

H₂O emission in high-*z* ultra-luminous infrared galaxies^{*}

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ABSTRACT

Using the IRAM Plateau de Bure interferometer (PdBI), we report the detection of water vapor in six new lensed ultra-luminous starburst galaxies at high redshift, discovered in the *Herschel* Astrophysical Terahertz Large Area Survey (H-ATLAS). The sources are detected either in the 2₀₂–1₁₁ or 2₁₁–2₀₂ H₂O emission lines with integrated line fluxes ranging from 1.8 to 14 Jy km s⁻¹. The corresponding apparent luminosities are $\mu L_{\text{H}_2\text{O}} \sim 3\text{--}12 \times 10^8 L_{\odot}$, where μ is the lensing magnification factor ($3 < \mu < 12$). These results confirm that H₂O lines are among the strongest molecular lines in high-*z* ultra-luminous starburst galaxies, with intensities almost comparable to those of the high-*J* CO lines, and similar profiles and line widths ($\sim 200\text{--}900$ km s⁻¹). With the current sensitivity of the PdBI, the water lines can therefore easily be detected in high-*z* lensed galaxies (with $F(500 \mu\text{m}) > 100$ mJy) discovered in the *Herschel* surveys. Correcting the luminosities for amplification, using existing lensing models, $L_{\text{H}_2\text{O}}$ is found to have a strong dependence on the infrared luminosity, varying as $\sim L_{\text{IR}}^{1.2}$. This relation, which needs to be confirmed with better statistics, may indicate a role of radiative (infrared) excitation of the H₂O lines, and implies that high-*z* galaxies with $L_{\text{IR}} \gtrsim 10^{13} L_{\odot}$ tend to be very strong emitters in water vapor, that have no equivalent in the local universe.

Key words. galaxies: high-redshift – galaxies: starburst – infrared: galaxies – submillimeter: galaxies – radio lines: galaxies

1. Introduction

Water plays an important role in the warm dense interstellar medium of galaxies. First, after CO, H₂O is the most abundant oxygen-bearing molecule, and, second, it can be an important coolant of the warm gas. Due to the Earth's atmosphere, bulk gas-phase water can only be detected from space or from the ground toward distant objects with redshifts that move the H₂O lines into atmospheric windows.

The study of water emission lines in nearby galaxies has recently made significant progress thanks to the availability of their infrared/submillimeter spectra using the spectrometer mode of the *Herschel* Spectral and Photometric Imaging REceiver (SPIRE, Griffin et al. 2010) and the Photodetecting Array Camera and Spectrometer (PACS, Poglitsch et al. 2010). The spectra of local ultra-luminous infrared galaxies (ULIRGs) and composite AGN/starbursts such as Mrk 231 (van der Werf et al. 2010), Arp 220 (Rangwala et al. 2011; González-Alfonso et al. 2012), and NGC 4418 (González-Alfonso et al. 2012) reveal a wealth of water lines and the presence of associated molecules such as OH⁺, H₂O⁺ and isotopologues. Together with high-*J* CO lines, these lines provide an important diagnostic of the warm dense cores of nearby ULIRGs.

In the cases of Mrk 231 and Arp 220, water emission lines up to energy levels of $E_{\text{up}}/k = 642$ K are detected with strong

line fluxes that reach 25%–75% of the neighboring CO emission line fluxes. Spectral surveys made with *Herschel* show that low-*z* ULIRGs always exhibit bright H₂O lines, whereas only one third of the sample of luminous infrared galaxies (LIRGs) displays luminous H₂O emission lines (van der Werf et al., in prep.), indicating that the strength of the water lines and the infrared luminosity, L_{IR} , must be related. The analysis of the H₂O emission lines in Mrk 231 shows that the excitation of the water molecules results from a combination of collisions and infrared excitation through far-infrared lines in warm dense gas (≥ 100 K, $\geq 10^5$ cm⁻³). Moreover, the far-infrared radiation field dominates the excitation of the high levels and their emission lines (González-Alfonso et al. 2010). Preliminary evidence from the comparison of the spectra of Arp 220, Mrk 231 and NGC 4418 shows that the properties of the water emission lines in their nuclear regions vary as a function of chemistry, nucleosynthesis, and inner motions (outflow/infall) – see Rangwala et al. (2011) and González-Alfonso et al. (2012).

First detections of H₂O megamasers at high redshift were reported in objects at $z = 0.66$ and 2.64 by Barvainis & Antonucci (2005) and Impellizzeri et al. (2008), while further searches have failed since then (McKean et al. 2011). Following initial attempts to detect H₂O rotational emission lines from high-*z* galaxies (Wagg et al. 2006; Riechers et al. 2006, 2009) and tentative detections in IRAS F10214+4724 at $z = 2.23$ (Casoli et al. 1994) and in the Cloverleaf at $z = 2.56$ (Bradford et al. 2009), a series of robust detections of non-maser H₂O emission lines were recently reported in high-*z* sources. The sources

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Table 1. Observation Log.

IAU name	ID	RA ^a	Dec ^b	Date	Frequency ^b	Beam (")	$t_{\text{on}}(h)^c$
HATLAS J090311.6+003906	SDP.81	09:03:11.61	+00:39:06.7	2011 Dec.	244.890	2.8×1.8	1.2
HATLAS J133649.9+291801	NA.v1.144	13:36:50.00	+29:17:59.6	2012 Mar.	235.330	1.6×1.1	0.8
HATLAS J090740.0-004200	SDP.9	09:07:40.05	-00:41:59.5	2011 Dec.	292.550	3.3×1.6	1.4
HATLAS J133008.4+245900	NB.v1.78	13:30:08.56	+24:58:58.3	2012 Jan. & Mar.	240.802	1.4×1.0	2.3
HATLAS J114637.9-001132	G12.v2.30	11:46:37.99	-00:11:32.0	2011 Mar.	232.734	2.0×1.0	1.2
HATLAS J142413.9+022303	G15.v2.779	14:24:13.90	+02:23:04.0	2011 Dec.	143.940	4.2×3.1	1.7
HATLAS J090302.9-014127	SDP.17b*	09:03:03.02	-01:41:26.9	2011 Jan.	298.148	2.9×1.1	0.6

Notes. ^(a) J2000 coordinates of the centers of the maps displayed in Fig. 2. ^(b) Central observed frequency (GHz). ^(c) Total on-source integration time for the PdBI array with six antennas. ^(*) The observations of SDP.17b are reported in Omont et al. (2011).

detected in H₂O emission lines include a strongly lensed galaxy, HATLAS J090302.9-014127 (SDP.17b) at $z = 2.3$ (Omont et al. 2011), which was uncovered in the *Herschel* Astrophysical Terahertz Large Area Survey (H-ATLAS; Eales et al. 2010); the gravitationally lensed quasar APM 08279+5255 at $z = 3.9$ (van der Werf et al. 2011; Bradford et al. 2011; Lis et al. 2011); HLSJ091828.6+514223, a lensed submillimeter galaxy at $z = 5.2$ in the field of Abell 773 (Combes et al. 2012), which was found in the *Herschel* Lensing Survey (Egami et al. 2010), IRAS F10214+4724 (Riechers et al., in prep.). For most of these sources, only one water emission line was reported, with the exception of APM 08279+5255, where a total of at least five emission lines with $E_u/k = 101$ to 454 K were detected (van der Werf et al. 2011; Lis et al. 2011; Bradford et al. 2011). These results underline the unique and powerful diagnostic power of H₂O emission lines, which give better insight into local conditions in distant galaxies than may be obtained by other means. As in Mrk 231, they reveal the presence of extended, warm and dense gas located in the infrared-opaque regions of the galactic cores.

This new window in the exploration of high- z sources is based on the combined availability of new instrumentation with improved sensitivities at key facilities and the increasing number of gravitationally lensed sources discovered in the *Herschel* and South Pole Telescope (SPT) cosmological surveys (see, e.g., Negrello et al. 2010; Vieira et al. 2010). Following the results reported by Omont et al. (2011), we present here a new study of water emission in six high- z lensed ULIRGs, which were selected from the H-ATLAS survey. The data, which were obtained using the IRAM Plateau de Bure Interferometer (PdBI), clearly show water emission lines in all the sources. Based on these results, we derive a clear relation between the H₂O and the infrared luminosities in high- z ULIRGs.

Throughout this paper, we adopt a cosmology with $H_0 = 71 \text{ km s}^{-1} \text{ Mpc}^{-1}$, $\Omega_M = 0.27$, $\Omega_\Lambda = 0.73$ (Spergel et al. 2003).

2. Sample selection and observations

The detection of strong H₂O emission in the lensed H-ATLAS galaxy SDP.17b by Omont et al. (2011) suggested that the H-ATLAS survey offered a unique opportunity to select an homogeneous sample of bright lensed galaxies spanning a broad range in luminosity and redshift to further study the properties of water emission in the early universe. We selected from the H-ATLAS catalog sources with strong flux densities (with $F_{500 \mu\text{m}} > 150 \text{ mJy}$) that were well characterized, i.e. had CO redshift measurements, additional sub/millimeter imaging and available deflector identification (see, e.g., Negrello et al. 2010; Bussmann et al., in prep.). In the selection, we also

somewhat preferred sources at $z > 3$ to shift the H₂O lines in the low frequency bands that are easier to observe. Finally, we chose sources spanning a wide range in intrinsic infrared luminosities from $\sim 5 \times 10^{12}$ to $\sim 2 \times 10^{13} L_\odot$.

