

Far-infrared colours of nearby late-type galaxies in the *Herschel* Reference Survey^{★,★★}

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Received 7 December 2011 / Accepted 10 January 2012

ABSTRACT

We study the far infrared (60–500 μm) colours of late-type galaxies in the *Herschel* Reference Survey, a *K*-band selected, volume limited sample of nearby galaxies. The far infrared colours are correlated with each other, with tighter correlations for the indices that are closer in wavelength. We also compare the different colour indices to various tracers of the physical properties of the target galaxies, such as the surface brightness of the ionising and non-ionising stellar radiation, the dust attenuation and the metallicity. The emission properties of the cold dust dominating the far infrared spectral domain are regulated by the properties of the interstellar radiation field. Consistent with that observed in nearby, resolved galaxies, our analysis shows that the ionising and the non-ionising stellar radiation, including that emitted by the most evolved, cold stars, both contribute to the heating of the cold dust component. This work also shows that metallicity is another key parameter characterising the cold dust emission of normal, late-type galaxies. A single modified black body with a grain emissivity index $\beta = 1.5$ better fits the observed SPIRE flux density ratios S_{250}/S_{350} vs. S_{350}/S_{500} than $\beta = 2$, although values of $\beta \approx 2$ are possible in metal rich, high surface brightness galaxies. Values of $\beta \leq 1.5$ better represent metal poor, low surface brightness objects. This observational evidence provides strong constraints for dust emission models of normal, late type galaxies.

Key words. infrared: galaxies – galaxies: spiral – galaxies: ISM

1. Introduction

The emission of late-type galaxies in the infrared domain ($5 \leq \lambda \leq 1000 \mu\text{m}$) is dominated by the dust component of the interstellar medium heated mainly by the stellar radiation field. Young stars are the dominant heating sources in star forming

* *Herschel* is an ESA space observatory with science instruments provided by European-led Principal Investigator consortia and with important participation from NASA.

** Appendices are available in electronic form at <http://www.aanda.org>

objects, while evolved stars can have a major contribution in quiescent systems. Although not dominant in mass, dust plays a major role in the equilibrium of the interstellar medium (Sauvage & Thuan 1992). Formed by the aggregation of metals produced by massive and asymptotic giant branch (AGB) stars at the latest stages of their evolution (Valiante et al. 2009) and injected into the interstellar medium through stellar winds and supernovae explosions, dust grains are important coolers of the gaseous phase (Bakes & Tielens 1994). They also act as catalysts in the process responsible for the formation of the molecular hydrogen within molecular clouds, and are thus of paramount importance in the process of star formation (Hollenbach & Salpeter 1971). Heated by the stellar radiation produced mainly by newly formed stars, the dust emission is often used as a powerful tracer of the star formation activity of galaxies, in particular in highly extinguished objects such as dusty starbursts and ultraluminous infrared galaxies (e.g. Kennicutt 1998).

Dust is composed of a mixture of carbonaceous and amorphous silicate grains of different size and composition (see for a review Draine 2003). Big grains, of size $a > 200 \text{ \AA}$, composed of carbon (graphite or amorphous) and silicates, are in thermal equilibrium with the UV and the optical radiation field and dominate the IR emission for $\lambda > 60 \mu\text{m}$. While the energy output of galaxies at wavelengths $\lambda \lesssim 200 \mu\text{m}$ is dominated by the warm dust, at longer wavelengths ($\lambda \gtrsim 200 \mu\text{m}$) is due to the cold dust emission dominating in mass (e.g. Whittet 1992).

The IRAS (Neugebauer et al. 1984), ISO (Kessler et al. 1996), *Spitzer* (Werner et al. 2004) and more recently AKARI (Murakami et al. 2007) space missions allowed us to gather data for hundreds of thousands of extragalactic sources in the spectral domain $\lambda \leq 170 \mu\text{m}$. Sub-millimetric ground based facilities such as SCUBA on the JCMT telescope allowed the observation of the coldest dust at $\lambda \approx 850 \mu\text{m}$. The *Herschel* Space Observatory (Pilbratt et al. 2010), launched in May 2009, has been designed to extend previous imaging and spectroscopic infrared observations in the spectral range 55–672 μm . Within the SPIRE and PACS guaranteed time consortia, we have started several observing programs using the three *Herschel* instruments (SPIRE, Griffin et al. 2010; PACS, Poglitsch et al. 2010; and HIFI, de Graauw et al. 2010) with the aim of studying the dust emission properties of different types of galaxies. The targets of these programmes have been selected to sample the largest parameter space in morphological type, luminosity, stellar and nuclear activity, metallicity and environment. One of these projects, the *Herschel* Reference Survey (HRS), has been designed to cover a complete, volume limited, statistically significant sample of nearby galaxies spanning the whole Hubble sequence of morphological type. The aim of this project is to derive the mean statistical properties such as the far infrared luminosity distribution, the scaling relations, the colours, and the spectral energy distributions of a representative sample of galaxies in the nearby universe (Boselli et al. 2010a). Its completeness makes the HRS the ideal sample for such a purpose.

The present paper studies the broad band far infrared colours of late-type systems from the HRS and their relationships with different tracers of the physical properties of the observed galaxies. SPIRE data obtained at 250, 350 and 500 μm are combined with IRAS 60 and 100 μm data and with other multifrequency data to trace the empirical properties of the dust emission of galaxies in this spectral range. The main aim of this work is that of identifying, through a multifrequency statistical analysis, the physical parameters driving the dust emission properties in galaxies. This approach is useful in the study of the interstellar medium for understanding the intrinsic properties of the emitting

dust in different types of galaxies. This analysis is also crucial to understand whether the standard assumptions taken in the interpretation of cosmological data are justified. The main result of this work is that the dust emission properties of galaxies are not universal but strongly depend on the metallicity and the intensity and the hardness of the interstellar radiation field. Understanding and quantifying these effects in the local universe is thus important for high redshift studies given the strong evolution of the physical properties of galaxies with cosmic time. Furthermore, observed far infrared colour indices can be easily compared to those of high redshift sources or of templates generally used to characterise the spectral energy distribution of galaxies at any redshift and thus to directly check their validity (Boselli et al. 2010b).

We limit the analysis to late-type systems, while those of early-types (ellipticals, lenticulars) will be addressed in another communication (Smith et al. 2012). This choice is dictated by the fact that both the dust properties and the nature of the heating sources in early-type systems are generally significantly different from those of spiral galaxies.

The present work extends the preliminary study of Boselli et al. (2010b) to the complete HRS sample. At the same time, it widens the analysis to a much larger and complete set of physical parameters necessary to characterise the physical properties of the target galaxies. Although a complete understanding of the far infrared emission of galaxies requires a comparison of observations with models, this empirical approach has the advantage of being free from any model dependent assumption and is thus particularly useful for identifying the main parameters governing the dust emission of galaxies. A detailed comparison of the data with models will be presented in upcoming communications.

2. The sample

The *Herschel* Reference Survey (HRS; Boselli et al. 2010a) provides us with an ideal sample for studying the relationship between the infrared colours and the physical properties of nearby galaxies spanning a large range in morphological type and luminosity. The HRS is a *K*-band selected, volume limited ($15 < D < 25 \text{ Mpc}$), complete sample of galaxies at high galactic latitude ($b > +55^\circ$; $A_B < 0.2$, to avoid cirrus contamination) composed of 322 objects¹. Distances (D) have been determined assuming galaxies in Hubble flow with $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$, outside the Virgo cluster. Within Virgo, galaxies are taken at fixed distance according to the subgroup membership as indicated in Gavazzi et al. (1999).

The present analysis is focused on the far infrared ($60 \mu\text{m} \leq \lambda \leq 500 \mu\text{m}$) properties of the late-type objects included in the sample. Several HRS galaxies belong to the Virgo cluster. To avoid any possible second order effect related to the perturbations induced by the cluster environment, we restrict our analysis to those objects with a normal HI gas content (149 objects). There are indeed indications that the cold dust component of the interstellar medium of galaxies is removed during their interactions with the hostile cluster environment (Boselli & Gavazzi 2006; Cortese et al. 2010a,b, 2012). We assume as normal, unperturbed objects those with an HI-deficiency parameter $HI - def \leq 0.4$, where $HI - def$ is defined as the difference in

¹ With respect to the original sample given in Boselli et al. (2010a), we removed the galaxy HRS 228 whose new redshift indicates it as a background object. We also revised the morphological type for three galaxies that moved from the early- to late-type class: NGC 5701, now classified as Sa, and NGC 4438 and NGC 4457 now Sb.

logarithmic scale between the expected and observed HI mass of a galaxy of given angular size and morphological type (Haynes & Giovanelli 1984). HI-deficiencies for all the target galaxies have been measured using the recent calibrations of Boselli & Gavazzi (2009).

The statistical properties of galaxies belonging to the HRS are extensively described in Boselli et al. (2010a). To summarise, the objects analysed in this work span a large range in morphological type (from Sa to Sm-Im-BCD), stellar mass ($8.5 \lesssim \log M_{\text{star}} \lesssim 11 M_{\odot}$), HI mass ($7.5 \lesssim \log M_{\text{HI}} \lesssim 10.5 M_{\odot}$), infrared ($7.2 \lesssim \log L_{60 \mu\text{m}} \lesssim 10.2 L_{\odot}$) and radio luminosity ($20 \lesssim \log L_{\text{Radio}} \lesssim 22.5 \text{ W Hz}^{-1}$) and star formation rate ($0.1 \lesssim SFR \lesssim 10 M_{\odot} \text{ yr}^{-1}$). Their metallicity is in the range $8.3 \lesssim 12 + \log(\text{O}/\text{H}) \lesssim 8.8$, where the mean metallicity of the Milky Way in the solar neighborhood is 8.67 (Rudolph et al. 2006), and that of the LMC and SMC is 8.40 and 8.23, respectively (Korn et al. 2002; Rolleston et al. 2003). The statistical properties of the dust component of the sample galaxies are described in a dedicated publication (Cortese et al. 2012).

3. The data

3.1. *Herschel*/SPIRE data

The HRS galaxies have been observed with the SPIRE instrument (Griffin et al. 2010) in the three bands at 250, 350 and 500 μm as part of a SPIRE guaranteed time key project the *Herschel* Reference Survey (Boselli et al. 2010a). Eighty three of the HRS galaxies lie in the footprint of the *Herschel* Virgo Cluster Survey (HeViCS; Davies et al. 2010), an open time key project aimed at covering with PACS and SPIRE ~ 60 sq deg of the Virgo cluster region. We took the SPIRE data of these 83 targets from HeViCS.

