwell's equations

$$abla \cdot \mathbf{B} = 0,$$
 (Gauss's law)  
 $abla \times \mathbf{E} = -\partial_t \mathbf{B},$  (Faraday's law)  
 $abla \times \mathbf{H} = \mathbf{J},$  (Ampère's law)

with the constitutive relations

$$m{B} = \mu m{H}, \ m{J} = \sigma m{E}.$$
 (Ohm's law)

Boundary condition

$$\mathbf{B} \cdot \mathbf{n} = 0.$$

Interface conditions

$$[\![\boldsymbol{B} \cdot \mathbf{n}]\!] = 0,$$

$$[\mathbf{B} \cdot \mathbf{n}] = 0, \qquad [\mathbf{H} \times \mathbf{n}] = \mathbf{0},$$

$$\llbracket \boldsymbol{E} \times \mathbf{n} \rrbracket = \mathbf{0}.$$



#### Nonlinear B - H relation

Suppose that

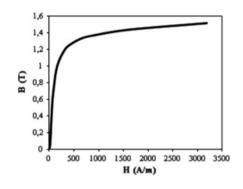
$$\mu_{\Sigma}(\boldsymbol{B}) = \hat{\mu}_{\Sigma} \ m^{-1}(|\boldsymbol{B}|),$$

then the magnetization curve

$$\boldsymbol{H} = \hat{\mu}^{-1} \boldsymbol{M}(\boldsymbol{B}),$$

where

$$\hat{\mu} = \begin{cases} \hat{\mu}_{\Sigma} & \text{in } \Sigma, \\ \mu & \text{in } \Pi \cup \Xi, \end{cases} \quad \text{and} \quad \boldsymbol{M}(\boldsymbol{b}) = \begin{cases} m(|\boldsymbol{b}|) \, \boldsymbol{b} & \text{in } \Sigma, \\ \boldsymbol{b} & \text{in } \Pi \cup \Xi. \text{ } \underline{\widehat{\text{im}}} \end{cases}$$



$$oldsymbol{M}(oldsymbol{b}) = egin{cases} m \| oldsymbol{b} \| oldsymbol{b$$

in 
$$\Sigma$$
,



## Assumptions on the function m

The function  $m: \mathbb{R}^+ \to \mathbb{R}^+$  is supposed to be differentiable and satisfy

$$m(s) + |m'(s)| s \le C_M \quad \forall s \in \mathbb{R}^+.$$

In addition, for each  $s \in \mathbb{R}^+$ , one of the following conditions is satisfied

$$m(s) - |m'(s)| s \ge c_M > 0,$$

or

$$m(s) \ge c_M$$
 and  $m'(s) \ge 0$ .

#### Lemma

Under these assumptions, the nonlinear operator  ${m M}$  is strongly monotone, coercive and Lipschitz continuous.



## $A-\phi$ potential formulation

The BVP for the electric scalar potential  $\phi$ 

$$\begin{cases} \nabla \cdot (-\sigma_{\Pi} \nabla \phi) = 0 & \text{in} \quad \Pi \times (0, T), \\ -\sigma_{\Pi} \nabla \phi \cdot \mathbf{n} = 0 & \text{on} \quad (\partial \Pi \setminus \Gamma) \times (0, T), \\ -\sigma_{\Pi} \nabla \phi \cdot \mathbf{n} = j & \text{on} \quad \Gamma \times (0, T). \end{cases}$$

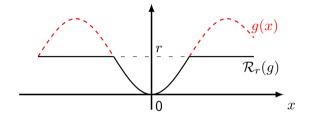
The IBVP for the magnetic vector potential  $m{A}$  ( $m{B} := 
abla imes m{A}$ )

$$\begin{cases} \sigma \partial_t \boldsymbol{A} + \nabla \times \hat{\mu}^{-1} \boldsymbol{M} \left( \nabla \times \boldsymbol{A} \right) + \chi_\Pi \sigma \nabla \phi = \boldsymbol{0} & \text{in} \quad \Omega \times (0,T), \\ \nabla \cdot \boldsymbol{A} = 0 & \text{in} \quad \Omega \times (0,T), \\ \boldsymbol{A} \times \mathbf{n} = \boldsymbol{0} & \text{on} \quad \partial \Omega \times (0,T), \\ \left[ \hat{\mu}^{-1} \boldsymbol{M} (\nabla \times \boldsymbol{A}) \times \mathbf{n} \right] = \boldsymbol{0} & \text{on} \quad (\partial \Theta \setminus \Gamma) \times (0,T), \\ \boldsymbol{A} \cdot (\cdot,0) = \boldsymbol{A}_0 & \text{in} \quad \Theta := \Sigma \cup \Pi. \end{cases}$$