In this paper, we report the first results of this survey and describe the properties of the H₂O emission for six new sources. The sample includes (see Table 1 for IAU names of the sources) i) sources reported in the initial H-ATLAS science demonstration phase paper by Negrello et al. (2010) – SDP.9 and SDP.81, together with SDP.17b, which was discussed in Omont et al. (2011); ii) two other well-studied sources from the H-ATLAS equatorial fields – G12.v2.30 (Fu et al. 2012) and G15.v2.779 (Cox et al. 2011; Bussmann et al. 2012); iii) two sources from the H-ATLAS NGP field, NA.v1.144 and NB.v1.78, for which CO observations (Harris et al. 2012; Riechers et al., in prep.; and this paper) and submillimeter imaging (Bussmann et al., in prep.) are available. Table 1 provides details on the H₂O observations of the seven H-ATLAS lensed galaxies, Table 2 lists their submillimeter and infrared properties together with the estimated amplification factors (see Sects. 3.2 and 4.1), and Table 3 gives their CO properties available from recent measurements using the PdBI.

The redshifts of the seven selected galaxies range from 1.57 to 4.24, with the majority in the range $2.0 < z < 3.3$. In the sample, G15.v2.779 is the only source at $z > 4$ and, for the time being, the highest redshift lensed galaxy spectroscopically confirmed in the H-ATLAS survey. For each source, at least one of the two strongest low-excitation H₂O lines is in an atmospheric window observable with the PdBI, either H₂O(2₁₁-2₀₂) with $\nu_{\text{rest}} = 752.033 \text{ GHz}$ and $E_{\text{up}} = 137 \text{ K}$, or H₂O(2₀₂-1₁₁) with $\nu_{\text{rest}} = 987.927 \text{ GHz}$ and $E_{\text{up}} = 101 \text{ K}$. As both lines have comparable intensities in Arp 220 and Mrk 231 (see Sect. 4), in the rare cases where both lines are observable, we chose the line at the lower frequency that is easier to observe, except in the case of SDP.17b (Omont et al. 2011), where the line H₂O(2₀₂-1₁₁) was selected to confirm the tentative detection that was reported by Lupu et al. (2012). Note that these two observed lines are both para-H₂O and they are adjacent lines in the H₂O level diagram.

The H₂O observations were conducted in the compact D-configuration from December 2011 to March 2012 in conditions of good atmospheric phase stability (seeing of 1") and reasonable transparency (PWV $\leq 1 \text{ mm}$). Except for SDP.81, which was observed with five antennas, all other sources were observed with the six antennas of the PdBI array. The total on-source integration ranges from ~ 1 to 2 hours (Table 1). We used Bands 2, 3 and 4, which cover the frequency ranges 129–174 GHz, 201–267 GHz and 277–371 GHz, and the band centers were tuned to the redshifted frequency of the

Table 2. Submillimeter and infrared properties of the lensed ultra-luminous galaxies.

Source ID	F_{250} (mJy)	F_{350} (mJy)	F_{500} (mJy)	μL_{IR} ($10^{13} L_{\odot}$)	μ	L_{IR} ($10^{12} L_{\odot}$)
SDP.81	129 ± 20	182 ± 28	166 ± 27	5.8	9.5 ± 0.8	6.1
NA.v1.144	295 ± 45	294 ± 45	191 ± 31	5.7	5.3 ± 2.9	11
SDP.9	485 ± 73	323 ± 49	175 ± 28	4.4	8.5 ± 1.8	5.2
NB.v1.78	273 ± 42	282 ± 43	214 ± 33	11	10.5 ± 1.4	10
G12.v2.30	323 ± 49	378 ± 57	298 ± 45	15.7	9.6 ± 0.9	16
G15.v2.779	115 ± 19	192 ± 30	204 ± 32	8.5	4.1 ± 0.2	21
SDP.17b	328 ± 50	308 ± 47	220 ± 34	6.9	4.3 ± 1.2	16

Notes. F_{250} , F_{350} , and F_{500} are the SPIRE flux densities at 250, 350, and 500 μm ; μL_{IR} is the apparent total infrared luminosity (8–1000 μm); μ is the adopted lensing magnification factor derived from lensing modeling of 880 μm continuum, mostly from Bussmann et al. (in prep.; see text Sect. 3 for details); L_{IR} is the intrinsic infrared luminosity, derived from the literature (Harris et al. 2012; Lupu et al. 2012; Fu et al. 2012; Bussmann et al. 2012) except for NB.v1.78, for which an approximate value is inferred from scaling the value of G12.v2.30 with the ratio of the 350 μm flux densities.

Table 3. Observed parameters of the CO emission lines.

Source ID	z	CO line	I_{CO} (Jy km s ⁻¹)	ΔV_{CO} (km s ⁻¹)	Other CO data
SDP.81	3.040	5–4	7.0 ± 0.4	560 ± 40	Harris et al. (2012)
NA.v1.144	2.202	4–3	12.3 ± 0.6	210 ± 10	Harris et al. (2012)
SDP.9	1.574	3–2	9.2 ± 0.3	420 ± 10	Iono et al. (2012)
NB.v1.78	3.111	3–2	3.3 ± 0.5	560 ± 70	Riechers et al. (in prep.)
G12.v2.30	3.259	3–2	7.8 ± 1.7	650 ± 160	Fu et al. (2012)
G15.v2.779	4.244	7–6	7.5 ± 0.6	840 ± 50	Cox et al. (2011)
SDP.17b	2.305	4–3	9.1 ± 0.3	320 ± 10	Harris et al. (2012)

Notes. z is the redshift measured from the PdBI CO spectra (Fig. 1); I_{CO} is the integrated flux of the CO line; ΔV_{CO} is the CO linewidth (FWHM); references to other CO measurements from the literature are given in the last column.

selected H₂O emission line using the redshifts estimated from CO observations (Table 3). The correlator (WideX) provided a contiguous frequency coverage of 3.6 GHz in dual polarization with a fixed channel spacing of 1.95 MHz, allowing us to detect the continuum as well as any additional emission lines, if present. In the compact configuration, the baselines extend from 24 to 179 m, resulting in synthesized beams of $\sim 1.5'' \times 1.0''$ to $\sim 4'' \times 3''$ (Table 1). Only SDP.81, G12.v2.30 and NB.v1.78 are resolved at this angular resolution (Fig. 2).

During the observations, the phase and bandpass were calibrated by measuring standards sources that are regularly monitored at the IRAM PdBI, including 3C84, 3C279, MWC349, CRL618, and 0851+202. The accuracy of the flux calibration is estimated to range from $\sim 10\%$ at 2 mm to $\sim 20\%$ at 0.8 mm.

To complement the H₂O data, we include in this paper new results on the CO emission for the seven sources. For five of them, the CO data were obtained in 2011 and 2012 using the PdBI in the A-configuration. Those data were acquired in the frame of a survey of lensed ULIRGs to map their CO and dust continuum emission at high-angular resolution (Cox et al., in prep.) and a full description of the results will be given in that paper. For two of the sources, NB.v1.78 and G12.v2.30, the CO data were obtained in August 2012 in the D-configuration with five antennas. The observations were made in good weather conditions and more details are provided in Table 3. In this paper, we only present the global CO spectra (i.e., integrated over the source’s extent) with the goal to compare the characteristics of the H₂O and CO emission line profiles. A detailed discussion of the morphology, the dynamics, and the lensing of these sources will be given in Cox et al. (in prep.).

3. Properties of the H₂O emission lines

3.1. General properties

The seven lensed high-*z* ULIRGs are all detected with high signal-to-noise ratios ($S/N \gtrsim 6$), except SDP.81, either in the $2_{02}-1_{11}$ or the $2_{11}-2_{02}$ H₂O emission line (see Fig. 1), as well as in the underlying redshifted submillimeter continuum emission ($S/N > 20$). The water emission lines are strong, with integrated fluxes ranging from 1.8 to 14 Jy km s⁻¹ (Fig. 1 and Table 4). This indicates that in high-*z* ULIRGs, they are the strongest molecular submillimeter emission lines after those of high-*J* CO (see Sect. 4.3; note, however, that 2 mm and 3 mm CO lines displayed in Fig. 1 may have weaker intensity than the 1 mm H₂O lines).

The H₂O linewidths (FWHM) range from 140 to 900 km s⁻¹ (Table 4). The lines have a variety of profiles including single Gaussian profiles (NA.v1.144, NB.v1.78), double-peaked profiles (G12.v2.779), or asymmetrical profiles (SDP.9, SDP.17b, SDP.81). The H₂O and CO line profiles share similar properties (shape, linewidth) in all sources, suggesting that there is no strong differential lensing effects between the CO and the H₂O emission lines.