239 HRS objects plus 4 galaxies in HeViCS targeted by the HRS during the *Herschel* Science Demonstration Phase, out of 322 galaxies of the HRS, have been observed using the SPIRE scan-map mode with a nominal scan speed of 30''/s. Late-type galaxies analysed in this work have been observed with three pairs of cross-linked scan maps to reach a pixel-by-pixel rms of $\sim 7, 8, 8$ mJy/beam. To secure the detection of the dust associated with the extended HI disc, all galaxies have been mapped at least up to 1.5 times their optical diameter. Galaxies with optical diameters smaller than ≈ 180 arcsec were observed using the small scan-map mode providing homogeneous coverage on a circular area of $\sim 5'$. Larger galaxies have been observed in scan map mode, with typical map sizes are $8' \times 8'$, $12' \times 12'$ and $16' \times 16'$. The remaining 83 galaxies² located within the HeViCS footprint have been observed using the PACS/SPIRE parallel scan-map mode with a scan speed of 60''/s in four 16 sq deg different fields.

Both the HRS and HeViCS data have been reduced using the map making pipelines developed within the Extragalactic Science Working Group (SAG2), extensively described in Bendo et al. (2011) and Davies et al. (2011). Data have been processed up to Level-1 using the standard SPIRE pipeline. We use the BriGAdE method (Smith et al., in prep.) to remove the temperature drift and bring all bolometers to the same level to secure the best baseline removal when temperature variations are present. The different scans were then combined using the standard SPIRE making pipeline. The resulting images have pixel

² Only 23/83 HeViCS galaxies match the conditions on the morphological type and the HI-deficiency parameter mentioned in the previous section. They are thus only 15% (23/149; see Table 1) of the objects analysed in this work.

Table 1. Completeness of the multifrequency data.

Data	N. obj.	Note
Complete sample	149	1
SPIRE	146	2
IRAS	128	3
SFR	138	4
<i>b</i>	138	4
$\Sigma(\text{H}\alpha)$	110	4
$\mu_c(\text{H})$	149	
$12 + \log(\text{O}/\text{H})$	126	5
$A(\text{H}\alpha)$	126	6
$A(\text{FUV})$	114	7

Notes. 1: Late-type galaxies with an HI-deficiency ≤ 0.4 . 2: Detected in the three SPIRE bands. 3: Detected by IRAS at 60 and 100 μm . 4: Determined from $\text{H}\alpha$ data accurately corrected for [NII] contamination and dust extinction as described in Boselli et al. (2009). Galaxies hosting an AGN are not included. 5: With a metallicity measured using the same calibration (Hughes et al., in prep.). 6: With a direct measure of the Balmer decrement from integrated spectra (Boselli et al., in prep.). 7: With $A(\text{FUV})$ directly measured from the far infrared to FUV flux ratio as described in Cortese et al. (2008).

sizes of 6, 8 and 12 arcsec with FWHM of 18.2'', 24.5'' and 36.0'' at 250, 350 and 500 μm , respectively.

Flux densities in the three SPIRE bands have been extracted using concentric elliptical apertures adapted to match the galaxy shape on the plane of the sky and to avoid unwanted contaminating sources (background objects, nearby companions, Galactic cirrus) in the annulus selected for the sky determination, as extensively described in Ciesla et al. (in prep.). The maximal elliptical aperture is taken, whenever possible, at 1.4 times the optical radius of the galaxy (measured at the 25 B mag arcsec⁻²), while the background in circular annuli at larger distances. For point like sources (2 objects matching the selection criteria), flux densities are measured directly from the time line data, following the prescription proposed by Bendo and presented in the SPIRE manual³. Uncertainties are measured following the prescription of Boselli et al. (2003). Typical uncertainties on the SPIRE flux densities of late-type galaxies due to the flux extraction procedure are of the order of 5.9, 7.6 and 10.5% at 250, 350 and 500 μm , respectively (Ciesla et al., in prep.). Fifteen galaxies have independent data since they have been observed during both the HRS and the HeViCS survey. The median difference in their flux densities is 2, 2 and 3% at 250, 350 and 500 μm , respectively. The uncertainty on the absolute calibration is 7%. This uncertainty should be added quadratically to the uncertainty on the extracted flux density of each single object.

The SPIRE pipeline transforms measured fluxes into monochromatic flux densities assuming that the dust emissivity of the emitting source changes as ν^{-1} . It also assumes that the sources are point-like. Ciesla et al. (in prep.) transformed monochromatic flux densities from point-like to extended sources following the prescription given in the SPIRE manual and described in Bendo et al. (2011). These monochromatic flux densities should be also corrected to take into account that the spectral energy distribution of the observed galaxies in this spectral domain is better represented by a modified black body of emissivity $S(\nu, T) = B(\nu, T)\nu^{\beta}$ with $\beta \sim 1.5-2$ and

³ http://Herschel.esac.esa.int/twiki/pub/Public/SpireCalibrationWeb/SPIREPhotometryCookbook_jul2011_2.pdf

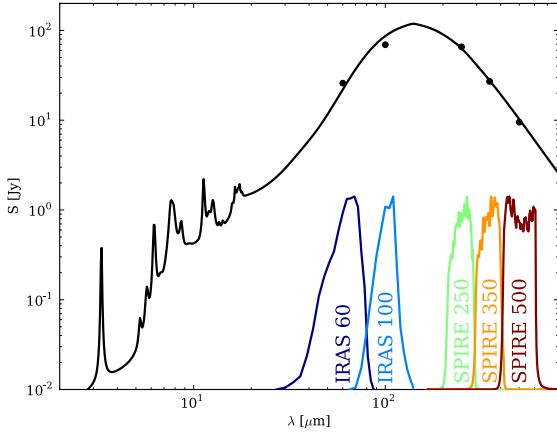


Fig. 1. The infrared (2–800 μm) spectral energy distribution of the spiral galaxy M100 (NGC 4321). The solid line is, among the Draine & Li (2007) dust models, the one better fitting the observed data (filled dots). The IRAS 60 and 100 μm bands sample the dust emitting at wavelengths shorter than the wavelength of the peak ($\sim 150 \mu\text{m}$) and are thus sensitive to the emission of the warm dust. The SPIRE bands at 250, 350 and 500 μm cover the Rayleigh-Jeans tail of the dust emission (cold dust).

$T = 15\text{--}30 \text{ K}$ for $\beta = 2$ and $T = 20\text{--}50 \text{ K}$ for $\beta = 1.5$ (see below). For these ranges of β and T the corrections for extended sources such as those analysed in this work are significantly smaller ($\leq 2\%$) than the uncertainty on the absolute calibration (7%) or on the measure of the flux density (6–11%), as indicated in Ciesla et al. (in prep.). Even considering that a fraction of the observed flux of late-type galaxies comes from unresolved sources within the galaxies (nucleus, HII regions...), for which larger corrections should be used, the mean colour corrections that should be applied to the total flux density of extended galaxies is $\leq 6.5\%$ (excluding absolute calibration), thus smaller than the mean uncertainty on the data. We thus decided not to apply any colour correction to the data. Only detected sources are considered in the following analysis (146 objects).

3.2. Ancillary data and the derived physical parameters

The SPIRE data have been combined with the IRAS 60 and 100 μm flux densities available for $\sim 90\%$ of the late-type galaxies of the sample taken from a large variety of sources (Sanders et al. 2003; Moshir et al. 1990; Thuan & Sauvage 1992; Soifer et al. 1989; Young et al. 1996) and collected on NED. These are integrated values thus directly comparable to the SPIRE flux densities used in this work. The typical uncertainty of the IRAS flux densities is of the order of 15%. We do not apply any colour correction to the IRAS data.

Different sets of multifrequency data are required to characterise the physical properties of the interstellar medium and their relationships with the dust emission properties in the far infrared spectral domain. The intensity of the general interstellar radiation field, i.e. the intensity of the radiation emitted by the bulk of the stellar component, is quantified by means of the H -band effective surface brightness $\mu_e(H)$ (in mag arcsec^{-2}), defined as the mean surface brightness within the effective radius (radius including half of the total stellar light; see Gavazzi et al. 2000). H -band magnitudes are available for all the target galaxies from 2MASS (Jarrett et al. 2003) or from our own observations (GOLDMine; Gavazzi et al. 2003).

UV GALEX data, necessary to quantify the UV attenuation and the activity of star formation, have been taken during the GALEX Virgo cluster survey (GUViCS, Boselli et al. 2011) and a dedicated Cycle 6 GALEX legacy program (Cortese et al., in prep.). $H\alpha + [\text{NII}]$ narrow band imaging data, recently obtained during a narrow band imaging survey (Boselli et al., in prep.), are used to quantify the intensity of the ionising radiation ($\Sigma(H\alpha)$) and the ongoing activity of star formation (SFR). FUV GALEX and $H\alpha + [\text{NII}]$ narrow band imaging data are corrected for dust extinction (and $[\text{NII}]$ contamination) as described in Boselli et al. (2009). Briefly, the $H\alpha + [\text{NII}]$ narrow band imaging data are corrected for $[\text{NII}]$ contamination and dust extinction using integrated spectroscopy data ($R \sim 1000$) obtained for 126/146 of the sample galaxies at the OHP 1.93 m telescope (Gavazzi et al. 2004; Boselli et al., in prep.). Standard recipes are used for those galaxies without spectroscopic data.

The Balmer decrement, here expressed as $A(H\alpha)$, is an independent, direct measure of the opacity of the interstellar medium. $A(\text{FUV})$, the attenuation in the FUV GALEX band at 1539 \AA , has been determined using the far infrared to FUV flux ratio following the prescription of Cortese et al. (2008). These recipes are indicated for normal, star forming galaxies such as those belonging to the HRS, where the general interstellar radiation field, and not only the radiation emitted by the youngest, most massive stars, contributes to the heating of the dust.

Extinction corrected $H\alpha$ fluxes and FUV flux densities are converted into star formation rates SFR using the standard calibration of Kennicutt (1998). We assume an escape fraction of zero and a fraction of ionising photons absorbed by dust before ionising the gas of zero ($f = 1$). Although unphysical (see Boselli et al. 2009), this choice has been done to allow a direct comparison with the results obtained in the literature using other star formation rates determined using $H\alpha$ data. These works generally assume $f = 1$. Our most recent results indicate that $f \sim 0.6$ (Boselli et al. 2009).

$H\alpha$ fluxes are also used to quantify the intensity of the UV ionising radiation using the $H\alpha$ surface brightness $\Sigma(H\alpha)$, where the $H\alpha$ emission is supposed to be as extended as the optical disc. This assumption is reasonable for unperturbed galaxies such as those analysed in this work (Boselli & Gavazzi 2006).

Near infrared imaging data combined with UV GALEX and $H\alpha + [\text{NII}]$ imaging data, are also used to quantify different direct tracers of the star formation history of the galaxies. This is done by measuring the birthrate parameter b (Kennicutt et al. 1994), that in a closed box model can be defined as in Boselli et al. (2001):

$$b = \frac{SFR}{\langle SFR \rangle} = \frac{SFR t_0 (1 - R)}{M_{\text{star}}} \quad (1)$$

with t_0 the age of the galaxy (13 Gyr) and R the returned gas fraction, here assumed to be $R = 0.3$ (Boselli et al. 2001). The total stellar mass M_{star} is estimated using the H -band data and recently determined colour-dependent recipes (Boselli et al. 2009). As defined,

$$b \propto \frac{L(H\alpha)}{L(H)} \quad (2)$$

measures the ratio of the ionising (photons with $\lambda < 912 \text{ \AA}$) to non ionising ($\lambda = 1.65 \mu\text{m}$) radiation and is thus a direct tracer of the hardness of the interstellar radiation field. Galaxies with a b parameter > 1 have a present day star formation activity more important than their mean star formation activity since their birth.

