A time discretization for solving an electromagnetic contact problem with moving conductor<sup>1</sup>

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 $^{1}$ This is a joint work with Marián Slodička and Karel Van Bockstal (Ghent University)

Mathematical model

- 2 Weak formulation
- **3** Time discretization
- Existence of a solution



Let us consider the eddy current equations

$$\nabla \cdot \boldsymbol{B} = 0,$$
  

$$\nabla \times \boldsymbol{E} = -\partial_t \boldsymbol{B},$$
  

$$\nabla \times \boldsymbol{H} = \boldsymbol{J},$$

with the constitutive relations

$$\boldsymbol{B} = \mu \boldsymbol{H}, \qquad \boldsymbol{J} = \sigma \left( \boldsymbol{E} + \mathbf{v} \times \boldsymbol{B} \right).$$

We acquire the following interface conditions

$$\begin{split} \llbracket \boldsymbol{B} \cdot \boldsymbol{n} \rrbracket &= \boldsymbol{0}, \\ \llbracket \boldsymbol{H} \times \boldsymbol{n} \rrbracket &= \boldsymbol{0}, \\ \llbracket (\boldsymbol{E} + \mathbf{v} \times \boldsymbol{B}) \times \boldsymbol{n} \rrbracket &= \boldsymbol{0}. \end{split}$$





# The jumping material coefficients

The material coefficients  $\sigma$  and  $\mu$  are defined by

$$\sigma(\boldsymbol{x},t) = \begin{cases} \sigma_{\Sigma}(\boldsymbol{x},t) & \text{ in } \Sigma(t), \\ \sigma_{\Pi}(\boldsymbol{x},t) & \text{ in } \Pi, \\ 0 & \text{ in } \Omega \setminus \overline{\Theta(t)}, \end{cases}$$

and

$$\mu(\boldsymbol{x},t) = \begin{cases} \mu_{\Sigma}(\boldsymbol{x},t) & \text{ in } \Sigma(t), \\ \mu_{\Pi}(\boldsymbol{x},t) & \text{ in } \Pi, \\ \mu_{air}(\boldsymbol{x},t) & \text{ in } \Omega \setminus \overline{\Theta(t)}. \end{cases}$$

At the interfaces of the different materials,  $\sigma$  and  $\mu$  have jumps, i.e.

$$\llbracket \sigma \rrbracket \neq 0, \qquad \qquad \llbracket \mu \rrbracket \neq 0$$

## $A - \phi$ formulation

Since  $\nabla \cdot \boldsymbol{B} = 0$ , there exists a unique vector potential  $\boldsymbol{A} \in \mathbf{H}^{1}(\Omega)$  of  $\boldsymbol{B}$  such that

$$\boldsymbol{B} = \nabla \times \boldsymbol{A}, \qquad \nabla \cdot \boldsymbol{A} = 0, \qquad \boldsymbol{A} \times \boldsymbol{n} = \boldsymbol{0} \quad \text{on} \quad \partial \Omega.$$

Therefore, from  $abla imes {m E} = -\partial_t {m B}$ , we get that

$$abla imes (oldsymbol{E} + \partial_t oldsymbol{A}) = oldsymbol{0}$$

or

$$\boldsymbol{E} + \partial_t \boldsymbol{A} = -\nabla\phi.$$

Thanks to  $J = \sigma(E + \mathbf{v} \times B)$ , the total current density J can be split into

$$\boldsymbol{J} = \boldsymbol{J}_{source} + \boldsymbol{J}_{eddy},$$

where

$$\boldsymbol{J}_{source} = -\sigma \nabla \phi, \qquad \quad \boldsymbol{J}_{eddy} = -\sigma \partial_t \boldsymbol{A} + \sigma \mathbf{v} \times (\nabla \times \boldsymbol{A}).$$



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# $A - \phi$ formulation (cont.)