From the spatially integrated H₂O line flux, $I_{\text{H}_2\text{O}}$, the apparent H₂O luminosity, $\mu L_{\text{H}_2\text{O}}$ where μ is the lensing magnification factor, can be derived using the relations given in, e.g., Solomon et al. (1992). $\mu L_{\text{H}_2\text{O}}$ varies by only a factor of ~ 4 from SDP.81 to NB.v1.78 or G12.v2.30 (Table 4). Figure 3 displays the relation between $\mu L_{\text{H}_2\text{O}}$ and μL_{IR} , including the values for local ULIRGs (Yang et al., in prep.). The infrared and H₂O apparent luminosities are one or two orders of magnitude higher in these high-*z* galaxies than in local ULIRGs. However, the ratios

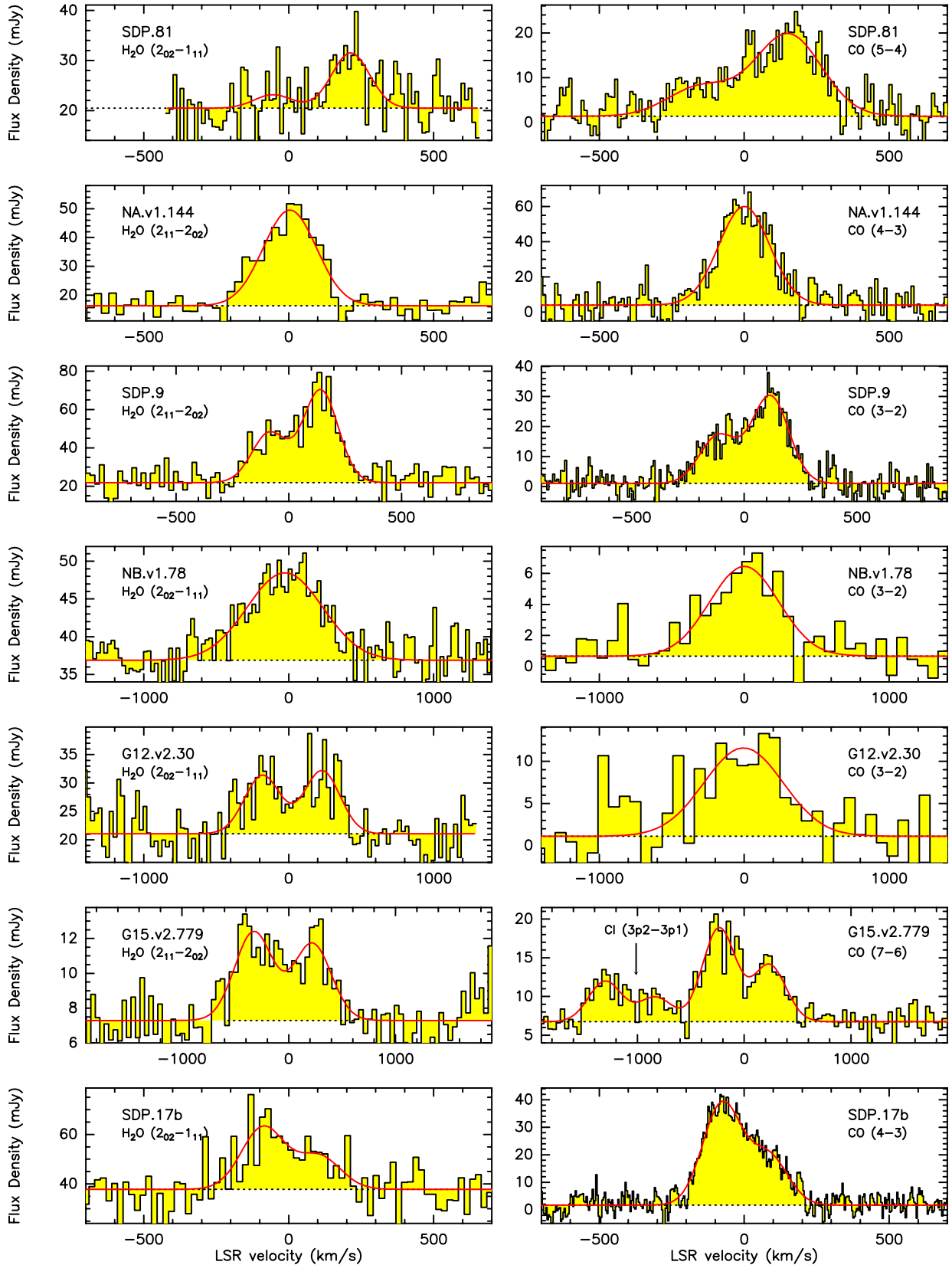


Fig. 1. Spectra of the H₂O and CO emission lines observed at PdBI toward the seven lensed high-redshift ultra-luminous galaxies discussed in this paper. The H₂O emission lines are displayed in the left column (showing the $(2_{11}-2_{02})$ or the $(2_{02}-1_{11})$ transition; the transitions are identified in each box). The CO spectra are displayed in the right column (showing 3–2, 4–3, 5–4 and 7–6 rotational transitions – for G15.v.2.779, the [CI] emission line is included in the band). The transitions are identified in each box. All spectra include the continuum, and the H₂O frequencies corresponding to the zero velocity are given in Table 4. The red lines trace the Gaussian fits to the H₂O and CO line profiles and the fit to the continuum emission level. In all cases, the match between the H₂O and CO spectra is excellent both in redshift and in the details of the line profile – see text.

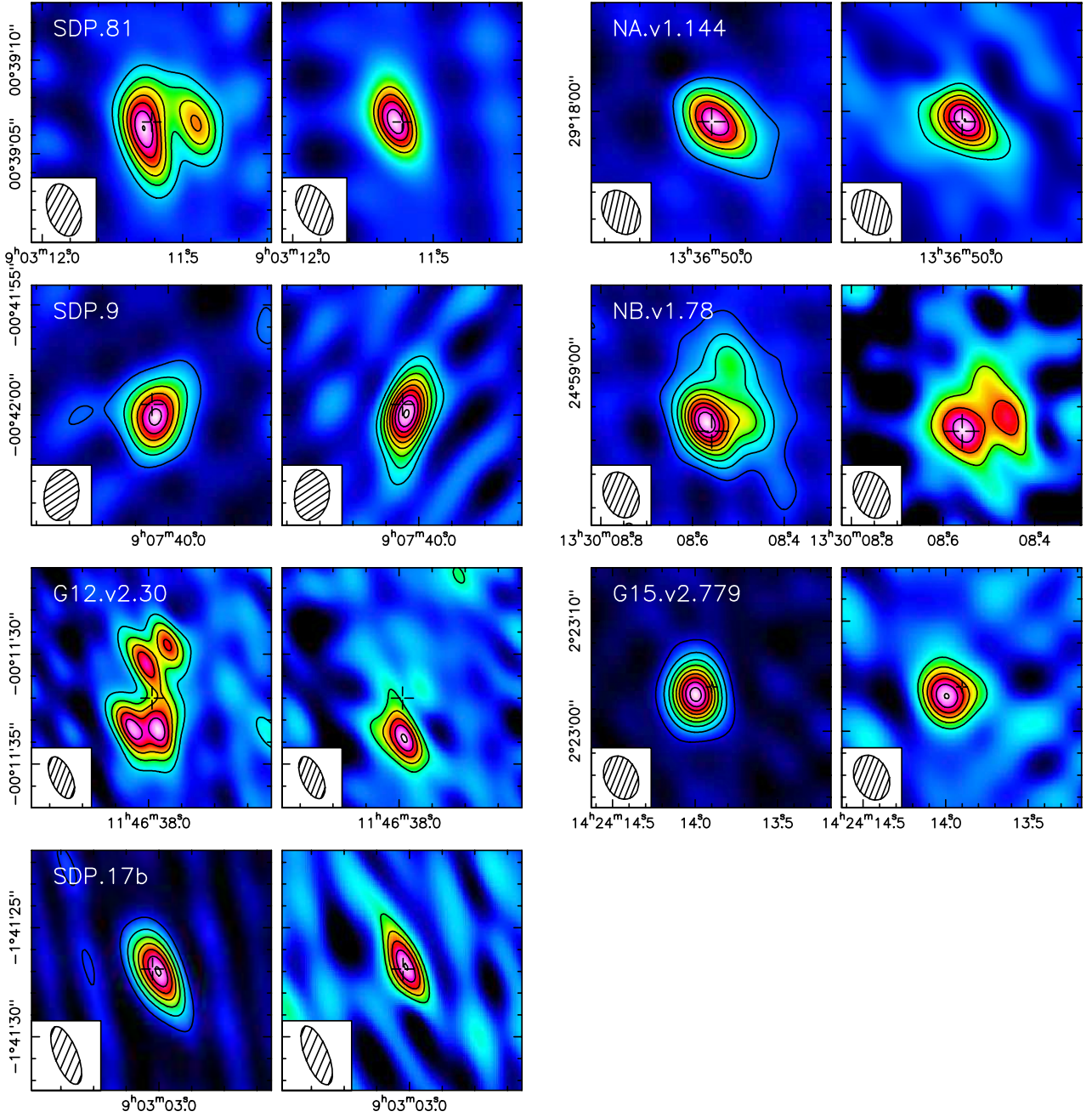


Fig. 2. Images of the submillimeter continuum (*left panel*) and H₂O line emission (*right panel* – see Fig. 1 and Table 4 for the identification of the transitions) for the seven high-*z* lensed ultra-luminous galaxies studied in this paper (sources are identified in the *left panels*). Synthesized beams are shown in the lower left corner of each map. The map centers correspond to the coordinates given in Table 1 and are indicated by a cross. For the continuum maps, the contours start from 3σ in steps of 3σ for SDP.81 and G12.v2.30, and in steps of 6σ for the other sources; for the H₂O line emission maps, the contours start from 3σ in steps of 1σ (for SDP.81) and 2σ for the other sources. For each source, the 1σ continuum (mJy/beam) and emission line (Jy km s⁻¹/beam) noise levels are as follows: SDP.81 (0.60/0.18), NA.v1.144 (0.35/0.31), SDP.9 (0.50/0.59), NB.v1.78 (0.28/0.30), G12.v2.30 (0.28/0.28), G15.v2.779 (0.11/0.24), and SDP.17b (0.73/0.56). Four sources remain unresolved at the current angular resolution, whereas SDP.81, NB.v1.78, and G12.v2.30 display complex lensing morphologies both in the continuum and line emission that are revealed even at the present low angular resolution. Detailed comments on the sources are given in the text (Sect. 3).