Table 2. Spearman correlation coefficients of the colour–colour relations (Fig. 2).

Y-colour	$S\ 60/S\ 100$	$S\ 60/S\ 250$	$S\ 100/S\ 250$	$S\ 250/S\ 350$	$S\ 250/S\ 500$
$S\ 60/S\ 250$	0.86				
$S\ 100/S\ 250$	0.65	0.94			
$S\ 250/S\ 350$	0.30	0.55	0.65		
$S\ 250/S\ 500$	0.28	0.52	0.63	0.98	
$S\ 100/S\ 500$	0.55	0.85	0.94	0.86	0.85

Notes. Only high quality detections are considered (see Ciesla et al., in prep.).

They are characterised by very blue colours and thus have hard interstellar radiation fields.

The birthrate parameter is proportional to the specific star formation rate $SSFR$ defined as (Brinchmann et al. 2004):

$$SSFR = \frac{SFR}{M_{\text{star}}} = \frac{b}{t_0(1-R)}. \quad (3)$$

This parameter is also important since it is often used to discriminate the far infrared properties of galaxies in blind infrared cosmological surveys such as H-GOODS and H-ATLAS (e.g. Elbaz et al. 2011; Smith et al. 2011). We thus use either of the two parameters in the following analysis.

Integrated spectroscopy is also used to measure gas metallicities $12 + \log(\text{O}/\text{H})$ (Hughes et al., in prep.) and Balmer decrement ($A(\text{H}\alpha)$; Boselli et al., in prep.) for most of the late-type HRS galaxies. Depending on the availability of several main emission lines, different calibrations have been adopted to convert line emissions into $12 + \log(\text{O}/\text{H})$. Following Kewley & Ellison (2008) we adopt the PP04 O3N2 calibration on the [NII] and [OIII] emission lines (Pettini & Pagel 2004) as the base metallicity. We then determine the average oxygen abundance $12 + \log(\text{O}/\text{H})$ for each galaxy.

HI gas data, available for almost the totality of the late-type galaxies of the sample, are used only to reject those objects that might have suffered any kind of perturbation induced by the Virgo cluster environment. Table 1 summarizes the completeness of the multifrequency data used for the present analysis.

4. Far infrared colours

4.1. Far infrared colour indices

Far infrared colours, here defined as the ratio of the flux densities measured in different far infrared bands, are useful quantitative tools to characterise the properties of the interstellar dust in galaxies. Those analysed in this work include:

- the widely used $S\ 60/S\ 100$ IRAS colour, often adopted as a direct tracer of the starburst activity of the target galaxies. This index is sensitive to the emission of the warm dust component principally heated by young stars (see Fig. 1). Objects with $S\ 60/S\ 100 \geq 0.5$ are generally considered as starbursts (Rowan-Robinson & Crawford 1989). In the following we will refer to this colour index as to the *warm dust sensitive index*;
- the colour indices $S\ 60/S\ 250$, $S\ 100/S\ 250$ and $S\ 100/S\ 500$ are sensitive to the relative weight of the warm and cold dust component since the peak of the emission of normal, star forming galaxies such as those analysed in this work lies in between 100 and 200 μm , thus in between the two photometric bands used to define these indices (see Fig. 1). They are related to the wavelength position of the peak of the dust

emission. We generally refer to these indices as the *dust peak sensitive indices*;

- the SPIRE colour indices $S\ 250/S\ 350$, $S\ 250/S\ 500$ and $S\ 350/S\ 500$ are indices tracing the emitting properties of the coldest dust component (*cold dust sensitive indices*). Indeed they sample the Rayleigh-Jeans tail of the dust emission dominating in mass in normal galaxies (see Fig. 1).

4.2. Colour–colour diagrams

Figure 2 shows the relationships between several far infrared colour indices for galaxies coded according to their morphological type. Table 2 lists the Spearman correlation coefficients (r) of the relations. Clearly, all colour indices are mutually related. The tightest correlations are present when the two plotted colour indices cover the closest spectral bands. The correlations between the colour indices sensitive to the emission of the coldest dust component ($S\ 250/S\ 350$, $S\ 250/S\ 500$) and the warm dust sensitive index $S\ 60/S\ 100$, however, are very marginal ($r \sim 0.3$). We do not observe any strong systematic difference between galaxies of different morphological type. The large dispersion in many colour–colour relations clearly indicates that the dust emission properties change significantly from galaxy to galaxy. We thus try to investigate which physical parameter is at the origin of these variations.

4.3. Colours vs. physical parameters

The dust of the interstellar medium is heated by the interstellar radiation field. The energy absorbed by the dust grains is re-radiated in the infrared spectral domain. It is thus expected that the dust emission properties are tightly related to the properties (intensity, hardness) of the interstellar radiation field. They can also depend on the chemical composition of the emitting dust grains, which in turn can be related to the mean metallicity of the interstellar medium. To quantify these effects, we plot in Fig. 3 the relationships between the far infrared colour indices and different tracers of the physical properties of the interstellar medium. The Spearman correlation coefficients of these relations are listed in Table 3. These tracers are:

- the star formation rate SFR , in solar masses per year;
- the birthrate parameter b that, as defined, is a tracer of the hardness of the incident interstellar radiation (hard for $b \geq 1$);
- the $\text{H}\alpha$ surface brightness $\Sigma(\text{H}\alpha)$, in $\text{erg s}^{-1} \text{kpc}^{-2}$, which measures the intensity of the ionising radiation. Corrected for dust extinction, this entity quantifies the surface density of the present day activity of star formation;
- the H -band effective surface brightness $\mu_e(H)$, in AB mag arcsec $^{-2}$, which traces the intensity of the interstellar radiation field produced by the bulk of the stellar population;

Table 3. Spearman correlation coefficients of the relations between the colour indices and the physical parameters (Fig. 3).

Y-colour Units	$\log SFR$ $M_{\odot} \text{ yr}^{-1}$	$\log b$	$\log \Sigma(H\alpha)$ $\text{erg s}^{-1} \text{ kpc}^{-2}$	$\mu_e(H)$ AB mag arcsec $^{-2}$	$12 + \log(O/H)$	$A(H\alpha)$ mag	$A(FUV)$ mag
$S60/S100$	-0.03	0.48	0.24	-0.09	-0.18	-0.20	0.11
$S60/S250$	0.01	0.44	0.44	-0.24	-0.05	-0.14	0.29
$S100/S250$	0.01	0.36	0.53	-0.31	0.01	-0.08	0.39
$S100/S500$	0.14	0.21	0.65	-0.46	0.29	0.08	0.55
$S250/S350$	0.35	-0.09	0.70	-0.62	0.56	0.28	0.70
$S250/S500$	0.29	-0.10	0.70	-0.61	0.55	0.26	0.68

Notes. Only high quality detections are considered (see Ciesla et al., in prep.).

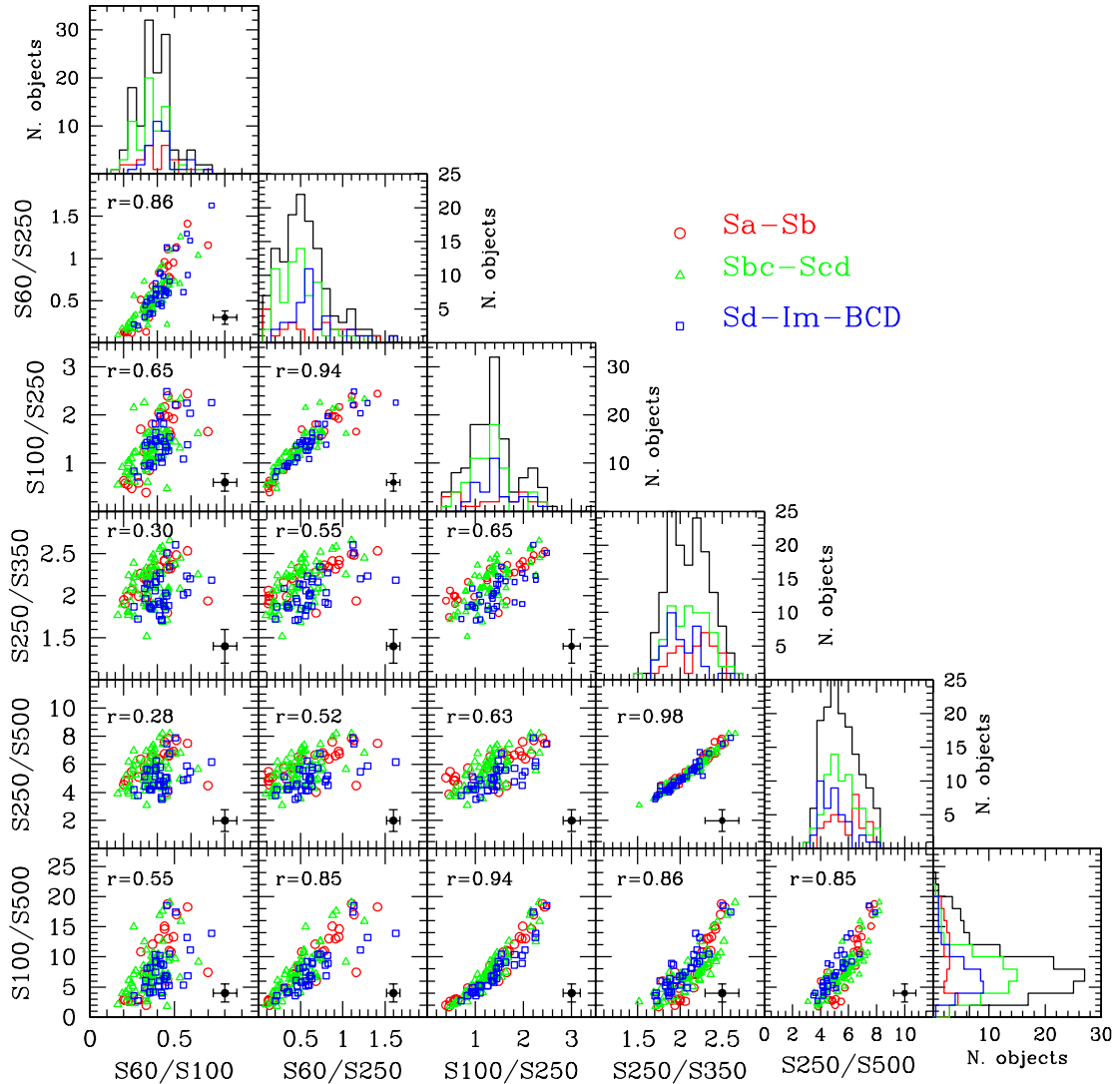


Fig. 2. Far infrared colour-colour diagrams, where colours are defined as the ratio of the flux densities measured in two different bands. Galaxies are coded according to their morphological type, red open circles for Sa-Sb (34 objects), green empty triangles for Sbc-Scd (72) and blue open squares for Sd-Im-BCD (40). The typical error bar is indicated with a black cross.