The boundary value problem for the scalar potential  $\phi$ 

$$\begin{cases} \nabla \cdot (-\sigma \nabla \phi) = 0 & \text{ in } \Pi \times (0,T), \\ -\sigma \nabla \phi \cdot \mathbf{n} = 0 & \text{ on } (\partial \Pi \setminus \Gamma) \times (0,T), \\ -\sigma \nabla \phi \cdot \mathbf{n} = j & \text{ on } \Gamma \times (0,T) \,, \end{cases}$$

where

$$\int_{\Gamma} j(s,t) \, \mathrm{d}s = 0.$$

From abla imes H = J and  $B = \mu H$ , we derive the initial-boundary value problem for the vector potential A

$$\begin{cases} \sigma \partial_t \mathbf{A} + \nabla \times (\mu^{-1} \nabla \times \mathbf{A}) + \chi_{\Pi} \sigma \nabla \phi \\ -\sigma \mathbf{v} \times (\nabla \times \mathbf{A}) = \mathbf{0} & \text{ in } \Omega \times (0, T), \\ \mathbf{A} \times \mathbf{n} = \mathbf{0} & \text{ on } \partial \Omega \times (0, T), \\ \mathbf{A}(\mathbf{x}, 0) = \mathbf{A}_0 & \text{ in } \Theta(0). \end{cases}$$

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### Functional setting

Let us consider the following Hilbert spaces

$$\mathbf{Z} = \left\{ \boldsymbol{\psi} \in \mathbf{H}^1(\boldsymbol{\Pi}) : (\boldsymbol{\psi}, 1)_{\boldsymbol{\Pi}} = \mathbf{0} \right\},\$$

$$\boldsymbol{W}_0 = \left\{ \boldsymbol{\varphi} \in \mathbf{L}^2(\Omega) : \nabla \times \boldsymbol{\varphi} \in \mathbf{L}^2(\Omega), \nabla \cdot \boldsymbol{\varphi} = 0, \boldsymbol{\varphi} \times \mathbf{n} = \mathbf{0} \text{ on } \partial \Omega \right\},$$

with the equipped norms

$$\begin{split} \|\psi\|_{\mathbf{Z}} &= \|\nabla\psi\|_{\mathbf{L}^{2}(\Pi)} \,, \\ \|\varphi\|_{\boldsymbol{W}_{0}} &= \|\nabla\times\varphi\|_{\mathbf{L}^{2}(\Omega)} \end{split}$$

We recall the well-known Reynolds transport theorem

$$\frac{\mathrm{d}}{\mathrm{d}t} \int_{\omega(t)} f \,\mathrm{d}\boldsymbol{x} = \int_{\omega(t)} \partial_t f \,\mathrm{d}\boldsymbol{x} + \int_{\partial\omega(t)} f \,\mathbf{v} \cdot \mathbf{n} \,\mathrm{d}s.$$

# Weak formulation

The weak problem for  $\phi$  reads as:

$$\left(\sigma(t)\nabla\phi(t),\nabla\psi\right)_{\Pi}+\left(j(t),\psi\right)_{\Gamma}=0 \qquad \qquad \forall\psi\in\mathbf{Z}\,.$$

For any  $oldsymbol{arphi} \in oldsymbol{W}_0$ , it holds that

$$\begin{split} \left(\sigma(t)\partial_t \boldsymbol{A}(t), \boldsymbol{\varphi}\right)_{\Theta(t)} &+ \left(\mu^{-1}(t)\nabla \times \boldsymbol{A}(t), \nabla \times \boldsymbol{\varphi}\right)_{\Omega} \\ &+ \left(\sigma(t)\nabla\phi(t), \boldsymbol{\varphi}\right)_{\Pi} - \left(\sigma(t)\mathbf{v}(t) \times (\nabla \times \boldsymbol{A}(t)), \boldsymbol{\varphi}\right)_{\Theta(t)} = 0. \end{split}$$