$L_{\text{H}_2\text{O}}/L_{\text{IR}}$ are similar, although slightly higher for the high-*z* galaxies.

Figure 2 displays the images of the H₂O emission line and the corresponding submillimeter continuum emission for the seven lensed high-*z* ULIRGs. The relatively low angular resolution (at best $\sim 1''$) of the current data limits for most of the sources a detailed study of the spatial properties of their signal. Four of the sources remain unresolved at the present angular resolution: SDP.9, SDP.17b, NA.v1.144, and G15.v2.779. The three

remaining sources are extended and display distinct lensed morphologies including two-image configuration system (SDP.81), a complex elongated structure (G12.v2.30), and an extended structure (NB.v1.78). Each source is described in detail below.

Table 4 reports the values of the H₂O velocity-integrated flux density (line flux) detected at the peak of the source in one synthesized beam, $I_{\text{H}_2\text{O}}pk$, and the velocity-integrated line flux integrated over the source size, $I_{\text{H}_2\text{O}}$. In most cases, the ratio $I_{\text{H}_2\text{O}}/I_{\text{H}_2\text{O}}pk$ is lower than 1.5, except for the extended sources

Table 4. Observed parameters of the H₂O emission lines and continuum emission.

Source	Line ^a	ν_{obs}^b [GHz]	$S_{\nu}pk^c$ [mJy/beam]	S_{ν}^c [mJy]	$S_{\nu\text{H}_2\text{O}}pk^d$ [mJy/beam]	$\Delta V_{\text{H}_2\text{O}}$ [km s ⁻¹]	$I_{\text{H}_2\text{O}}pk^d$ [Jy km s ⁻¹ /beam]	$I_{\text{H}_2\text{O}}^d$ [Jy km s ⁻¹]	$L_{\text{H}_2\text{O}}^e$ [10 ⁷ L _⊙]
SDP.81	2	244.5000	12.5	27.0 ± 1.2	10.7 ± 1.9	140 ± 50	1.6	1.8 ± 0.5	32 ± 9
NA.v1.144	1	234.8343	9.4	16.4 ± 0.7	22.8 ± 1.7	200 ± 50	4.9	7.5 ± 0.9	57 ± 7
SDP.9	1	292.1425	17.4	21.7 ± 0.9	20.3/37.8 ± 2.9 ^f	410 ± 50	11.3	14.4 ± 1.1	60 ± 5
NB.v1.78	2	240.2896	15.4	36.9 ± 0.4	5.0 ± 0.8	510 ± 90	2.7	6.7 ± 1.3	122 ± 24
G12.v2.30	2	231.9458	4.9	21.0 ± 0.8	4.5/4.5 ± 0.7 ^f	690 ± 80	2.7	6.5 ± 1.2	128 ± 24
G15.v2.779	1	143.3837	5.9	7.3 ± 0.2	3.2/6.6 ± 0.5 ^f	890 ± 140	3.4	4.1 ± 0.6	94 ± 9
SDP.17b	2	298.8646	28.6	37.9 ± 1.1	22.9 ± 2.6	300 ± 30	7.2	7.8 ± 0.9	85 ± 10

Notes. ^(a) Lines 1 and 2 correspond to the H₂O (2₁₁–2₀₂) ($\nu_{\text{rest}} = 752.03$ GHz, $E_{\text{up}} = 137$ K) and (2₀₂ – 1₁₁) ($\nu_{\text{rest}} = 987.93$ GHz, $E_{\text{up}} = 101$ K) transitions, respectively. ^(b) ν_{obs} is the frequency corresponding to the zero of the velocity scale of the spectra of Fig. 1. ^(c) $S_{\nu}pk$ corresponds to the peak continuum flux density (mJy/beam); S_{ν} refers to the total spatially integrated continuum flux density (mJy). ^(d) $S_{\nu\text{H}_2\text{O}}pk$ corresponds to the peak H₂O flux density (mJy/beam); $I_{\text{H}_2\text{O}}pk$ is the peak H₂O velocity-integrated flux density (flux in one beam, in Jy km s⁻¹/beam) and $I_{\text{H}_2\text{O}}$ the spatially integrated H₂O line flux (in Jy km s⁻¹). The values derived for SDP.81 and G12.v2.30. are underestimated (see text). ^(e) $\mu L_{\text{H}_2\text{O}}$ is the apparent luminosity of the observed H₂O line using the relations given in, e.g., Solomon et al. (1992). ^(f) In the cases of SDP.9, G12.v2.30 and G15.v2.779, the H₂O emission lines are double-peaked and the parameters in the table refer to double Gaussian fits.

G12.v2.30 and NB.v1.78. For the other sources studied in this paper, except for SDP.81, it is thus unlikely that significant flux is missed. This is also confirmed when comparing our total spatially integrated continuum flux density S_{ν} with measurements published elsewhere at the nearby frequency of 250 GHz, using the bolometer MAMBO camera at the IRAM 30-m telescope (Negrello et al. 2010; Dannerbauer et al., in prep.). NA.v1.144 and SDP.81, which were observed at the PdBI close to 250 GHz, agree very well with MAMBO observations. For SDP.9 and SDP.17b, which were observed at the PdBI close to 300 GHz, the comparison is more difficult; however, S_{ν} is somewhat higher than the values expected from extrapolations of the MAMBO values. Although there is no MAMBO observation for NB.v1.78, S_{ν} agrees fairly well with the flux density measured at 880 μm using the Smithsonian Millimeter Array (SMA; Bussmann et al. in prep.). Finally, the continuum flux density measured at 143 GHz for G15.v2.779 is in excellent agreement with the value reported by Cox et al. (2011) at 154 GHz.

The 3.6 GHz bandwidth covered by the PdBI receivers (and correlator), which was generally centered on the redshifted frequency of the H₂O emission line, encompasses other potential emission lines of interest, including i) ¹³CO, H₂¹⁸O, H₂¹⁷O, H₃O⁺ in the vicinity of H₂O(2₀₂–1₁₁); and ii) H₂O⁺(2₀₂–1₁₁, $J_{3/2,3/2}$ and $J_{5/2,3/2}$) in the vicinity of H₂O(2₁₁–2₀₂). However, the current sensitivity of our spectra is too low to detect these lines, which are not seen in the *Herschel* spectra of Mrk 231 and Arp 220 except for the two H₂O⁺ lines, which are present in the spectrum of Arp 220 (Rangwala et al. 2011). Apart from the H₂O emission lines, no other emission line is detected at the S/N of the spectra of the seven high- z galaxies studied in this paper.

3.2. Individual sources

The following describes the properties of the individual lensed galaxies.

3.2.1. SDP.81 at $z = 3.040$

The source SDP.81 is at the highest redshift among the SDP lensed galaxies of Negrello et al. (2010). It displays two arcs in the SMA 880 μm data, which are also seen in the *Hubble* Space Telescope (HST) (Negrello et al., in prep.) and the PdBI

CO imaging results (Cox et al., in prep.). Preliminary lensing modeling indicated a very high lensing magnification factor of $\mu = 19$ (Negrello et al. 2010), but this value has been revised downward to $\mu = 9.5 \pm 0.5$ based on recent HST data (Dye et al., in prep.). This is confirmed by the SMA 880 μm imaging results which yield $\mu = 9.5 \pm 0.8$ (Bussmann et al., in prep. – hereafter B13). As discussed in Sect. 4.1, we adopted the values of μ derived from the 880 μm SMA data by B13 as basic reference for the lensing magnification factors listed in Table 2.

The CO line profile is asymmetrical with a red component that is much stronger than the blue one. The line profile of the H₂O emission line is also dominated by a main red component (peaking at $\sim +200$ km s⁻¹ with a width of $\Delta v = 140$ km s⁻¹). Due to the low S/N of the H₂O spectrum, the blue emission component of the profile (peaking at -100 km s⁻¹) that is clearly seen in the CO spectrum is only barely visible in the H₂O emission spectrum (Fig. 1).

Despite the low angular resolution, the dust continuum image is clearly resolved (Fig. 2), showing both gravitational arcs of the 880 μm image of Negrello et al. (2010), with a stronger emission in the eastern arc. This structure is also seen in the CO map of Cox, Ivison et al. which indicates that the stronger red emission traces the bright eastern arc, whereas the weaker blue emission corresponds to the western arc. Due to the low S/N of the data, the H₂O emission is seen only in the eastern arc, consistent with the only clear detection of the red part of the spectrum.