- e) the mean metallicity of the gaseous component, inferred from the $12 + \log(O/H)$ metallicity index;
- f) the Balmer decrement $A(H\alpha)$, in magnitudes, which measures the attenuation of the ionising radiation within the gaseous component of the interstellar medium;
- g) the attenuation of the non ionising UV radiation $A(FUV)$, in magnitudes, determined from the far infrared to UV flux ratio.

Figure 3 shows that only a few far infrared colour indices analysed here are correlated with some of the physical entities used to characterise the properties of the interstellar medium. The analysis of Fig. 3 brings to the following conclusions:

- 1) surprisingly, the warm dust sensitive index $S60/S100$ is not related to the direct tracers of the star formation activity (SFR , $\Sigma(H\alpha)$), while it is only (and marginally; $r = 0.48$) correlated to the birthrate parameter b . Galaxies still active

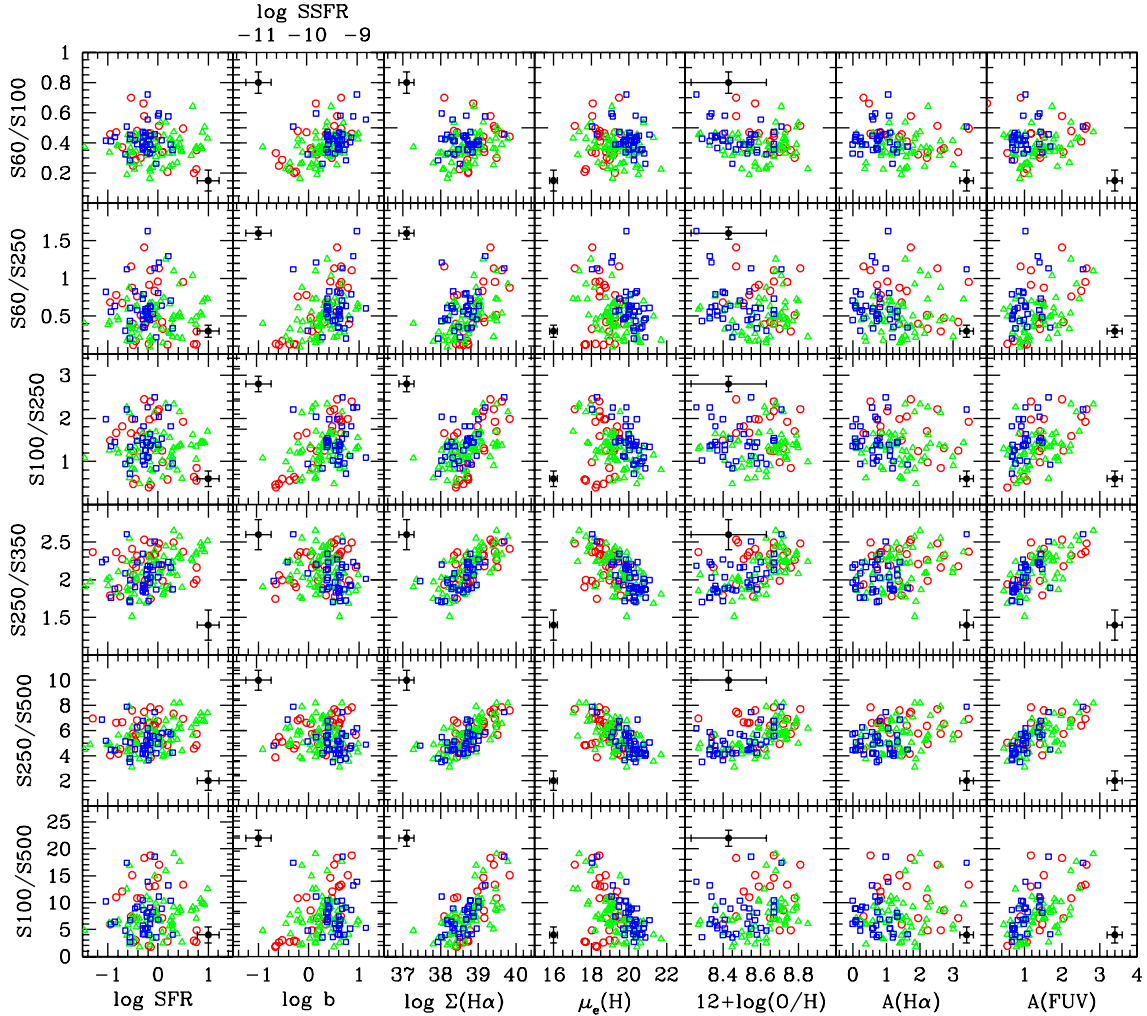


Fig. 3. The relationship between the far infrared colour indices and different tracers of the physical properties of the interstellar medium, *from left to right*: *first column*: the logarithm of the star formation rate (in $M_{\odot} \text{ yr}^{-1}$) measured as described in Boselli et al. (2009); *second column*: the logarithm of the birthrate parameter b (or the specific star formation rate $SSFR$); *third column*: the logarithm of the $H\alpha$ effective surface brightness (in $\text{erg s}^{-1} \text{ kpc}^{-2}$); *fourth column*: the H -band effective surface brightness (in $\text{AB mag arcsec}^{-2}$); *fifth column*: the metallicity index $12 + \log(O/H)$; *sixth column*: the Balmer extinction $A(H\alpha)$ (in magnitudes); *seventh column*: the FUV attenuation $A(\text{FUV})$ (in magnitudes). Red open circles for Sa-Sb, green empty triangles for Sbc-Scd and blue open squares for Sd-Im-BCD. The typical error bar is indicated with a black symbol.

- in forming stars at the present epoch ($b \geq 1$) have, on average, $S60/S100$ flux density ratios ($S60/S100 \approx 0.6$) slightly larger than quiescent objects ($b \leq 1$; $S60/S100 \approx 0.4$). This indicates that, whenever integrated values are used, the contribution of the warm dust component to the total emission of late-type galaxies is principally controlled by the history of star formation (or the hardness of the radiation field) rather than by the present day activity;
- 2) poorly defined and dispersed relations ($r \sim 0.3\text{--}0.6$) are also observed between the peak sensitive $S100/S500$, $S100/S250$ and $S60/S250$ colour indices (in order of decreasing r) and the surface brightness of the ionising ($\Sigma(H\alpha)$) and non-ionising ($\mu_e(H)$) interstellar radiation field and the FUV attenuation ($A(\text{FUV})$). Among these colour indices, only $S60/S250$ and $S100/S250$ ($r \sim 0.4$) correlate very mildly with the birthrate parameter. This evidence suggests that the position of the peak of the dust emission, or in other words the mean temperature of the dust, is governed, as expected, by the general interstellar radiation field;
 - 3) strong correlations ($r \geq 0.6$) are instead observed between the cold dust sensitive $S250/S350$ and $S250/S500$ colour indices and $\Sigma(H\alpha)$, $\mu_e(H)$ and $A(\text{FUV})$. Warmer colours

- are observed in those galaxies with higher radiation fields and higher extinction. The same colour indices are only barely related to the present day star formation activity (SFR ; $r \sim 0.3$) and the metallicity ($r \sim 0.55$). Again, these plots indicate that the emission of the cold dust component is also controlled by the general interstellar radiation field and partly by the metallicity. What is surprising, however, is that the correlations with these cold dust sensitive indices are significantly stronger than with the peak sensitive indices which are sampling a warmer dust component;
- 4) no evident correlations are observed between any of the far infrared colours and the attenuation $A(H\alpha)$;
 - 5) we also do not observe any evident systematic differences among galaxies of different morphological type, maybe with the exception that intermediate type spirals (Sbc-Scd, green empty triangles) covers a larger range in the parameter space than the other classes.

Figure 3 shows that the far infrared colour indices can be correlated at the same time with apparently different tracers of the properties of galaxies. The most striking example are the strong correlations observed between the cold dust sensitive colour