## Energy balance in the steel workpiece

The Joule heat caused by the eddy current

$$\boldsymbol{J} = \sigma_{\Sigma} |\partial_t \boldsymbol{A}|^2$$
.



The IBVP for the temperature  $\mathcal{T}$ 

$$\begin{cases} \rho c_p(\mathcal{T}) \partial_t \mathcal{T} - \nabla \cdot (\kappa(\mathcal{T}) \nabla \mathcal{T}) = \mathcal{R}_r \left( \sigma_{\Sigma} |\partial_t \mathbf{A}|^2 \right) - \rho f(z, \mathcal{T}) \partial_t z \\ -\kappa(\mathcal{T}) \nabla \mathcal{T} \cdot \mathbf{n} = \alpha(\mathcal{T} - u^c) \\ \mathcal{T}(\cdot, 0) = \mathcal{T}_0 \end{cases}$$

$$\quad \text{in} \quad \Sigma \times (0,T),$$

$$\quad \text{on} \quad \partial \Sigma \times (0,T),$$

in 
$$\Sigma$$

where

$$f(z, \mathcal{T}) = -L\left(z_{\mathsf{eq}}(\mathcal{T}) - z - \mathcal{T}z'_{\mathsf{eq}}(\mathcal{T})\right).$$



#### Kirchhoff transformation

We introduce the primitive function  $\hat{\kappa}$  from  $\kappa$ 

$$\hat{\kappa}(x) = \int_{0}^{x} \kappa(s) \, \mathrm{d}s.$$

Let  $u := \hat{\kappa}(\mathcal{T})$  and  $u_0 := \hat{\kappa}(\mathcal{T}_0)$ . The IBVP for the transformed temperature u

$$\begin{cases} \beta(\hat{\kappa}^{-1}(u))\partial_t u - \Delta u = \mathcal{R}_r \left(\sigma_{\Sigma} |\partial_t \mathbf{A}|^2\right) - \rho f(z, \hat{\kappa}^{-1}(u))\partial_t z & \text{in } \Sigma \times (0, T), \\ -\nabla u \cdot \mathbf{n} = \alpha(\hat{\kappa}^{-1}(u) - u^c) & \text{on } \partial\Sigma \times (0, T), \\ u(\cdot, 0) = u_0 & \text{in } \Sigma. \end{cases}$$

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#### Leblond-Devaux model

Leblond-Devaux model for the austenite proportion z

$$\begin{cases} \partial_t z = \frac{z_{\text{eq}}(\mathcal{T}) - z}{\vartheta(\mathcal{T})} & \text{in } \Sigma \times (0, T), \\ z(\cdot, 0) = z_0 = 0 & \text{in } \Sigma. \end{cases}$$

or with the transformed temperature u

$$\begin{cases} \partial_t \mathbf{z} = \frac{z_{\text{eq}}(\hat{\kappa}^{-1}(u)) - \mathbf{z}}{\vartheta(\hat{\kappa}^{-1}(u))} & \text{in } \Sigma \times (0, T), \\ \mathbf{z}(\cdot, 0) = z_0 = 0 & \text{in } \Sigma. \end{cases}$$





## Induction hardening problem

#### Electric potential $\phi$

$$\begin{cases} \nabla \cdot (-\sigma_{\Pi} \nabla \phi) = 0, \\ -\sigma_{\Pi} \nabla \phi \cdot \mathbf{n} = 0, \\ -\sigma_{\Pi} \nabla \phi \cdot \mathbf{n} = j. \end{cases}$$