The H₂O line flux $I_{\text{H}_2\text{O}}$ and the apparent $L_{\text{H}_2\text{O}}$ luminosity, which are given in Table 4, take into account the total emission including the red and blue components. However, due to the low S/N of the H₂O data, these numbers are less accurate than for most other sources in our sample. Because the emission is very extended, the spatially integrated H₂O line flux $I_{\text{H}_2\text{O}}$ is probably underestimated, as indicated by the absence of significant H₂O emission seen in the western arc, and the very small difference between $I_{\text{H}_2\text{O}}$ and $I_{\text{H}_2\text{O}}pk$ compared to the continuum ratio $S_{\nu}/S_{\nu}pk = 2.2$.

3.2.2. NA.v1.144 at $z = 2.202$

The narrow H₂O linewidth of NA.v1.144 is the same as is seen in the CO emission line ($\Delta v \sim 200$ km s⁻¹). It is the narrowest CO linewidth of our sample (Table 3).

The estimate of the lensing magnification factor from detailed modeling based on the SMA image of the 880 μm continuum emission indicates a low value of $\mu \sim 4.6 \pm 1.5$ (B13). However, the SMA imaging of this object has limited *uv* coverage (only the extended array was used) which might compromise the measurement of the lensing magnification factor. The corresponding values for the IR and CO luminosities appear to be surprisingly high compared to the narrow linewidth. Indeed, the uncertain derivation of μ from the CO linewidth and luminosity (Harris et al. 2012) yields $\mu_{\text{CO}} \sim 23 \pm 15$. Therefore, we have doubled the uncertainty on the SMA value of μ and adopted the rms²-weighted average of the estimate from the SMA dust-continuum and the CO line, i.e. $\mu = 5.3 \pm 2.9$.

This bright submillimeter galaxy is clearly unresolved with our present angular resolution, a result that is confirmed by the small extension observed in the SMA results (B13).

3.2.3. SDP.9 at $z = 1.574$

The galaxy SDP.9 is the brightest 250 μm and lowest redshift of the five lenses discussed in Negrello et al. (2010). Its lensing magnification factor is estimated to be $\mu = 8.7 \pm 0.7$ from HST imaging (Dye et al., in prep.). We adopted here the SMA 880 μm value from B13, $\mu = 8.5 \pm 1.8$, which agrees with the HST value and also with $\mu_{\text{CO}} = 11$.

The H₂O intensity of SDP.9 is by far the largest of our sample, due to the low redshift of the source. The corresponding apparent H₂O luminosity is among the highest found so far. The emission line profiles in H₂O and CO are similar, including the strong red emission spike and the weaker, broader blue emission (Fig. 1). The high angular CO data from Cox et al. (in prep.) as well as the HST/WFC3 imaging data (Negrello et al., in prep.) show that SDP.9 is a good example of an Einstein ring, with a clear asymmetry between the east and west arcs. However, our angular resolution is not high enough to resolve the emission.

3.2.4. NB.v1.78 at $z = 3.111$

This source has one of the highest redshifts and highest apparent luminosities in our sample. The intermediate linewidth of the CO and H₂O lines ($\Delta v \sim 550 \text{ km s}^{-1}$) is consistent with those measured in the mid-*J* CO lines used to initially determine the redshift with CARMA (Riechers et al., in prep.). The H₂O line is indeed the brightest line measured in this source so far and thus provides the best constraints on the line profile. The moderate CO intensity and linewidth indicates intermediate CO, H₂O and infrared intrinsic luminosities, $L_{\text{IR}} \approx 10^{13} L_{\odot}$.

The lensing magnification factor, derived from the SMA 880 μm imaging results, is $\mu = 10.5 \pm 1.4$ (B13), a value which is adopted in this paper and is consistent with $\mu_{\text{CO}} = 12$ (inferred from the CO data of Table 3).

The source is resolved and shows a relatively complex morphology in the H₂O and continuum maps. The peak value in a single beam, $I_{\text{H}_2\text{O}} pk$, reported in Table 4, is lower than the spatially integrated value $I_{\text{H}_2\text{O}}$ by a factor ~ 2.5 . The ratio is similar for the 240 GHz continuum, where the peak value is 15.4 mJy/beam and the integrated flux density is 36.9 mJy (see Table 4). In addition to the southeast main peak, there are northern and southwestern extensions (Fig. 2). This morphology is also seen in recent SMA 880 μm imaging results (Bussmann et al., in prep.). However, the southwest peak is very strong in H₂O compared to the main southeast peak and the peak ratio in

the continuum. This could suggest that the H₂O emitting region is more compact than the continuum.

3.2.5. G12.v2.30 at $z = 3.259$

G12.v2.30 is another prominent source among the H-ATLAS brightest lenses. Its complex differential lensing has been analyzed by Fu et al. (2012), showing that the lensing is very different in the near-infrared compared to the submillimeter or CO emission. The various estimates of the lensing magnification factor based on submillimeter images or the CO linewidths are comparable: $\mu = 7.6 \pm 1.5$ for the dust submillimeter and 6.9 ± 1.6 for CO(1–0) imaging by Fu et al.; 9.6 ± 0.9 for the same submillimeter data by B13 (the value adopted here); $\mu_{\text{CO}} = 7 \pm 2$ by Harris et al. (2012) from CO(1–0) linewidth and luminosity. These values confirm that G12.v2.30 is one of the most luminous infrared galaxies in H-ATLAS with an estimated $L_{\text{IR}} \sim 1.6 \times 10^{13} L_{\odot}$.

The H₂O profile appears to be double peaked with a total width of $\sim 700 \text{ km s}^{-1}$. Due to the short integration time of the observations, the CO results for G12.v2.30 are the noisiest of our sample and the details of the profile, in particular, its double-peaked nature remain unclear. Within the noise, the CO profile is compatible with the H₂O emission, although additional observations are needed to confirm this result.

The source is resolved with the present angular resolution of $2.0'' \times 1.0''$ and shows a complex structure that is most clearly seen in the continuum map (Fig. 2): an elongated arc to the south shows two peaks and emission is seen about $2''$ to the north. Due to the low S/N, only the southern component is detected in the H₂O emission line. G12.v2.30 appears to be a complex source. The dust and gas emissions as well as the stellar emission have very different morphologies, which indicates that the lensing potential is unusual (Fu et al. 2012). Estimating the total intensity of the H₂O line remains difficult for G12.v2.30. The peak value in a single beam, $I_{\text{H}_2\text{O}} pk$, reported in Table 4, is lower than the spatially integrated value $I_{\text{H}_2\text{O}}$ by a factor ~ 2.4 . The situation is similar for the 240 GHz continuum, where the peak value is 4.9 mJy/beam and the integrated flux density is 21.0 mJy (see Table 4). This latter value is itself much lower than the flux density of $\sim 36 \text{ mJy}$ measured at 1.2 mm using MAMBO at the 30-m telescope (Dannerbauer et al., in prep.). However, the reported value for the spatially integrated line flux $I_{\text{H}_2\text{O}}$ is less accurate than for most other sources in our sample. Because the emission is very extended, the spatially integrated H₂O line flux $I_{\text{H}_2\text{O}}$ is probably underestimated, similarly to the continuum, as indicated for the latter by the factor ~ 1.5 lower compared to the MAMBO value.

3.2.6. G15.v2.779 at $z = 4.244$

With a redshift higher than 4 and $F_{500 \mu\text{m}} = 204 \pm 32 \text{ mJy}$, G15.v2.779 (also referred to as ID.141) remains unique in the H-ATLAS survey (Cox et al. 2011). Its lensing magnification factor has recently been estimated to be only 4.1 ± 0.2 (Bussmann et al. 2012), which makes it the most luminous H-ATLAS source, with $L_{\text{IR}} = 2 \times 10^{13} L_{\odot}$. This is consistent with the broad CO lines, the widest of our sample ($760\text{--}890 \text{ km s}^{-1}$), which would imply $\mu_{\text{CO}} \sim 8$. The H₂O line has one of the largest S/N thanks to good observing conditions in the 2 mm window.

The width and the characteristics of the complex, double-peaked H₂O line agree well with those of the CO(7–6) and [CI] lines (Fig. 1). G15.v2.779 has the highest intrinsic

Table 5. Properties of the dust continuum and the H₂O line emission.

Source ID	line _{obs} ^a	L _{IR} (10 ¹² L _⊙)	L _{H₂O-1m} (10 ⁷ L _⊙)	L _{H₂O-1a} (10 ⁷ L _⊙)	L _{H₂O-2m} (10 ⁷ L _⊙)	L _{H₂O-2a} (10 ⁷ L _⊙)
SDP.81	2	6.1	(1.5) ^b	(2.3)	3.3	3.3
NA.v1.144	1	11	9.7	9.7	(25)	(16)
SDP.9	1	5.2	7.0	7.0	(16)	(10)
NB.v1.78	2	10	(5.1)	(7.9)	12	12
G12.v2.30	2	16	(5.8)	(9.0)	13	13
G15.v2.779	1	21	23	23	(55)	(34)
SDP.17b	2	16	(8.7)	(14)	20	20

Notes. ^(a) Line 1 is H₂O(2₁₁-2₀₂) 752.03 GHz and line 2 is H₂O(2₀₂-1₁₁) 987.93 GHz; ^(b) luminosity values for unobserved lines are written in parenthesis, they are inferred from the other observed lines assuming the line intensity ratio from Mrk 231 (González-Alfonso et al. 2010, subscript m) or Arp220 (Rangwala et al. 2011, subscript a). As detailed in the text, we consider that these values significantly underestimate the H₂O luminosity for SDP.81 and G12.v2.30.