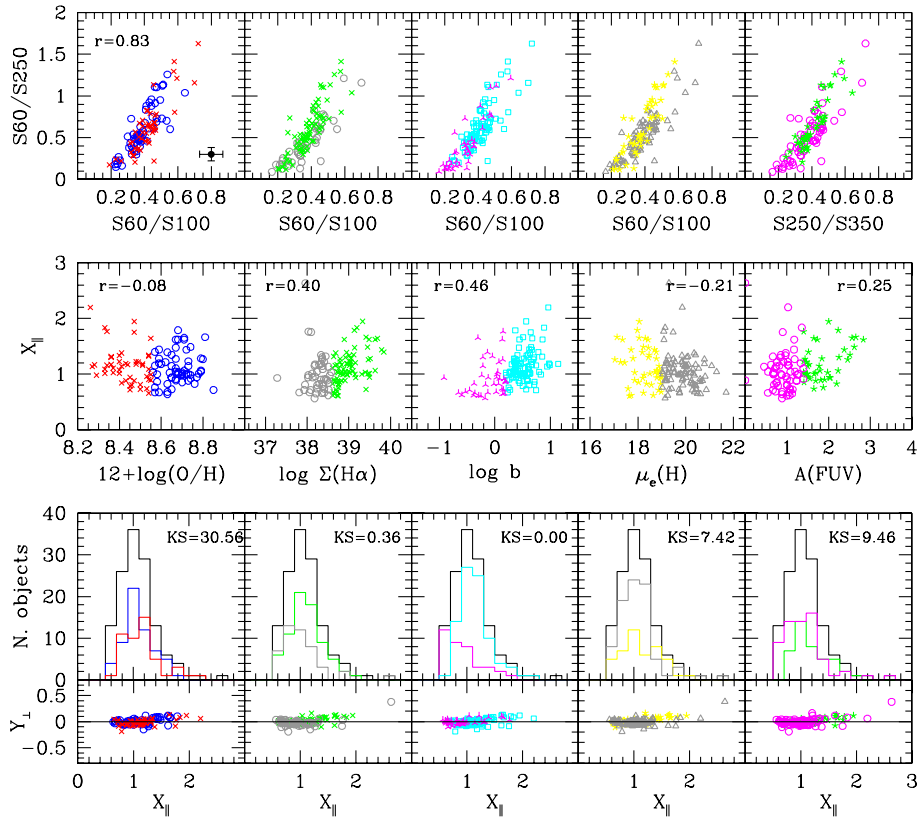
S60/S250 vs. S60/S100_{||}

Fig. 4. *Upper panel:* the $S60/S250$ vs. $S60/S100$ colour–colour relation with galaxies coded according to their physical parameters (from left to right) as a function of the metallicity (red crosses for $12 + \log(\text{O}/\text{H}) \leq 8.55$, blue circles for $12 + \log(\text{O}/\text{H}) > 8.55$), $H\alpha$ surface brightness (grey circles for $\log \Sigma(H\alpha) \leq 38.6 \text{ erg s}^{-1} \text{ kpc}^{-2}$, green crosses for $\log \Sigma(H\alpha) > 38.6 \text{ erg s}^{-1} \text{ kpc}^{-2}$), birthrate parameter (magenta three points stars for $\log b \leq 0.2$, cyan squares for $\log b > 0.2$), H band effective surface brightness (yellow stars for $\mu_e(H) \leq 19 \text{ AB mag arcsec}^{-2}$, grey triangles for $\mu_e(H) > 19 \text{ AB mag arcsec}^{-2}$) and UV attenuation (magenta circles for $A(\text{FUV}) < 1.4 \text{ mag}$, green stars for $A(\text{FUV}) \geq 1.4$). *Middle panel:* the dependence of the position of galaxies along the $S60/S250$ vs. $S60/S100$ colour–colour relation ($X_{||}$) on the different physical parameters. *Lower panel:* the histogram of the distribution of galaxies along the colour–colour relation (*upper*) and the relation between the position of galaxies in the direction perpendicular (Y_{\perp}) and that along ($X_{||}$) the colour–colour relation. r gives the Spearman correlation coefficient of the different relations, KS the probability that the two galaxy populations are driven by the same parent population (Kolmogorov-Smirnov test: for $\text{KS} \leq 5$ the two galaxy populations are statistically different). The de-projection of the $S60/S100$ vs. $S60/S250$ colour–colour relation is done using the measured linear best fit: $S60/S250 = 2.64(\pm 0.16) \times S60/S100 - 0.45(\pm 0.06)$; $r = 0.83$.

indices and $\Sigma(H\alpha)$ and $\mu_e(H)$, the former tracing the surface density of the radiation emitted by newly formed stars, the latter that of the very evolved stars. Although measured using independent data, the different physical parameters used in this analysis are sometime related with each other, as discussed in Appendix A. The mutual relationships between the different physical parameters must be considered for a complete and coherent interpretation of Fig. 3.

To better understand the role of the different physical parameters in the definition of the various colour–colour relations plotted in Fig. 2 we analyse in detail a few colour–colour diagrams selected to represent at the same time those colour indices sensitive to the warm dust component, to the position of the peak of the dust emission, and the cold dust emission. These relations are shown in Figs. 4–7 (upper panel)⁴. To this aim, galaxies are divided into two subsamples and coded according to different

⁴ For such a purpose here we analyse in detail only the less dispersed colour–colour relations shown in Fig. 2 because: i) only here the linear fit (and its residual) can be accurately determined; ii) variations in the more dispersed colour–colour relations can be directly determined from the inspection of Fig. 2 once galaxies are coded according to different physical parameters, as done in Appendix B (see Figs. B.1–B.5).

physical parameters, namely the metal content (with a threshold for selecting galaxies taken at $12 + \log(\text{O}/\text{H}) = 8.55$), the $H\alpha$ surface brightness ($\log \Sigma(H\alpha) = 38.6 \text{ erg s}^{-1} \text{ kpc}^{-2}$), the birthrate parameter ($\log b = 0.2$), the H -band effective surface brightness ($\mu_e(H) = 19 \text{ AB mag arcsec}^{-2}$) and the FUV attenuation ($A(\text{FUV}) = 1.4 \text{ mag}$). We choose these physical parameters because they seem to be the most important in defining the observed relations in Fig. 3. The adopted thresholds have been chosen to split the sample in two sub-samples each with approximately the same number of objects. For simplicity we focus here only on some representative colour–colour relations, referring the interested reader to Appendix B for the rest of the colour–colour diagrams shown in Fig. 2. To quantify the role of the different physical parameters in defining the observed colour–colour relations, we first calculate the best fit of the relations shown in Figs. 4–7 (upper panel). We then use the best fit to de-project these colour–colour relations, and then study how the obtained variables running along ($X_{||}$) or perpendicular (Y_{\perp}) to these relations depend on the different physical parameters (middle panel). We then quantify the difference in the various subsamples of objects through a Kolmogorov-Smirnov test (lower panel). The position of galaxies along or in the direction

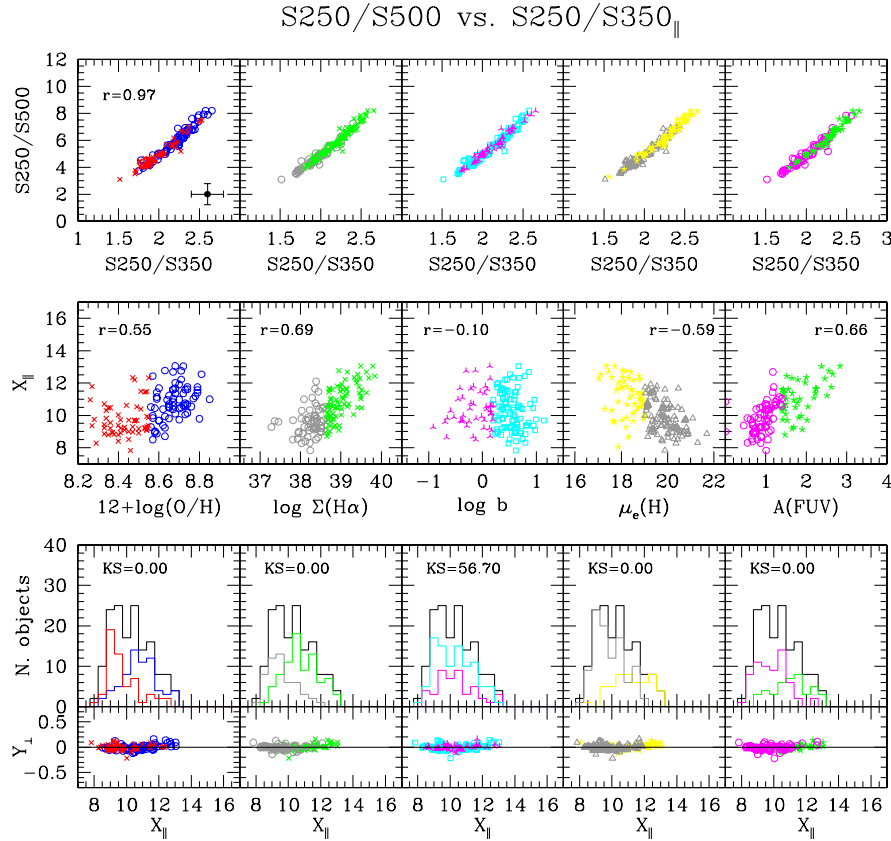


Fig. 5. *Upper panel:* the $S250/S500$ vs. $S250/S350$ colour–colour relation with galaxies coded according to their physical parameters as in Fig. 4. *Middle panel:* the dependence of the position of galaxies along the $S250/S500$ vs. $S250/S350$ colour–colour relation (X_{\parallel}) on the different physical parameters. *Lower panel:* the histogram of the distribution of galaxies along the colour–colour relation (*upper*) and the relation between the position of galaxies in the direction perpendicular (Y_{\perp}) and that along (X_{\parallel}) the colour–colour relation. The de-projection of the $S250/S500$ vs. $S250/S350$ colour–colour relation is done using the measured linear best fit: $S250/S500 = 4.78(\pm 0.09) \times S250/S350 - 4.59(\pm 0.19)$; $r = 0.97$.

perpendicular to that of these colour–colour diagrams indicate variations in their spectral energy distribution due to either a variation of the mean dust temperature or of the dust grain emissivity in the sampled photometric bands.

Figures 4–7 clearly show that galaxies coded with different symbols populate these colour–colour diagrams following a well defined order. Interestingly, galaxies can be at the same time well mixed and clearly segregated in a given colour–colour diagram according to their different physical parameters (indicated by different symbols and colour codes in the figures). An evident example is the mixed vs. segregated distribution of galaxies in Fig. 5 when coded according to their birthrate parameter b or $H\alpha$ surface brightness: while galaxies with different b populate the whole dynamic range of the $S250/S500$ vs. $S250/S350$ relation, those of high $\Sigma(H\alpha)$ have systematically warmer infrared colours than objects with low ionising radiation fields. All together, this is a further evidence that the emission properties of the dust are tightly related to the properties of the interstellar radiation field (hardness, intensity), to the internal attenuation and to the mean metallicity of the interstellar medium in a quite complex manner which can change as a function of wavelength.

A detailed, comparative analysis of Figs. 3–7 at the same time confirms our previous findings and brings to light several new and interesting results:

- 1) the commonly used colour index $S60/S100$, tracer of starburst activity sensitive to the emission of the warm dust component, depends only very marginally on the star formation

history of the galaxy b (or equivalently, on the hardness of the interstellar radiation field) and on the surface brightness of the youngest stars ($\Sigma(H\alpha)$; Fig. 4). Galaxies are indeed located along the $S60/S250$ vs. $S60/S100$ colour–colour relation according to the birthrate parameter b ($r = 0.46$) and $\Sigma(H\alpha)$ ($r = 0.40$);

- 2) galaxies are distributed in the direction perpendicular to the colour–colour relations defined with indices sensitive to the position of the peak of the far infrared emission (Fig. 6) according to all the physical parameters used in this work but b , where any correlation, if present, is very weak ($r = -0.38$);
- 3) the properties of the coldest dust component, as traced by the $S250/S350$ vs. $S250/S500$ colour index, do not depend on the birthrate parameter b (magenta three points stars and cyan open squares are well mixed in Fig. 5). The distribution of galaxies along this colour–colour relation changes with $\Sigma(H\alpha)$ ($r = 0.69$), $\mu_e(H)$ ($r = -0.59$) and $A(FUV)$ ($r = 0.66$) and only very marginally on metallicity ($r = 0.30$);
- 4) in colour–colour relations that are sensitive to the cold dust and their wavelength peak, the perpendicular distance to these trends are only moderately related to a galaxy’s star formation history ($r = 0.49$) and metallicity ($r = -0.38$).