Using the Reynolds transport theorem, the variational formulation for A is defined as: Find  $A(t) \in W_0$  such that the following identity holds true for any  $\varphi \in W_0$ 

$$\begin{split} \frac{\mathrm{d}}{\mathrm{d}t} \left( \boldsymbol{\sigma}(t) \boldsymbol{A}(t), \boldsymbol{\varphi} \right)_{\Theta(t)} &- \left( \boldsymbol{\sigma}(t) \boldsymbol{A}(t), \boldsymbol{\varphi}(\mathbf{v} \cdot \mathbf{n})(t) \right)_{\partial \Theta(t)} \\ &- \left( \boldsymbol{\gamma}(t) \boldsymbol{A}(t), \boldsymbol{\varphi} \right)_{\Theta(t)} + \left( \boldsymbol{\mu}^{-1}(t) \nabla \times \boldsymbol{A}(t), \nabla \times \boldsymbol{\varphi} \right)_{\Omega} \\ &+ \left( \boldsymbol{\sigma}(t) \nabla \boldsymbol{\phi}(t), \boldsymbol{\varphi} \right)_{\Pi} - \left( \boldsymbol{\sigma}(t) \mathbf{v}(t) \times (\nabla \times \boldsymbol{A}(t)), \boldsymbol{\varphi} \right)_{\Theta(t)} = 0, \\ & \underset{\mathsf{Output}}{\underset{\mathsf{Output}}{\underset{\mathsf{Output}}{\underset{\mathsf{Output}}{\underset{\mathsf{Output}}{\underset{\mathsf{Output}}{\underset{\mathsf{Output}}{\underset{\mathsf{Output}}{\underset{\mathsf{Output}}{\underset{\mathsf{Output}}{\underset{\mathsf{Output}}{\underset{\mathsf{Output}}{\underset{\mathsf{Output}}{\underset{\mathsf{Output}}{\underset{\mathsf{Output}}{\underset{\mathsf{Output}}{\underset{\mathsf{Output}}{\underset{\mathsf{O}(t)}}{\underset{\mathsf{O}(t)}{\underset{\mathsf{O}(t)}{\underset{\mathsf{O}(t)}{\underset{\mathsf{O}(t)}{\underset{\mathsf{O}(t)}{\underset{\mathsf{O}(t)}{\underset{\mathsf{O}(t)}{\underset{\mathsf{O}(t)}}{\underset{\mathsf{O}(t)}{\underset{\mathsf{O}(t)}}{\underset{\mathsf{O}(t)}}{\underset{\mathsf{O}(t)}}}}}}}}}}}{}}}}}}}$$

where  $\gamma = \partial_t \sigma$ .

### Theorem

Let the material functions  $\mu \in C^1$  and  $\sigma \in C^2$  on each component (in all variables). Moreover, assume that  $A_0 \in L^2(\Theta_0)$ ,  $j \in Lip([0,T], H^{-1/2}(\Gamma))$ ,  $\mathbf{v} \in \mathbf{C}^1(\overline{\Omega} \times [0,T])$  and

$$0 < \sigma_* \le \sigma \le \sigma^*$$
 in  $\Theta$ ;  $\sigma = 0$  in air

and

$$0 < \mu_* \le \mu \le \mu^*$$
 in  $\Omega$ .

The variational system admits at most one solution  $(\phi, \mathbf{A})$  satisfying  $\phi \in L^2((0,T), \mathbb{Z})$ ,  $\mathbf{A} \in L^2((0,T), \mathbf{W}_0)$  and  $\|\sqrt{\sigma}\mathbf{A}\|_{\mathbf{L}^2(\Omega)} \in C([0,T])$ .