H₂O luminosity of our sample by a factor ~2. This implies that both should have very large H₂O luminosities if the two components corresponded to two galaxies.

At the present angular resolution G15.v2.779 remains unresolved, in agreement with the compactness of the CO (Cox et al., in prep.) and the 880 μm continuum emissions (Bussmann et al. 2012).

3.2.7. SDP.17b at z = 2.305

The detection of the H₂O emission line in SDP.17b was reported and discussed in Omont et al. (2011). Its H₂O line profile is in perfect agreement with the CO emission profile – see Fig. 1 for the CO(4–3); also Harris et al. (2011) for the CO(1–0) and Iono et al. (2012) and Riechers et al. (in prep.) for the CO(3–2) profiles – including the asymmetry with the distinct blueshifted emission wing. The source is not resolved in the present data and higher angular resolution data indicate that the source remains compact and displays a velocity gradient (Cox et al., in prep.).

Again the μ value derived by B13, μ = 4.0 ± 0.6, is much lower than approximately expected from the CO linewidth and luminosity, which yields μ_{CO} = 18 ± 8 (Harris et al. 2012) – note that this latter value is uncertain because of the asymmetric CO line profile. However, the image separations are at the limit of what can be resolved by the SMA. Therefore, we have again doubled the uncertainty on the SMA value of μ and adopted the rms²-weighted average of the estimate from the SMA dust-continuum and the CO line, i.e. μ = 4.3 ± 1.2. However, the line of sight toward SDP.17b could include an intervening system, as indicated by a CO line at z = 0.942 reported by Lupu et al. (2012) based on Z-Spec data, in addition to the z = 2.3 line seen by the GBT (Harris et al. 2012) and the PdBI (Fig. 1). But this CO line at z = 0.942 was not confirmed by Omont et al. (2011) and more observations are needed to clarify this question. If there was such a dusty object at z = 0.942, the lens model derived from the SMA data of this very compact source would have to be revised.

4. Discussion

4.1. Lensing

The main purpose of this paper is to investigate the importance of the H₂O emission in high-z ULIRGs, but not to study its spatial structure by separating it from lensing images. This is the reason why the observations were carried out with the most

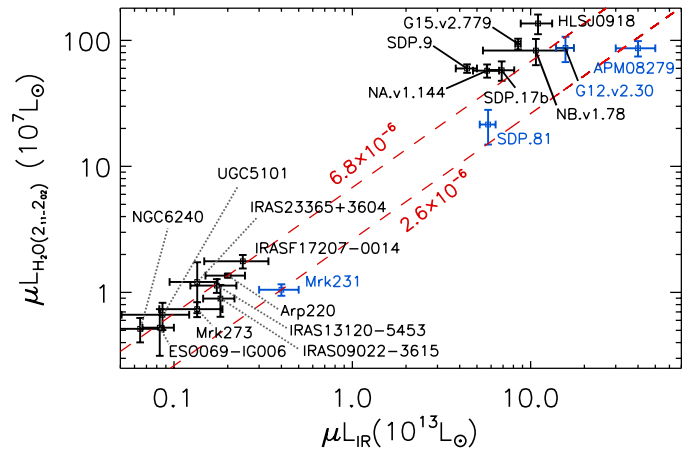


Fig. 3. Apparent H₂O luminosity versus apparent infrared luminosity of high-redshift sources, for H₂O Line 1 (2₁₁-2₀₂, 752.03 GHz) with the Arp 220 H₂O line ratio (see text), compared to local ULIRGs (μL_{IR} < 4 × 10¹² L_⊙, Yang et al., in prep.). The dashed diagonal lines have constant L_{H₂O}/L_{IR} ratios 6.8 × 10⁻⁶ and 2.6 × 10⁻⁶, equal to those of Arp 220 and Mrk 231. Because of their peculiarities, values for SDP.81 and G12.v2.30 (underestimated by factors possibly up to ~1.5–2 because they are very extended, see text), and those of APM 08279+5255 and Mrk 231 (type 1 and type 2 QSOs, respectively) are printed in blue.

compact configuration of PdBI to maximize the sensitivity for detection. Detailed lensing modeling is thus out of the scope of this paper. However, a good control of lensing effects is mandatory for various reasons: i) as seen above, the H₂O emission could be resolved out in sources that are extended, even with the limited resolution of the observations reported in this paper. ii) Differential lensing effects must be controlled for comparison with emission from other H₂O lines or other molecules or dust. As discussed, e.g., by Serjeant (2012), the importance of differential lensing is enhanced for the highest magnifications which take place close to caustics, and, especially, for blank fields, such as H-ATLAS where the deflectors are either single massive galaxies or small groups of galaxies. A striking example is the large difference observed in lensed images in the near-infrared compared to the submillimeter or CO emission of G12.v2.30 (Fu et al. 2012). iii) Lensing magnification must be corrected for deriving intrinsic properties of the sources such as luminosities or line ratios.

To infer intrinsic infrared and H₂O luminosities of our sources, a first approximation is to derive a (mean) lensing

magnification factor μ . We systematically preferred values of μ inferred from submillimeter imaging rather than near-infrared HST imaging, since H₂O emission should be more related to submillimeter dust emission than to near-infrared stellar emission, and submillimeter and near-infrared lensing can be very different as exemplified in the case of G12.v2.30 (Fu et al. 2012). As discussed in Sect. 3.2, the lensing magnification factor μ is thus determined from existing lensing models based on 880 μ m SMA imaging results (B13; see also Bussmann et al. 2012). For two sources these values were slightly corrected by also using the values of μ derived following a method described in Harris et al. (2012), which uses the relation between the CO linewidth and luminosity. As the low angular resolution and limited S/N of our H₂O data do not allow us to trace possible differential effects of lensing, we adopted a minimum uncertainty of $\pm 20\%$ for all values of μ .

4.2. Relation between H₂O and infrared luminosities

The intrinsic H₂O luminosities, based on the magnification factors listed in Table 2, are reported for each source in Table 5. The information on the H₂O emission is limited by the observation of a single emission line per source. This prevents us from comparing the H₂O excitations between different sources. It is nevertheless interesting to compare their H₂O line luminosity and to check their dependency on other properties, such as the infrared intrinsic luminosity, L_{IR} , since infrared excitation is thought to play a key role in the excitation of the H₂O levels.

In addition, the H₂O lines that were observed differ from source to source: three sources were measured in the H₂O(2₁₁–2₀₂) emission line (“Line 1”) and four sources in the H₂O(2₀₂–1₁₁) emission line (“Line 2”) To deal with this disparity, we approximated the integrated line flux, $I_{\text{H}_2\text{O}}$, of the emission line that was not observed, by adopting the value of the line ratio of “Line 2” to “Line 1” measured either in Arp 220, $r_{21} = 3440/2970 = 1.16$ (Rangwala et al. 2011) or in Mrk 231, $r_{21} = 718/415 = 1.73$ (González-Alfonso et al. 2010), listed as “a” and “m” in Table 5. The corresponding values of $L_{\text{H}_2\text{O}}$ are reported in parenthesis in Table 5. Finally, the full set of values of the intrinsic luminosities $L_{\text{H}_2\text{O}-1/2}$ for lines 1 and 2 and the line ratios of Arp 220 or Mrk 231 are reported in Table 5.

There are several arguments to favor the conversion factor of Arp 220 rather than Mrk 231: for all galaxies, except G12.v2.30, the CO excitation seems significantly lower than in Mrk 231 (Lupu et al. 2012); using the Mrk 231 conversion factor seems to yield anomalously low values for the line 1 luminosity of sources only observed in line 2 (Table 5).

Therefore, keeping all uncertainties in mind, the relationship between $L_{\text{H}_2\text{O}}$ and L_{IR} is displayed in Fig. 4 for $L_{\text{H}_2\text{O}-1a}$ (line 1 with Arp 220 line ratio). This relationship is inferred from all sources studied in this paper - with the exception of SDP.81 and G12.v2.30 whose H₂O intensities are uncertain and probably underestimated by factors possibly up to ~ 1.5 –2 because they are very extended – together with HLSJ0918 from Combes et al. (2012), and local ULIRGs ($L_{\text{IR}} < 4 \times 10^{12} L_{\odot}$) from Yang et al. (in prep.). A best-fit power-law to this relation following

$$L_{\text{H}_2\text{O}} = L_{\text{IR}}^{\alpha} \quad (1)$$

yields $\alpha = 1.15 \pm 0.10$ using “line 1” and the Arp 220 line ratio. Using the Mrk 231 line ratio would yield $\alpha = 1.11 \pm 0.10$, whereas using “line 2” would give $\alpha = 1.22 \pm 0.10$ and 1.28 ± 0.12 for the Arp 220 and Mrk 231 line ratios, respectively. Independent of the selected lines or ratios, these results

indicate that $L_{\text{H}_2\text{O}}$ increases rapidly with L_{IR} , possibly faster than linear. We note, however, that since systematic errors were not taken into account in this analysis, adding the statistical and systematic errors could result in a relation that is closer to being linear. A steep increase is consistent with a significant role of infrared radiation in the excitation of the water vapor lines. It is interesting to note that a similar behavior with the infrared luminosity has been found for the intensity of local H₂O and OH megamasers, which could be interpreted as a result of the importance of infrared pumping on these masers (e.g. Lo 2005).