#### Magnetic potential A

$$\begin{cases} \sigma(\boldsymbol{z})\partial_{t}\boldsymbol{A} + \nabla \times \hat{\mu}^{-1}\boldsymbol{M}(\nabla \times \boldsymbol{A}) \\ + \chi_{\Pi}\sigma\nabla\phi = \boldsymbol{0}, \\ \nabla \cdot \boldsymbol{A} = 0, \\ \boldsymbol{A} \times \mathbf{n} = \boldsymbol{0}, \\ [\hat{\mu}^{-1}\boldsymbol{M}(\nabla \times \boldsymbol{A}) \times \mathbf{n}] = \boldsymbol{0}, \\ \boldsymbol{A}(\cdot,0) = \boldsymbol{A}_{0}. \end{cases}$$

#### Temperature u

$$\begin{cases} \beta(\mathbf{z}, \hat{\kappa}^{-1}(u))\partial_t u - \Delta u \\ = \mathcal{R}_r \left( \sigma_{\Sigma}(\mathbf{z}) |\partial_t \mathbf{A}|^2 \right) - \rho(\mathbf{z}) f(\mathbf{z}, \hat{\kappa}^{-1}(u)) \partial_t \mathbf{z}, \\ -\nabla u \cdot \mathbf{n} = \alpha(\hat{\kappa}^{-1}(u) - u^c), \\ u(\cdot, 0) = u_0. \end{cases}$$

#### Austenite proportion *z*

$$\begin{cases} \partial_t z = \frac{z_{\text{eq}}(\hat{\kappa}^{-1}(u)) - z}{\vartheta(\hat{\kappa}^{-1}(u))}, \\ z(\cdot, 0) = z_0 = 0. \end{cases}$$





## Function spaces

Let us introduce the following Hilbert spaces

$$\begin{split} & Z := H^1(\Pi)/\mathbb{R}, \\ & \boldsymbol{W}_0 := \left\{ \boldsymbol{f} \in \mathbf{H}(\mathsf{div},\Omega) \cap \mathbf{H}_0(\mathsf{curl},\Omega) : \nabla \cdot \boldsymbol{f} = 0 \right\}, \end{split}$$

equipped with the norms

$$\begin{aligned} & \|f\|_{\mathbf{Z}} := & \|\nabla f\|_{\mathbf{L}^2(\Pi)} \,, \\ & \|f\|_{\mathbf{W}_0} := & \|\nabla \times f\|_{\mathbf{L}^2(\Omega)} \,. \end{aligned}$$

Remark: The space  $W_0$  is continuously embedded into  $\mathbf{H}^1(\Omega)$  if the domain  $\Omega$  is bounded and either the boundary  $\partial\Omega$  is of class  $C^{1,1}$  or  $\Omega$  is a convex polyhedron.

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#### Variational formulation

Find  $\phi \in \mathbf{Z}, z \in \mathbf{L}^2(\Sigma), \mathbf{A} \in \mathbf{W}_0$  and  $u \in \mathbf{H}^1(\Sigma)$  such that

$$\begin{split} \sigma_{\Pi} \left( \nabla \phi, \nabla \psi \right)_{\Pi} + (j, \psi)_{\Gamma} &= 0, \\ \left( \partial_{t} z, \varsigma \right)_{\Sigma} &= \left( \frac{z_{\text{eq}} (\hat{\kappa}^{-1}(u)) - z}{\vartheta(\hat{\kappa}^{-1}(u))}, \varsigma \right)_{\Sigma}, \\ \left( \sigma(z) \partial_{t} \boldsymbol{A}, \boldsymbol{\varphi} \right)_{\Theta} + \left( \hat{\mu}^{-1} \boldsymbol{M} (\nabla \times \boldsymbol{A}), \nabla \times \boldsymbol{\varphi} \right)_{\Omega} + \sigma_{\Pi} \left( \nabla \phi, \boldsymbol{\varphi} \right)_{\Pi} &= 0, \\ \left( \beta(z, \hat{\kappa}^{-1}(u)) \partial_{t} u, w \right)_{\Sigma} + (\nabla u, \nabla w)_{\Sigma} + \alpha \left( \hat{\kappa}^{-1}(u) - u^{c}, w \right)_{\partial \Sigma} \\ &= \left( \mathcal{R}_{r} \left( \sigma_{\Sigma}(z) |\partial_{t} \boldsymbol{A}|^{2} \right), w \right)_{\Sigma} - \left( \rho(z) f(z, \hat{\kappa}^{-1}(u)) \partial_{t} z, w \right)_{\Sigma} \end{split}$$