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We partition the time range [0,T] into  $n \in \mathbb{N}$  equidistant subintervals with the time step

$$au = \frac{T}{n}.$$

The time discretization scheme is defined as: Find  $\phi_i \in \mathbb{Z}$  and  $A_i \in W_0$  such that the following identities are valid for any  $\psi \in \mathbb{Z}$  and  $\varphi \in W_0$ 

$$\begin{aligned} (\sigma_i \nabla \phi_i, \nabla \psi)_{\Pi} + (j_i, \psi)_{\Gamma} &= 0, \\ \delta \left( \sigma_i \boldsymbol{A}_i, \boldsymbol{\varphi} \right)_{\Theta_i} - \left( \sigma_i \boldsymbol{A}_i, \boldsymbol{\varphi} (\mathbf{v}_i \cdot \mathbf{n}) \right)_{\partial \Theta_i} - (\gamma_i \boldsymbol{A}_i, \boldsymbol{\varphi})_{\Theta_i} \\ &+ \left( \mu_i^{-1} \nabla \times \boldsymbol{A}_i, \nabla \times \boldsymbol{\varphi} \right)_{\Omega} + (\sigma_i \nabla \phi_i, \boldsymbol{\varphi})_{\Pi} - \left( \sigma_i \mathbf{v}_i \times (\nabla \times \boldsymbol{A}_i), \boldsymbol{\varphi} \right)_{\Theta_i} = 0, \end{aligned}$$

where

$$\delta\left(\sigma_{i}\boldsymbol{A}_{i},\boldsymbol{\varphi}\right)_{\Theta_{i}}=\frac{1}{\tau}\left(\left(\sigma_{i}\boldsymbol{A}_{i},\boldsymbol{\varphi}\right)_{\Theta_{i}}-\left(\sigma_{i-1}\boldsymbol{A}_{i-1},\boldsymbol{\varphi}\right)_{\Theta_{i-1}}\right).$$

#### Lemma

For any  $\tau < \tau_0$  and for any i = 1, 2, ..., n, there exists a unique couple  $(\phi_i, \mathbf{A}_i) \in \mathbb{Z} \times \mathbf{W}_0$  solving the time discrete system.

#### Lemma

There exist positive constants C and  $\tau_0$  such that the following estimates hold true for any  $\tau<\tau_0$ 

(i) 
$$\max_{1 \le i \le n} \|\nabla \phi_i\|_{\mathbf{L}^2(\Pi)}^2 \le C,$$
  
(ii) 
$$\max_{1 \le l \le n} \|A_l\|_{\mathbf{L}^2(\Theta_l)}^2 + \sum_{i=1}^n \|\nabla \times A_i\|_{\mathbf{L}^2(\Omega)}^2 \tau + \sum_{i=1}^n \|A_i - A_{i-1}\|_{\mathbf{L}^2(\Theta_{i-1})}^2 \le C.$$

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### Rothe's method

### We introduce the following Rothe's functions and domain

$$\overline{\phi}_n(t) = \phi_i, \quad \overline{A}_n(t) = A_i,$$

$$g_n(t) = (\sigma_{i-1}A_{i-1}, \varphi)_{\Theta_{i-1}} + (t - t_{i-1}) \,\delta \left(\sigma_i A_i, \varphi\right)_{\Theta_i},$$

$$\overline{j}_n(t) = j_i, \quad \overline{\sigma}_n(t) = \sigma_i, \quad \overline{\gamma}_n(t) = \gamma_i, \quad \overline{\mu}_n(t) = \mu_i, \quad \overline{\mathbf{v}}_n(t) = \mathbf{v}_i, \quad \overline{\Theta}_n(t) = \Theta_i,$$

with the initial data  $\overline{\phi}_n(0) = \phi_0, \overline{A}_n(0) = A_0, g_n(0) = (\sigma_0 A_0, \varphi)_{\Theta_0}$ . The discrete problems can be rewritten in the continuous sense: for any  $\psi \in \mathbb{Z}$  and  $\varphi \in W_0$ , it holds that