However, it is important to stress that, as shown e.g. by the modeling of H₂O excitation in Mrk 231 by González-Alfonso et al., the intensity of (optically thick) H₂O lines results from a complex interplay between various parameters that includes gas temperature and density, H₂O abundance, infrared radiation field, spatial distribution of H₂O and dust, etc. The present data, with only one H₂O emission line per source, prevents us for constraining the water excitation. Multi-line H₂O observations are required for a detailed modeling of the excitation. Therefore, we defer the study of the water excitation in these high-*z* sources to a future paper, which will report measurements of multiple water transitions.

The high-*J* HCN lines seem to display a similar behavior with the infrared luminosity, with L_{HCN} growing possibly faster than linear with L_{IR} . For instance, Bussmann et al. (2008) reported that $L_{\text{HCN}(3-2)}$ varies as $L_{\text{IR}}^{1.3 \pm 0.1}$ in a sample of 30 nearby galaxies, although correcting for the difference in emission extension seems compatible with a linear relation (Gao et al., in prep.). The strong HCN(5–4) emission observed by Weiss et al. (2005) in APM 08279+5255 seems to require infrared excitation. The high HCN 6–5/5–4 ratio observed by Riechers et al. (2010) ultimately confirmed that infrared excitation is important for the HCN excitation in this source. Note that this is not true for HCN(1–0), which gives a strictly linear relation, with an index of 1.0, consistent with the interpretation that HCN(1–0) is proportional to the mass of dense gas (Solomon et al. 1992; Gao & Solomon 2004). However, because the H₂O and high-*J* HCN line intensities may be increasing faster than linear with L_{IR} may indicate the importance of IR excitation. This needs to be confirmed by better statistics and by future observations of higher transitions.

4.3. Comparison of H₂O lines with high-*J* CO lines and the continuum

In all cases where the information is available on adjacent CO lines – CO(6–5) or (7–6) for H₂O(2₁₁–2₀₂) (line 1), and CO(8–7) or (9–8) for H₂O(2₀₂–1₁₁) (line 2) – the line flux $I_{\text{H}_2\text{O}-1/2}$ is at least ~ 0.3 –0.5 that of the neighboring high-*J* CO lines. The lack of detailed data on high-*J* CO emission for most of the sources limits any detailed comparative study.

The three SDP sources and G12.v2.30 benefit from a complete spectral coverage from 200 to 300 GHz with ZSpec (Lupu et al. 2012, and in prep.). This range includes 2 to 5 high-*J* CO lines, depending on the redshift. Even though the sensitivity and the spectral resolution of ZSpec are limited, these data allow us to derive useful trends on the strength of CO as the approximate spectral line energy distributions (SLEDs) derived by Lupu et al. (2012).

In SDP.17b, the ratio $I(\text{H}_2\text{O}(2_{02}-1_{11}))/I(\text{CO}(8-7))$ is about 0.5 (Omont et al. 2011). For SDP.9, information is lacking on CO emission from levels $J > 6$ because of the low redshift ($z = 1.57$) of the source. However, the CO(6–5) emission line is strong with a ratio $I(\text{H}_2\text{O}(2_{11}-2_{02}))/I(\text{CO}(6-5)) \sim 0.36$. In the

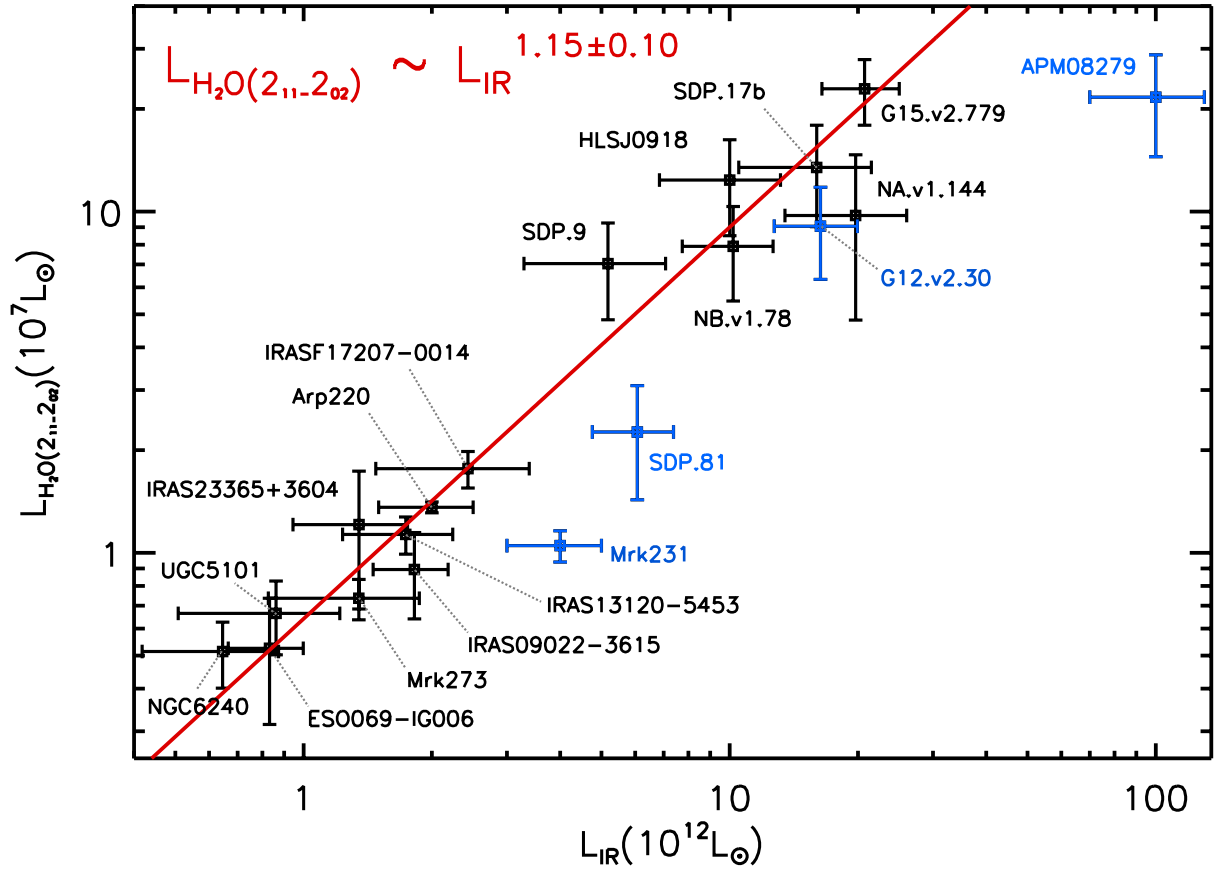


Fig. 4. Relation of the intrinsic H₂O luminosity with the infrared luminosity for the seven ultra-luminous infrared galaxies reported in this paper, other well-known high-*z* sources taken from the literature, and local ULIRGs ($L_{\text{IR}} < 4 \times 10^{12} L_{\odot}$, Yang et al., in prep.). The luminosity is reported to the H₂O ($2_{11}-2_{02}$, 752.03 GHz) emission line (referred to as “line 1” in this paper), where we adopted an Arp 220 H₂O line ratio when another H₂O line (“line 2”) was measured (see text). The solid line shows a best-fit power-law, $L_{\text{H}_2\text{O}} = L_{\text{IR}}^{\alpha}$, where $\alpha = 1.15 \pm 0.10$. Error bars are defined in the following way: i) for L_{IR} , they combine the uncertainties on the values of μL_{IR} and the uncertainties on μ given in Table 2; ii) for $L_{\text{H}_2\text{O}}$, they include the errors on $I_{\text{H}_2\text{O}}$ of Table 4, and the same uncertainties on μ except that when the latter is $< 20\%$ it is replaced by 20% (see Sect. 4.1). Because of their peculiarities, values (in blue) for SDP.81 and G12.v2.30 (both underestimated by factors possibly up to $\sim 1.5-2$ because they are very extended, see text), and those of APM 08279+5255 and Mrk 231 (type 1 and type 2 QSOs, respectively) are not taken into account for the fit.

case of SDP.81, the H₂O and CO emission lines are both weak. Accounting for the part of the H₂O emission that is not detected in the intrinsically weaker western arc (see Sect. 3.2.1), we estimate that the ratio $I(\text{H}_2\text{O}(2_{02}-1_{11}))/I(\text{CO}(7-6))$ is ≈ 0.3 – note that the CO $J > 7$ lines are not detected in the ZSpec data reported by Lupu et al. (2012). The ZSpec spectrum of G12.v2.30 displays strong lines of H₂O($3_{21}-3_{12}$) and CO($10-9$), with about similar intensities (Lupu et al. in prep.), but their ratio may be uncertain by a factor up to two because of the noise in this frequency range of ZSpec.