4.4. Comparison with a single modified black body emission

The relationships observed between the different far infrared colour indices can be due either to variations of the mean

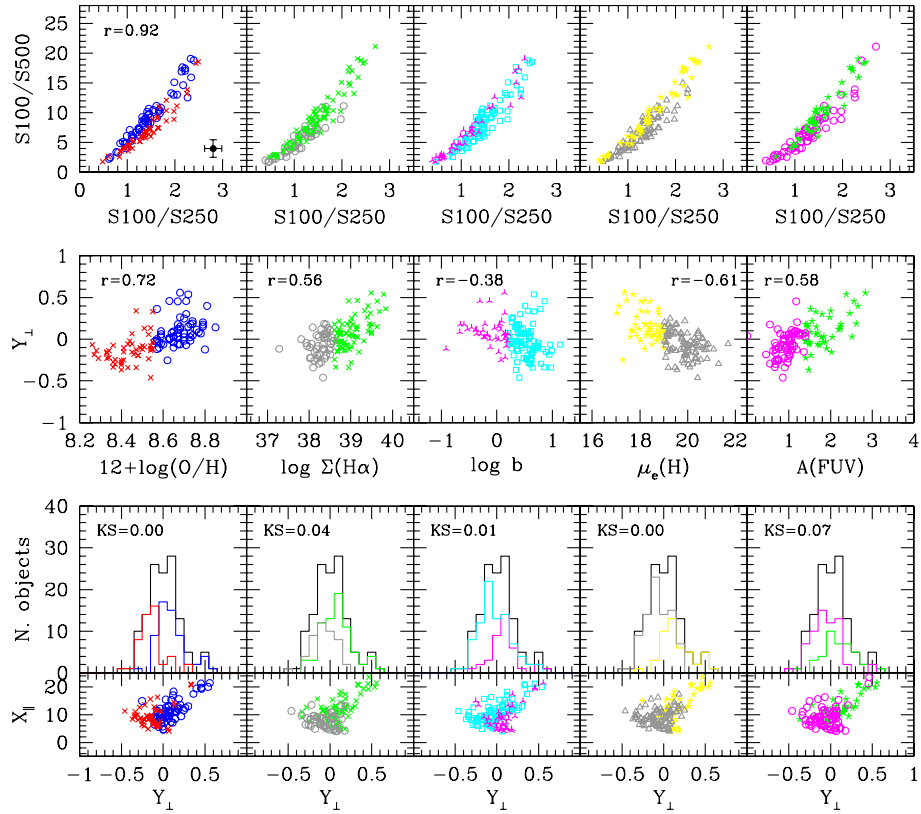
S100/S500 vs. S100/S250_⊥

Fig. 6. *Upper panel:* the $S100/S500$ vs. $S100/S250$ colour–colour relation for galaxies coded according to their physical parameters as in Fig. 4. *Middle panel:* the dependence of the position of galaxies in the direction perpendicular to the $S100/S500$ vs. $S100/S250$ colour–colour relation (Y_{\perp}) on the different physical parameters. *Lower panel:* the histogram of the distribution of galaxies in the direction perpendicular (X_{\parallel}) and that in the direction perpendicular (Y_{\perp}) to the colour–colour relation. The de-projection of the $S100/S500$ vs. $S100/S250$ colour–colour relation is done using the measured linear best fit: $S100/S500 = 7.37(\pm 0.28) \times S100/S250 - 2.23(\pm 0.41)$; $r = 0.92$.

temperature of the emitting dust or to variations in the dust grain emissivity properties that might change in different types of galaxies. Assuming that dust grains are in thermal equilibrium with the radiation, which is probably the case for $\lambda \geq 250 \mu\text{m}$, the far infrared emission of galaxies is generally approximated by one (or more) modified black body $B(\nu, T)$, with a resulting emitted flux density:

$$S_{\text{dust}}(\nu, T) \propto \nu^{\beta} B(\nu, T) \quad (4)$$

where β is the grain emissivity index whose value ranges between ~ 1.5 and ~ 2 (Hildebrand 1983). In particular, a single modified black body with a temperature of the order of ~ 20 K is generally used to reproduce the rapid decrease of the far infrared emission observed at long wavelengths (Reach et al. 1995; Bianchi et al. 1999; Dunne et al. 2000; Dunne & Eales 2001; Bendo et al. 2003; Draine 2003; Davies et al. 2011; Magrini et al. 2011). This assumption is not physical since it is not indicated to represent a continuum in the dust grain distribution in matter of size and temperature (Mathis et al. 1977; Draine & Li 2007). However this simple analytical prescription is often used both in extragalactic and cosmological studies since it reproduces fairly well the observed far infrared spectral energy distributions of galaxies. Our new homogeneous and complete set of data, in particular those obtained by SPIRE in the spectral domain $250\text{--}500 \mu\text{m}$, can be used to see whether this simple modified black body assumption is realistic. At the same time

this dataset can help us to understand which of the two main parameters regulating the dust emissivity, the grain emissivity index β or the dust temperature T , is the main driver of the observed colour–colour far infrared relations and for their scatter. To this aim we plot in Fig. 8 the SPIRE colour–colour indices $S250/S350$ vs. $S350/S500$ for galaxies coded according to the different physical parameters already used in the previous figures and compare them to the expected values obtained for a modified black body with a grain emissivity index $\beta = 2$ (solid line) and $\beta = 1.5$ (dotted line). Clearly, the observed relations can not be represented by a modified black body with a dust grain emissivity index of $\beta = 2$, but are better reproduced when $\beta = 1.5$. Systematic errors of the order of $\sim 15\%$ on a colour are necessary to shift the data onto a $\beta = 2$ single black body. This result seems thus robust vs. calibration uncertainties, that are at present estimated to be of the order of $\simeq 7\%$ for each band, or colour corrections ($\leq 6.5\%$). As expected, these relationships might just be due to a variation of the mean temperature of the dust grains which, for a $\beta = 1.5$, might range from ~ 15 to ~ 50 K (from ~ 15 to ~ 30 K for $\beta = 2$). These values of β and T are consistent with that found in the Milky Way using COBE data by Reach et al. (1995) or in other nearby galaxies observed in the submillimetre domain (Bianchi et al. 1998, 1999; Amblard et al. 2010; Hwang et al. 2010; Dunne et al. 2011; Planck Collaboration 2011b).

The observed relationship between the two colour indices, however, has a slope slightly flatter than the one predicted by a

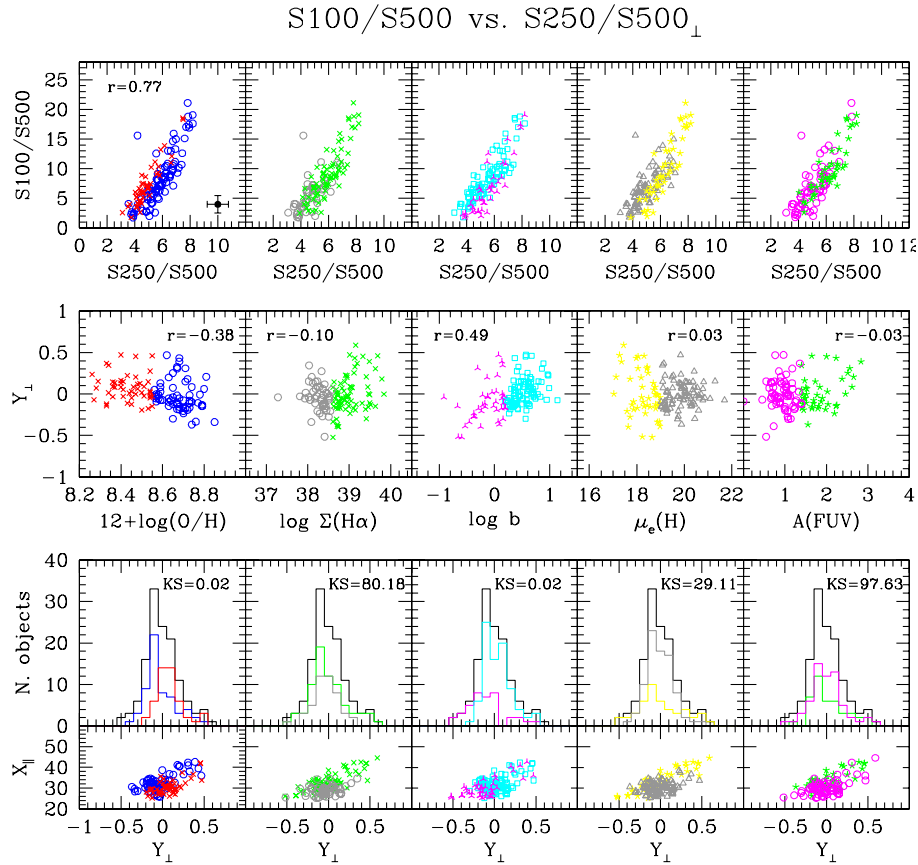


Fig. 7. *Upper panel:* the $S100/S500$ vs. $S250/S500$ colour–colour relation for galaxies coded according to their physical parameters as in Fig. 4. *Middle panel:* the dependence of the position of galaxies in the direction perpendicular to the $S100/S500$ vs. $S250/S500$ colour–colour relation (Y_{\perp}) on the different physical parameters. *Lower panel:* the histogram of the distribution of galaxies in the direction perpendicular to the colour–colour relation (upper) and the relation between the position of galaxies along (X_{\parallel}) and that in the direction perpendicular (Y_{\perp}) to the colour–colour relation. The de-projection of the $S100/S500$ vs. $S250/S500$ colour–colour relation is done using the measured linear best fit: $S100/S500 = 12.22(\pm 0.90) \times S250/S500 - 23.43(\pm 2.32)$; $r = 0.77$.

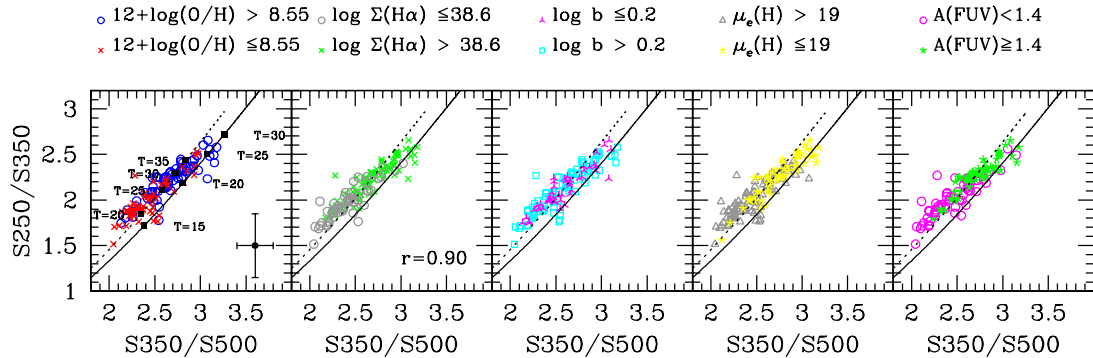


Fig. 8. The relationship between the SPIRE colour–colour indices $S250/S350$ vs. $S350/S500$ for galaxies coded according to their different physical parameters as in Fig. 4. The colour–colour relation is compared to the expected relations obtained for a modified black body with a grain emissivity index $\beta = 2$ (solid line) and $\beta = 1.5$ (dotted line). Black squares indicate different temperatures for the two modified black bodies (*left panel*). The Spearman correlation coefficient of this relation is $r = 0.90$.

single modified black body of fixed grain emissivity index. We can thus not exclude values of $\beta > 1.5$ in galaxies with the highest flux density ratios. Similarly, in the objects with the lowest flux density ratios, β might be < 1.5 . This evidence might indicate that both β and T vary along the sequence. Indeed, there is a quite strong degeneracy between β and T given that at these low temperatures a variation of T of ~ 5 K is almost equivalent

to a variation of β of ~ 0.5 . Furthermore, we have to remember that β and T might also be inversely correlated (Desert et al. 2008; Shetty et al. 2009; Veneziani et al. 2010; Anderson et al. 2010; Planck Collaboration 2011a; Bracco et al. 2011).