$$\begin{split} \left(\overline{\sigma}_{n}(t)\nabla\overline{\phi}_{n}(t),\nabla\psi\right)_{\Pi} &+ \left(\overline{j}_{n}(t),\psi\right)_{\Gamma} = 0,\\ g_{n}'(t) &- \left(\overline{\sigma}_{n}(t)\overline{A}_{n}(t),\varphi(\overline{\mathbf{v}}_{n}\cdot\mathbf{n})(t)\right)_{\partial\overline{\Theta}_{n}(t)} - \left(\overline{\gamma}_{n}(t)\overline{A}_{n}(t),\varphi\right)_{\overline{\Theta}_{n}(t)} \\ &+ \left(\overline{\mu}_{n}^{-1}(t)(\nabla\times\overline{A}_{n}(t)),\nabla\times\varphi\right)_{\Omega} + \left(\overline{\sigma}_{n}(t)\nabla\overline{\phi}_{n}(t),\varphi\right)_{\Pi} \\ &- \left(\overline{\sigma}_{n}(t)\overline{\mathbf{v}}_{n}(t)\times(\nabla\times\overline{A}_{n}(t)),\varphi\right)_{\overline{\Theta}_{n}(t)} = 0. \end{split}$$

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### Lemma

There exists a constant C > 0 such that the following relations hold true

$$\begin{split} \max_{t\in[0,T]} \left\| \overline{j}_n(t) - j(t) \right\|_{\mathbf{H}^{-1/2}(\Gamma)} &\leq C\tau, \\ \max_{t\in[0,T]} \left\| \overline{\mathbf{v}}_n(t) - \mathbf{v}(t) \right\|_{\mathbf{H}(\mathsf{div},\Omega)} &\leq C\tau, \\ \max_{t\in[0,T]} \left\| \overline{\mu}_n(t) - \mu(t) \right\|_{\mathbf{L}^2(\Omega)} &\leq C\sqrt{\tau}, \\ \max_{t\in[0,T]} \left\| \overline{\gamma}_n(t) - \gamma(t) \right\|_{\mathbf{L}^2(\Omega)} &\leq C\sqrt{\tau}, \\ \max_{t\in[0,T]} \left\| \overline{\sigma}_n(t) - \sigma(t) \right\|_{\mathbf{L}^2(\Omega)} &\leq C\sqrt{\tau}, \\ \max_{t\in[0,T]} \left\| \overline{\nabla}\overline{\sigma}_n(t) - \nabla\sigma(t) \right\|_{\mathbf{L}^2(\Omega)} &\leq C\sqrt{\tau}. \end{split}$$

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### Theorem

There exists a unique weak solution  $(\phi, \mathbf{A})$  to the variational system satisfying  $\phi \in L^2((0,T), \mathbb{Z})$  and  $\mathbf{A} \in L^2((0,T), \mathbf{W}_0)$  with  $\|\sqrt{\sigma}\mathbf{A}\|_{L^2(\Omega)} \in C([0,T])$ . Moreover, the following convergences hold true

(i)	$\overline{\phi}_n \rightharpoonup \phi$	in	$\mathbf{L}^{2}\left((0,T),\mathbf{Z}\right),$
(ii)	$\overline{oldsymbol{A}}_n  ightarrow oldsymbol{A}$	in	$\mathrm{L}^{2}\left((0,T),\boldsymbol{W}_{0}\right),$
	$g_n  ightarrow (\sigma {oldsymbol A}, {oldsymbol arphi})_{\Theta}$	in	$\mathrm{L}^{1}\left((0,T)\right).$



*Proof.* The boundedness of  $\{\overline{\phi}_n\}$  and  $\{\overline{A}_n\}$  implies the existence of  $\phi \in L^2((0,T), \mathbb{Z})$  and  $A \in L^2((0,T), W_0)$  such that

$$\begin{split} \overline{\phi}_n &\rightharpoonup \phi & \qquad & \text{in} \quad \operatorname{L}^2\left((0,T),\operatorname{Z}\right), \\ \overline{A}_n &\rightharpoonup A & \qquad & \text{in} \quad \operatorname{L}^2\left((0,T), W_0\right). \end{split}$$

Thanks to the stability of  $\overline{A}_n$ , the following strong convergence holds true

$$g_n - \left(\overline{\sigma}_n \overline{A}_n, \varphi\right)_{\overline{\Theta}_n} \to 0 \qquad \text{in} \quad \mathrm{L}^1((0,T)).$$

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# Proof of the existence (cont.)