The best-documented comparison of H₂O and CO lines is found for G15.v2.779 because of its high redshift ($z = 4.24$) and the sensitive PdBI CO observations of Cox et al. (2011). For G15.v2.779, the ratio $I(\text{H}_2\text{O}(2_{11}-2_{02}))/I(\text{CO}(7-6)) = 0.52$. For the two other sources, NA.v1.144 and NB.v1.78, the lack of information on the high- J CO levels prevents any comparison with the H₂O emission lines.

Moreover, the intensities of H₂O lines ($I_{\text{H}_2\text{O}}$ – Table 4) and of the lower- J CO 2 mm and 3 mm lines (Table 3 and Fig. 1) are roughly comparable. Another direct comparison is the relation between the intensity of the H₂O emission line (hence the H₂O luminosity) with the underlying continuum emission (hence the infrared luminosity) – see Table 4 and Figs. 3 and 4. But the comparison between $I_{\text{H}_2\text{O}}$ and the continuum flux density S_{ν} significantly depends on the redshift and the rest frequency.

Finally, as shown by Fig. 1 and Table 4, the H₂O and continuum flux densities, e.g. $S_{\nu\text{H}_2\text{O}} \text{ pk}$ and $S_{\nu} \text{ pk}$, are roughly comparable, although their ratio obviously depends on the linewidth.

4.4. Implication of ubiquitous detection of H₂O in high-*z* ultra-luminous infrared galaxies

As in the case of SDP.17b (Omont et al. 2011), the strength of the H₂O $2_{02}-1_{11}$ and $2_{11}-2_{02}$ emission lines in six new high-*z* ultra-luminous galaxies implies that water vapor has a high abundance in these sources – probably $n(\text{H}_2\text{O})/n(\text{H}_2) \gtrsim 10^{-6}$ as in Mrk 231 (González-Alfonso et al. 2010), although this has to be confirmed by detailed modeling – and that its rotation level ladder is excited up to energies of at least ~ 100 K. This is consistent with the recent results obtained on local ULIRGs such as Arp 220 and Mrk 231 (Rangwala et al. 2011; van der Werf et al. 2010, and in prep.; Yang et al., in prep.).

Our results seem to indicate that high-*z* ultra-luminous galaxies are analogs of their local counterparts scaled by factors up to an order of magnitude in infrared luminosity and star formation rate. The similitude in excitation of high-energy H₂O levels, observed in local ULIRGs or in APM 08279+5255 (van der Werf et al. 2011), still needs to be confirmed in high-*z* sources by future observations of the higher frequency H₂O lines connected

to these levels. Nevertheless, the results reported here already indicate many similarities, such as the possibly slightly faster-than-linear increase of the intrinsic H₂O luminosity, $L_{\text{H}_2\text{O}}$, with the infrared luminosity L_{IR} ; the corresponding fact that the ratio $L_{\text{H}_2\text{O}}/L_{\text{IR}}$ in high-*z* sources is slightly higher than in local ULIRGs (Fig. 3); and the high intensity ratio of H₂O lines to nearby high-*J* CO lines, in between 0.3 to 0.5, as observed in local ULIRGs. It is therefore likely that the infrared radiation plays a significant role in the H₂O excitation, as demonstrated in the case of Mrk 231 by the detailed modeling of its spectrum by Gonzalez-Alfonso et al. (2010).

The ubiquity of H₂O emission lines in the submillimeter spectra of luminous high-*z* starburst galaxies underlines the importance of water vapor in probing their warm, dense, and dusty interstellar cores and, in addition, their strong local infrared radiation fields. Indeed, the H₂O lines constitute a totally different diagnostic from, e.g., CO or low-*J* HCN emission lines, that are excited by collisions and hence trace only temperature and density. Because of the large electric dipole of H₂O and its large rotational constant, the critical density for collisional excitation of its excited levels is extremely high, typically 10^8 cm^{-3} , much higher than expected densities in our sources. The relatively easy detection of H₂O emission lines in the sample of high-*z* starburst galaxies studied in this paper is therefore indicative of the presence of dusty, infrared-opaque nuclear disks in their centers as in local ULIRGs and of a significant role of infrared excitation of H₂O (e.g. González-Alfonso et al. 2010).

The role of a strong active galactic nucleus (AGN) in the excitation of H₂O of local ULIRGs is not yet entirely clear (van der Werf et al. 2010). There is no definite evidence of the presence of a strong AGN in any of the seven sources discussed in this paper, neither from their optical/near-infrared spectra nor from their X-ray emission. There is also no indication of a high radio or mid-infrared excess that could have been detected in the FIRST radio survey (Becker et al. 1995) or the Wide-field Infrared Survey Explorer (WISE; Wright et al. 2010) – see Omont et al. (2011) for a detailed discussion of the radio emission of SDP.17b. However, in the case of SDP.81, the radio spectral index suggests that a fraction of the radio emission is powered by an AGN (Valtchanov et al. 2011).

We stress again that most of our sources are much more powerful than local ULIRGs and that they probe a regime of infrared luminosities up to a few $10^{13} L_{\odot}$, i.e., an order of magnitude higher than local ULIRGs. It is probable that the local radiation field is also stronger, and that it might reach the limit of stability of such cores against radiation pressure, possibly close to the limit of maximum starbursts. As starbursts with intrinsic far-infrared luminosities much higher than a few $10^{13} L_{\odot}$ do not appear to exist (see e.g. Karim et al. 2013), it seems that some of the objects that we are studying, such as G15.v2.779, might be close to this critical regime.

All sources studied in this paper are highly magnified by gravitational lensing. This makes the detection of lines such as H₂O much easier, especially for the highest lensing magnification factors, i.e., the lowest intrinsic luminosities. However, for the strongest luminosities, a few $10^{13} L_{\odot}$, lensing magnification factors go down to a few units; e.g. $\mu \sim 4$ for G15.v2.779. Therefore, for similarly luminous sources, H₂O lines should be easily detectable even in the absence of lensing magnification. For instance, a source identical to G15.v2.779 ($L_{\text{IR}} \sim 2 \times 10^{13} L_{\odot}$ at $z > 4$) would have $I_{\text{H}_2\text{O}} \sim 1 \text{ Jy km s}^{-1}$ without any magnification, which is well within the reach of the PdBI in its current configuration.

5. Conclusion

This paper reports a significant step forward in the study of water vapor emission in the warm and dense cores of ultra-luminous starburst galaxies at high redshift. The seven sources studied here are strongly lensed galaxies discovered in the *Herschel* H-ATLAS survey and, using the PdBI, all were detected in H₂O either in the (2₀₂–1₁₁) or the (2₁₁–2₀₂) rotational line. One of the sources (SDP.17b) was previously reported in Omont et al. (2011). In all cases, the H₂O emission lines are strong, with integrated line fluxes ranging from 1.8 to 14 Jy km s⁻¹, comparable both in strength and in the characteristics of their profiles to the CO emission lines. The derived apparent infrared luminosities, $\mu L_{\text{H}_2\text{O}}$, range from ~ 3 to $12 \times 10^8 L_{\odot}$. When correcting for the magnification factors μ , which are estimated from lensing models based on submillimeter imaging and corroborated by the empirical relationship between CO luminosity and linewidth, the derived H₂O luminosities show a strong dependence on the infrared luminosity. It is found that $L_{\text{H}_2\text{O}}$ varies as $\sim L_{\text{IR}}^{1.2}$, the exact shape of the relation remaining uncertain mainly due to the uncertainties in the derivation of μ . This slightly non-linear relation may indicate that infrared pumping plays a role in the excitation of H₂O lines. Water emission is therefore expected to become very strong in the most ultra-luminous ($L_{\text{IR}} > 10^{13} L_{\odot}$) galaxies at high redshift.

The results described in this paper underline the fact that the most powerful high-*z* infrared galaxies are analogs of local ULIRGs scaled by factors up to 10 in infrared luminosity and star formation rate. The fact that the excitation of H₂O levels with $E_{\text{up}} \sim 100\text{--}150 \text{ K}$ in high-redshift luminous galaxies appears to be similar to what is seen in local ULIRGs is supported by the present results, including the rapid increase of $L_{\text{H}_2\text{O}}$ with L_{IR} , the slightly higher ratio $L_{\text{H}_2\text{O}}/L_{\text{IR}}$ in high-*z* sources than in local ULIRGs, or the high intensity ratio of H₂O lines to nearby CO lines. However, future observations of the higher frequency H₂O lines connected to upper levels in high-*z* galaxies are needed to compare the excitation of the high-energy H₂O levels with $E_{\text{up}} > 150\text{--}200 \text{ K}$ to that of local ULIRGs.

The detection of H₂O emission lines in seven high-redshift luminous galaxies reported in this paper demonstrates that the millimeter/submillimeter lines of water vapor are a new and key diagnostic tool to probe both the extreme physical conditions and the infrared radiation in the warm dense environments of luminous high-*z* starburst galaxies. Follow-up observations of such sources, taking advantage of the increased sensitivities that are and will become available with ALMA and the upgraded PdBI (NOEMA) facilities, will allow more comprehensive studies of water vapor in high-*z* luminous galaxies, including unlensed candidates, providing essential and new insights into the physical conditions of these galaxies, which were present when the universe was young and have no equivalent in the local universe.

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