By identifying galaxies according to their physical properties we have shown that high flux density ratios are mainly observed in objects with strong ionising radiation fields, which are

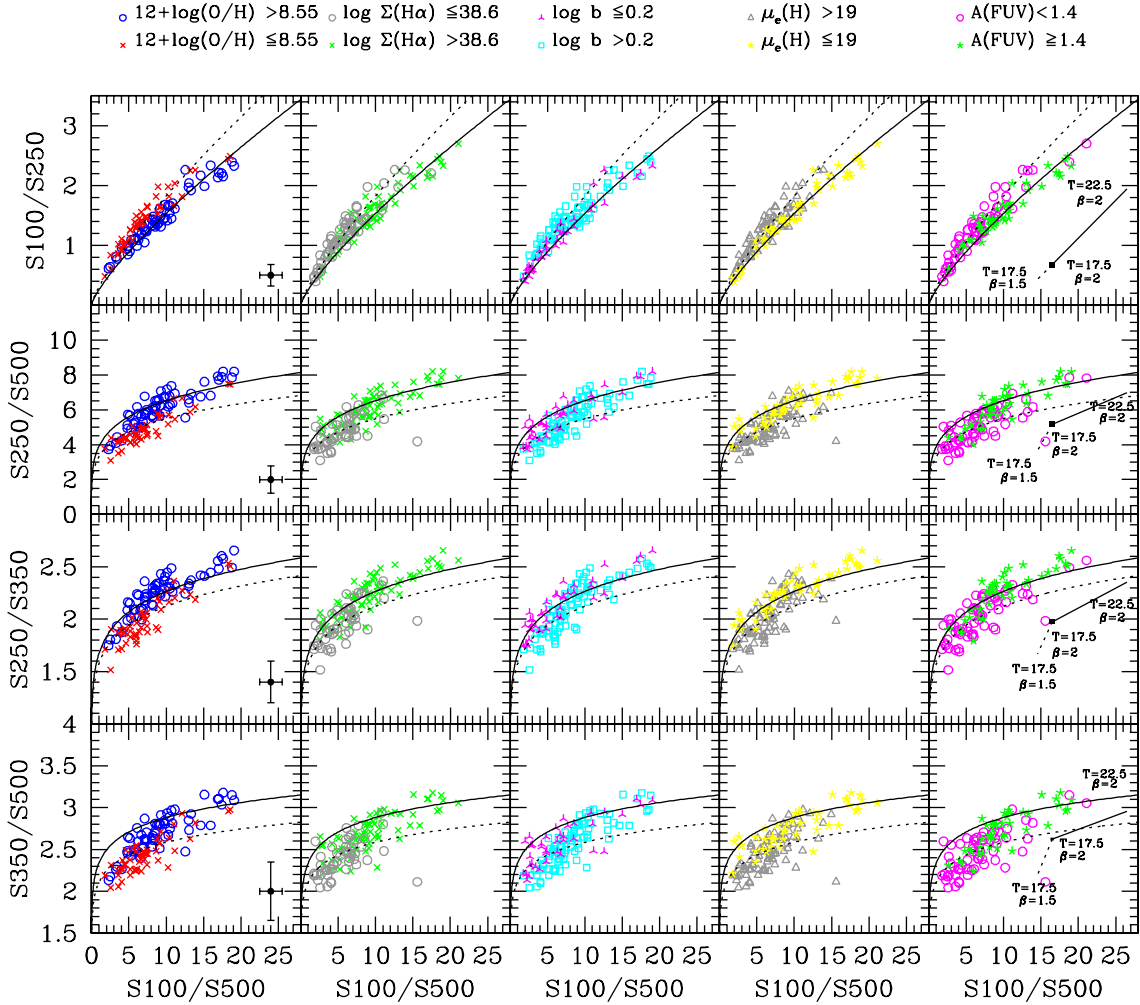


Fig. 9. The comparison of relationships between different far infrared colour indices for galaxies coded according to their different physical parameters as in Fig. 4. The solid and dotted lines indicate the expected ratios for a single modified black body with grain emissivity $\beta = 2$ and 1.5, respectively. The short solid and dotted lines in the last column show the expected variations when the temperature T changes from 17.5 to 22.5 K for $\beta = 2$ (solid line), and from 17.5 K and $\beta = 2$ to $T = 17.5$ K and $\beta = 1.5$ (dotted line). The black square is the expected value (shifted on the X-axis for clarity) for a black body of $T = 17.5$ K and $\beta = 2$.

generally also metal rich galaxies with strong non-ionising stellar radiation fields. If this single modified black body configuration is valid, the grain emissivity index of metal poor objects with low ionising radiation fields is $\beta \lesssim 1.5$. It is worth mentioning, however, that all the previous considerations on the cold dust emission are valid if the contribution of the warm dust component to the far infrared emission ($\lambda \geq 250 \mu\text{m}$) is negligible, which is probably the case given that our sample is composed of normal, late-type galaxies and do not include strong starbursts such as M 82 or Arp 220 (Boselli et al. 2010a).

The same analysis can be done by comparing other far infrared colour–colour relations defined for galaxies coded according to the same physical parameters (Fig. 9). Figure 9 confirms the trend observed in the previous figures with the different physical parameters. Only the $S100/S250$ vs. $S100/S500$ colour–colour relation for metal poor galaxies or for objects with a high $H\alpha$ or H band surface brightness (upper row in Fig. 9) can be easily reproduced by a single black body with a grain emissivity parameter $\beta = 2$. It is however clear from both Figs. 8 and 9 that all the other colour–colour relations made using at least one cold dust sensitive colour index cannot be well reproduced by a single modified black body. This evidence confirms that either the grain emissivity parameter β or the temperature significantly change

from galaxy to galaxy, as indeed expected given the variety of physical conditions that characterise these objects. Fitting the infrared spectral energy distribution with different dust models might help us to make a further step in understanding the very nature of the emitting dust. This will be the topic of a future communication.

5. Discussion

The analysis of the far infrared colours of nearby late-type galaxies have been presented in the literature using data taken from the IRAS (e.g. Helou 1986) and *Spitzer* (Dale et al. 2005, 2007). This is the first dedicated work on the analysis of the far infrared colours of nearby late-type galaxies in the SPIRE/*Herschel* bands. Our results can thus be compared only to those obtained for a few nearby galaxies recently observed by *Herschel*. Variations of the infrared colour indices with the $H\alpha$ and the H -band surface brightness similar to those presented in this work have been also observed on galactic scales over the disc of M 81, M 83, NGC 2403 (Bendo et al. 2011) and of M 33 (Boquien et al. 2011). The analysis of the *Spitzer* and *Herschel* data of these well known galaxies confirmed that while the far infrared emission in the 70–160 μm spectral range is principally governed

by the activity of star formation, the cold dust responsible for the emission at longer wavelengths is predominantly heated by the general interstellar radiation field. Variations of the far infrared colour indices as a function of metallicity have been also reported on galactic scales in M99 and M100 by Pohlen et al. (2010) and in unresolved objects by Boselli et al. (2010b).

The new observational evidence presented in this work puts a strong constraint for the study of the physical properties of the dust grains emitting in the far infrared spectral domain. It shows that the properties of the dust emitting in this spectral domain are not universal but rather vary in galaxies due to their large range of parameter space. Qualitatively, the results discussed in the previous section can be interpreted as follow:

- 1) The colour–colour diagrams and the relationships between colours and the different physical parameters are tighter when SPIRE bands are used instead of the 60 and 100 μm IRAS bands. This evidence might be related to a lower photometric quality of the IRAS data with respect to the new SPIRE observations, or to the larger wavelength range sampled by SPIRE (250–500 μm) with respect to IRAS (60–100 μm).
- 2) The weak dependence of S_{60}/S_{100} on the birthrate parameter b (Figs. 3 and 4) is a further indication of a well known result, i.e. that this colour index is related to the star formation history of galaxies. This tendency, however, is much less clear than that observed in other samples (e.g. Chapman et al. 2003) just because the dynamic range in the star formation history of mass-selected HRS galaxies is quite small ($0.1 \lesssim SFR \lesssim 10 M_{\odot} \text{yr}^{-1}$). Indeed the increase of the S_{60}/S_{100} ratio as a function of the infrared luminosity reported in Chapman et al. (2003) is evident only for total infrared luminosities $L_{\text{TIR}} \geq 10^{10} L_{\odot}$, which roughly corresponds to the upper limit in the infrared luminosity sampled by the HRS galaxies (Boselli et al. 2010a).
- 3) The dependence of the SPIRE colour indices S_{250}/S_{350} , S_{350}/S_{500} and S_{250}/S_{500} on metallicity (Fig. 3) is probably a different indication of the widely known submillimetre excess in metal poor, low surface brightness dwarf galaxies (e.g. Galametz et al. 2011). Indeed, these far infrared flux density ratios are low in metal poor objects characterised by a low FUV attenuation and a low ionising and non-ionising radiation field. The observed variations in the colour–colour diagrams might be due to variations in the grain emissivity parameter β . A low value of β has been already proposed to explain the excess of the submillimetre emission at $\lambda \geq 500 \mu\text{m}$ observed in metal poor, dwarf galaxies (Galliano et al. 2005; Bendo et al. 2006; Galametz et al. 2010; O’Halloran et al. 2010; Boselli et al. 2010b). Silicates and amorphous carbons dominate the emission at long wavelengths, so if their fraction changes, different values of β might be expected (Compiegne et al. 2011). Stellar evolution models, however, do not predict for a given metallicity strong differences in the relative production of C and Si for galaxies undergoing a different star formation history that might justify a systematic difference in the abundance of silicate and carbonaceous dust grains. Despite a large uncertainty in the dust production and destruction processes, models of dust formation in galaxies determined for different star formation histories do not predict a significant difference in the fraction of carbonaceous and silicate dust grains (Dwek 1998; Calura et al. 2008). The excess of the cold dust emission in low luminosity, metal poor galaxies might have also other origins, as extensively discussed in Galametz et al. (2011).

These include a dependence of the dust grain emissivity with temperature (Dupac et al. 2003; Meny et al. 2007; Desert et al. 2008), variation of the grain emissivity with the fractal aggregation of individual amorphous grains inside molecular clouds (Paradis et al. 2009), grain coagulation (Bazell & Dwek 1990), overabundance of very small grains with grain emissivity $\beta = 1$ in extreme environments (Lisenfeld et al. 2001; Zhu et al. 2009) or spinning dust (Ferrara & Dettmar 1994).