are valid for any  $\psi \in \mathbf{Z}, \varsigma \in \mathrm{L}^2(\Sigma), \boldsymbol{\varphi} \in \boldsymbol{W}_0$  and  $w \in \mathrm{H}^1(\Sigma)$ .



#### Time discretization

Let [0,T] be partitioned into n equidistant subintervals with time step

$$\tau = \frac{T}{n}.$$

For any function v, we introduce the notations

$$v_i = v(t_i),$$
  $\delta v_i = \frac{v_i - v_{i-1}}{\tau},$ 

which stand for the value of function v and the approximation of its time derivative at the time-point  $t_i := i\tau, i = 1, 2, ..., n$ .





## Time discretization (cont.)

For every  $i=1,2,\ldots,n$ , find  $\phi_i\in \mathbb{Z}, z_i\in \mathrm{L}^2(\Sigma), \boldsymbol{A}_i\in \boldsymbol{W}_0$  and  $u_i\in \mathrm{H}^1(\Sigma)$  such that

$$\begin{split} \sigma_{\Pi} \left( \nabla \phi_{i}, \nabla \psi \right)_{\Pi} + \left( j_{i}, \psi \right)_{\Gamma} &= 0, \\ \left( \delta z_{i}, \varsigma \right)_{\Sigma} &= \left( \frac{z_{\text{eq}} (\hat{\kappa}^{-1} (u_{i-1})) - z_{i}}{\vartheta (\hat{\kappa}^{-1} (u_{i-1}))}, \varsigma \right)_{\Sigma}, \\ \left( \sigma(z_{i}) \delta \boldsymbol{A}_{i}, \boldsymbol{\varphi} \right)_{\Theta} + \left( \hat{\mu}^{-1} \boldsymbol{M} (\nabla \times \boldsymbol{A}_{i}), \nabla \times \boldsymbol{\varphi} \right)_{\Omega} + \sigma_{\Pi} \left( \nabla \phi_{i}, \boldsymbol{\varphi} \right)_{\Pi} &= 0, \\ \left( \beta(z_{i}, \hat{\kappa}^{-1} (u_{i-1})) \delta u_{i}, w \right)_{\Sigma} + \left( \nabla u_{i}, \nabla w \right)_{\Sigma} + \alpha \left( \hat{\kappa}^{-1} (u_{i-1}) - u_{i}^{c}, w \right)_{\partial \Sigma} \\ &= \left( \mathcal{R}_{r} \left( \sigma_{\Sigma}(z_{i}) |\delta \boldsymbol{A}_{i}|^{2} \right), w \right)_{\Sigma} - \left( \rho(z_{i}) f(z_{i}, \hat{\kappa}^{-1} (u_{i-1})) \ \delta z_{i}, w \right)_{\Sigma} \end{split}$$

are valid for any  $\psi \in \mathbf{Z}, \varsigma \in \mathrm{L}^2(\Sigma), \varphi \in \mathbf{W}_0$  and  $w \in \mathrm{H}^1(\Sigma)$ .



## Solvability of the discretized problem

#### Lemma

There exists a unique solution  $(\phi_i, z_i, \mathbf{A}_i, u_i) \in \mathbb{Z} \times L^2(\Sigma) \times \mathbf{W}_0 \times H^1(\Sigma)$  to the discretized variational problem for any i = 1, 2, ..., n.

#### Proof:

- The existence of a unique solution  $(\phi_i, z_i, u_i)$  is guaranteed by using the Lax-Milgram lemma.
- Thanks to the properties of the nonlinear operator M, we invoke the Browder-Minty theorem to prove the existence of a unique solution  $A_i$ .