Moreover, we can also deduce the following convergences in  $L^1((0,T))$ 

$$\begin{split} & \left(\overline{\gamma}_{n}\overline{\boldsymbol{A}}_{n},\boldsymbol{\varphi}\right)_{\overline{\Theta}_{n}} \rightarrow \left(\gamma\boldsymbol{A},\boldsymbol{\varphi}\right)_{\Theta}, \\ & \left(\overline{\mu}_{n}^{-1}\nabla\times\overline{\boldsymbol{A}}_{n},\nabla\times\boldsymbol{\varphi}\right)_{\Omega} \rightarrow \left(\mu^{-1}\nabla\times\boldsymbol{A},\nabla\times\boldsymbol{\varphi}\right)_{\Omega}, \\ & \left(\left(\nabla\overline{\sigma}_{n}\cdot\overline{\mathbf{v}}_{n}\right)\overline{\boldsymbol{A}}_{n},\boldsymbol{\varphi}\right)_{\overline{\Theta}_{n}} \rightarrow \left(\left(\nabla\sigma\cdot\mathbf{v}\right)\boldsymbol{A},\boldsymbol{\varphi}\right)_{\Theta}, \\ & \left(\overline{\sigma}_{n}\left(\nabla\cdot\overline{\mathbf{v}}_{n}\right)\overline{\boldsymbol{A}}_{n},\boldsymbol{\varphi}\right)_{\overline{\Theta}_{n}} \rightarrow \left(\sigma\left(\nabla\cdot\mathbf{v}\right)\boldsymbol{A},\boldsymbol{\varphi}\right)_{\Theta}, \\ & \left(\overline{\sigma}_{n}\overline{\mathbf{v}}_{n},\nabla\left(\overline{\boldsymbol{A}}_{n}\cdot\boldsymbol{\varphi}\right)\right)_{\overline{\Theta}_{n}} \rightarrow \left(\sigma\mathbf{v},\nabla\left(\boldsymbol{A}\cdot\boldsymbol{\varphi}\right)\right)_{\Theta}, \\ & \left(\overline{\sigma}_{n}\overline{\mathbf{v}}_{n}\times\left(\nabla\times\overline{\boldsymbol{A}}_{n}\right),\boldsymbol{\varphi}\right)_{\overline{\Theta}_{n}} \rightarrow \left(\sigma\mathbf{v}\times\left(\nabla\times\boldsymbol{A}\right),\boldsymbol{\varphi}\right)_{\Theta}. \end{split}$$

Therefore, we can pass to the limit and use the Reynolds transport theorem to conclude that  $(\phi, A)$  solves the variational system.

Finally, the following relation can be obtained

$$\left| (\sigma \boldsymbol{A}, \boldsymbol{A})_{\Theta} (\eta) - (\sigma \boldsymbol{A}, \boldsymbol{A})_{\Theta} (\xi) \right| \lesssim |\eta - \xi| + \int_{\xi}^{\eta} \| \nabla \times \boldsymbol{A} \|_{\mathbf{L}^{2}(\Omega)}^{2} dt$$

which implies the continuity in time, i.e.

$$\lim_{\eta \to \xi} \left( \sigma \boldsymbol{A}, \boldsymbol{A} \right)_{\Theta} (\eta) = \left( \sigma \boldsymbol{A}, \boldsymbol{A} \right)_{\Theta} (\xi)$$

and thus  $\left\|\sqrt{\sigma} \boldsymbol{A}\right\|_{\mathbf{L}^{2}(\Omega)} \in \mathrm{C}([0,T]).$ 



- A mathematical model for an electromagnetic contact problem with moving conductor is investigated.
- A time discretization for solving the variational problem is introduced.
- The convergence of the numerical scheme is shown.
- The existence of a unique weak solution is present.

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### V. C. Le, M. Slodička, and K. Van Bockstal.

A time discrete scheme for an electromagnetic contact problem with moving conductor.

Applied Mathematics and Computation, 404:125997, sep 2021.



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# Thank you for your attention!



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