## A priori estimates

#### Lemma

There exists a constant C>0 such that

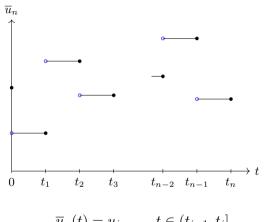
(i) 
$$\max_{1 \le l \le n} \|\nabla \delta \phi_l\|_{\mathbf{L}^2(\Pi)} \le C,$$

(ii) 
$$0 \le z_i(\boldsymbol{x}) \le 1$$
,  $\left| \delta z_i(\boldsymbol{x}) \right| \le C$ ,  $\max_{1 \le l \le n} \|\nabla z_l\|_{\mathbf{L}^2(\Sigma)} \le C$ ,

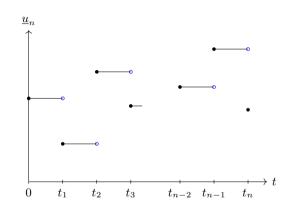
(iii) 
$$\max_{1 \le l \le n} \|\delta \mathbf{A}_l\|_{\mathbf{L}^2(\Theta)}^2 + \sum_{i=1}^n \|\delta \mathbf{A}_i - \delta \mathbf{A}_{i-1}\|_{\mathbf{L}^2(\Theta)}^2 + \sum_{i=1}^n \|\nabla \times \delta \mathbf{A}_i\|_{\mathbf{L}^2(\Omega)}^2 \tau \le C,$$

$$(iv) \sum_{i=1}^{n} \|\delta u_i\|_{\mathrm{L}^2(\Sigma)}^2 \tau + \max_{1 \le l \le n} \|\nabla u_l\|_{\mathbf{L}^2(\Sigma)}^2 + \sum_{i=1}^{n} \|\nabla u_i - \nabla u_{i-1}\|_{\mathbf{L}^2(\Sigma)}^2 + \max_{1 \le l \le n} \|u_l\|_{\mathrm{L}^2(\partial \Sigma)}^2 \le C.$$

## Rothe's piecewise constant in time functions



$$\overline{u}_n(t) = u_i, \qquad t \in (t_{i-1}, t_i].$$

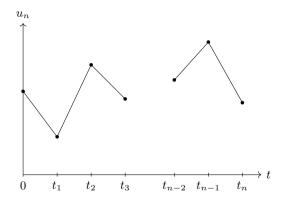


$$\underline{u}_n(t) = u_{i-1}, \qquad t \in [t_{i-1}, t_i)_{\text{fill}}$$

$$t \in [t_{i-1}, t_i)$$

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## Rothe's piecewise affine in time function



$$u_n(t) = u_{i-1} + (t - t_{i-1})\delta u_i, \qquad t \in [t_{i-1}, t_i].$$



## Discretized problem in the continous sense

The following identities are valid for any  $\psi \in Z, \varsigma \in L^2(\Sigma), \varphi \in W_0$  and  $w \in H^1(\Sigma)$ 

#### Existence of a weak solution

#### Theorem

There exists a weak solution

$$(\phi, z, u, \boldsymbol{A}) \in \operatorname{Lip}([0, T], \mathbf{Z}) \times \left[ \mathbf{C}([0, T], \mathbf{L}^2(\Sigma)) \cap \mathbf{L}^{\infty}((0, T), \mathbf{H}^1(\Sigma)) \right]^2 \times \mathbf{C}([0, T], \boldsymbol{W}_0)$$

to the variational system. Moreover, the following convergences hold for subsequences

#### Reference



V. C. Le, M. Slodička, and K. Van Bockstal.

Existence of a weak solution to a nonlinear induction hardening problem with Leblond-Devaux model for a steel workpiece.

Communications in Nonlinear Science and Numerical Simulation, 107:106156, 2022.





#### Conclusions

- A mathematical model has been introduced for the induction hardening process.
- A time discretization scheme has been proposed.
- Numerical analysis have been performed, and the existence of a weak solution has been proved.
- The uniqueness of a solution to the nonlinear induction hardening problem is still a challenge.





## THANK YOU FOR YOUR ATTENTION!