- 4) The presence of parallel colour–colour relations for galaxies of different physical parameters in the diagrams done with flux density ratios sensitive to the position of the peak (defined by the different symbols in the relations shown in Fig. 6) indicate that the peak of the dust emission in the spectral energy distribution, and thus the mean temperature of the dust, shifts from short to long wavelengths as a function of these physical parameters.
- 5) Far infrared colours ($\lambda \geq 250 \mu\text{m}$) change with the surface brightness of the ionising and non-ionising radiation, and the FUV dust attenuation (Figs. 3, 5). This result is expected since the integrated emission of galaxies includes both HII regions along the spiral arms, where the heating sources are mainly young and massive OB stars, and the more quiescent interarm regions where dust is principally heated by the general interstellar radiation field. Although $A(\text{FUV})$ and the far infrared colour indices are not completely independent entities since $A(\text{FUV})$ is measured using the total far infrared to UV flux ratio ($A(\text{FUV}) \propto L_{\text{FIR}}/L_{\text{FUV}}$), the strong correlations observed in the S_{250}/S_{500} and S_{250}/S_{350} vs. $A(\text{FUV})$ cannot result from this underlying relation. This is due to the fact that i) the far infrared luminosity L_{FIR} used to estimate $A(\text{FUV})$ strongly depends on the total far infrared emission at $\lambda \lesssim 100 \mu\text{m}$, where most of the energy is radiated and only weakly on the dust emission at longer wavelengths; and ii) the SPIRE bands flux density ratio only depend at a second order on the total infrared luminosity that in our sample spans only three orders of magnitude. This last consideration is also valid for the S_{60}/S_{100} colour index. The dependences of the IRAS/SPIRE flux density ratio on $A(\text{FUV})$, on the contrary, might partly result from the fact that only IRAS flux densities are strongly related to the total far infrared luminosity of galaxies. The weak dependence with metallicity observed in Fig. 5 suggests that the far infrared colours of galaxies might also depend on the chemical properties of the dust grains.

Overall, this analysis indicates that the emission properties of the cold dust dominating the far infrared spectral domain are regulated by the properties of the interstellar radiation field. Consistent with that observed in nearby, resolved galaxies (Bendo et al. 2011; Boquien et al. 2011), our analysis has shown that the ionising and the non-ionising stellar radiation, including that emitted by the most evolved, cold stars, both contribute to the heating of the cold dust component. This work, however, indicates that the mean metallicity of the gaseous phase of the interstellar medium is another key parameter characterising the cold dust emission of normal, late-type galaxies.

The simple interpretation outlined in this section, however, although based on solid statistical observational results, remains in a speculative stage given the strong degeneracy in the different parameters responsible for the emission of dust (dust grain emissivity, temperature, composition), and could be confirmed only after an accurate comparison with model predictions. It is however clear that this observational evidence will be a major

constraint in the years to come for models of dust emission in galaxies.

Considering the grain emissivity parameter, we obtain consistent results ($\beta \approx 1.5$) with those obtained using similar *Herschel* data in other nearby galaxies such as the LMC (Gordon et al. 2010) and M 33 (Kramer et al. 2010). Values of $\beta = 1.2, 1.5$ and 1.8 have been determined by the Planck collaboration for the SMC, LMC and the Milky Way, respectively (Planck Collaboration 2011c). Our value is however slightly larger than the one obtained by Planck for external galaxies once their SED are fitted with a single modified black body ($\beta = 1.2, T = 26.3$ K; Planck Collaboration 2011b). All these values of $\beta \approx 1.5$ are significantly smaller than those obtained including far infrared data at shorter wavelengths ($\lambda \leq 250 \mu\text{m}$) on a similar set of galaxies ($\beta = 2$) by Davies et al. (2011; see also Fig. 9). A decrease of β with λ has been previously reported by Paradis et al. (2009). They are, however, perfectly consistent with the values obtained by Dale et al. (2012) to fit the 100–500 μm SED of KINGFISH galaxies.

6. Conclusions

We have studied for the first time the far infrared (60–500 μm) colours of normal, late-type galaxies using new data recently obtained by the *Herschel* space mission. We have determined the relationships between different far infrared colour indices, defined as flux density ratios using the two IRAS bands at 60 and 100 μm and the three SPIRE bands at 250, 350 and 500 μm , and compared them to various tracers of the physical properties of the target galaxies. These are the present day star formation activity *SFR*, the birthrate parameter *b* or equivalently the specific star formation rate *SSFR*, the surface brightness of the bulk of the stellar emission $\mu_e(H)$ and that of the ionising stellar radiation $\Sigma(H\alpha)$, the metallicity $12+\log(O/H)$, and the internal attenuation $A(H\alpha)$ and $A(FUV)$. Our analysis has shown that:

- These far infrared colour indices are correlated with each other. Tight correlations are observed only for infrared indices close in λ .
- The far infrared colour indices are only partly related to the different tracers of the physical properties of the analysed galaxies. A tight correlation is observed only between the cold dust sensitive flux density ratios S_{250}/S_{350} and S_{250}/S_{500} and the surface brightness of both the ionising and non-ionising radiation and $A(FUV)$. More dispersed relations are observed between the flux density ratios sensitive to the position of the peak of the far infrared emission (S_{60}/S_{250} , S_{100}/S_{250} , S_{100}/S_{500}) and $\Sigma(H\alpha)$ and $\mu_e(H)$. Coarse relations are observed between the cold dust colour indices S_{250}/S_{350} and S_{250}/S_{500} and the metallicity $12 + \log(O/H)$, while the relationships between the other far infrared colours and the other physical tracers, if present, are very dispersed.
- We have shown that galaxies can be well segregated in the direction perpendicular to the main colour–colour relations determined using far infrared bands ($\lambda \geq 100 \mu\text{m}$) by means of their physical parameters, whose relative importance changes with the sampled spectral domain. Among these, the metallicity and the history of star formation, or equivalently the hardness of the interstellar radiation field, seem the most important parameters. The variation along the observed colour–colour relations, on the contrary, depends only weakly on all the physical parameters.
- We have also shown that a single modified black body with a grain emissivity index $\beta = 1.5$ better fits the observed

SPIRE flux density ratios S_{250}/S_{350} vs. S_{350}/S_{500} than $\beta = 2$. Values of β close to 2 are possible only in metal rich, high surface brightness galaxies, while $\beta \lesssim 1.5$ are more representative of metal poor, low surface brightness objects.

These results are strong constraints for the study of the physical properties of the dust grains emitting in the far infrared spectral domain. They first show that the properties of the large dust grains responsible for the emission in this spectral domain are not universal but rather change according to the physical properties of galaxies. The metallicity, the intensity of the ionising and non-ionising radiation fields as well as the hardness of the incident radiation are indeed key parameters in regulating the dust emission in the far infrared spectral domain.

The observational evidence presented in this work is statistically significant. Its interpretation, however, here done only on phenomenological bases, requires an accurate comparison with different dust models for a more complete understanding of the properties of the cold dust dominating the emission of galaxies in the far infrared spectral domain. To this purpose, we have collected *Spitzer* data at shorter wavelengths ($3 \leq \lambda \leq 160 \mu\text{m}$) to map the dust emission on the widest possible spectral range. These data, that will be soon combined with new PACS observations (Davies et al. 2011), will allow us to reconstruct the UV to radio continuum spectral energy distribution of this complete sample of galaxies necessary for a combined study of the heating sources and of the emitting dust (Ciesla et al., Boquien et al., in prep.). The comparison of the observed spectral energy distributions with different models available in the literature will definitely help us in making a further step in the comprehension of the emitting properties of the interstellar dust and their interrelationships with the other physical parameters characterising the properties of the interstellar medium. The first observational evidence presented in this work might have strong implications in the study of galaxies at high redshift. They indeed show that the far infrared colours of galaxies, and thus their spectral energy distributions, might change as a function of the intensity of the interstellar radiation field and the metallicity, physical parameters that strongly evolve during the cosmic time. If possible variations in the star formation activity, and indirectly of the stellar radiation field, are often considered in cosmological studies (e.g. Elbaz et al. 2011), metallicity variations are still neglected. At the same time our results indicate that the assumption of a single modified black body with a grain emissivity parameter $\beta = 1.5$ used to characterise the far infrared-submillimetre emission of galaxies at high redshift (e.g. Magdis et al. 2010a,b; Hwang et al. 2010) is justified by the properties of nearby objects.

Acknowledgements. We thank the referee, D. Dale, for precious comments which helped improving the quality of the manuscript. SPIRE has been developed by a consortium of institutes led by Cardiff Univ. (UK) and including Univ. Lethbridge (Canada); NAOC (China); CEA, LAM (France); IFSI, Univ. Padua (Italy); IAC (Spain); Stockholm Observatory (Sweden); Imperial College London, RAL, UCL-MSSL, UKATC, Univ. Sussex (UK); Caltech, JPL, NHSC, Univ. Colorado (USA). This development has been supported by national funding agencies: CSA (Canada); NAOC (China); CEA, CNES, CNRS (France); ASI (Italy); MCINN (Spain); SNSB (Sweden); STFC, UKSA (UK); and NASA (USA). This research has made use of the NASA/IPAC Extragalactic Database (NED) which is operated by the Jet Propulsion Laboratory, California Institute of Technology, under contract with the National Aeronautics and Space Administration. The research leading to these results has received funding from the European Community's Seventh Framework Programme (FP7/2007-2013) under grant agreement No. 229517. This research has made use of the NASA/IPAC Extragalactic Database (NED) which is operated by the Jet Propulsion Laboratory, California Institute of Technology, under contract with the National Aeronautics and Space Administration and of the GOLDMine database (<http://goldmine.mib.infn.it/>).

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Appendix A: Mutual correlations among the different physical parameters

The physical parameters used to trace the properties of the galaxies analysed in this work, although determined using independent data, can trace non completely independent variables. The *SFR*, the birthrate parameter b and the surface brightness of the ionising radiation $\Sigma(\text{H}\alpha)$, for instance, are all strongly related to the present day star formation activity. The effective surface brightness $\mu_e(H)$ traces the distribution of the evolved stars dominating the stellar mass of galaxies. The total stellar mass is also necessary to estimate the birthrate parameter. $A(\text{H}\alpha)$ and $A(\text{FUV})$ are two independent tracers of the dust extinction within galaxies. These different physical parameters might thus be mutually related, as shown in Fig. A.1 and Table A.1.

Figure A.1 shows a tight correlation between the surface brightness of the very evolved, cold stars $\mu_e(H)$ and the UV attenuation $A(\text{FUV})$ ($r = -0.57$) indicating that in normal, star forming galaxies dust absorbs at the same time the photons emitted by the newly formed stars and by the most evolved stellar populations. $\Sigma(\text{H}\alpha)$, *SFR* and b , as previously mentioned, are not fully independent tracers since all related with the present day star formation activity of galaxies. Surprisingly, however, the mutual relations between these parameters are very dispersed. The correlation between $\mu_e(H)$ and $\Sigma(\text{H}\alpha)$ suggests that the surface density of all kind of stars increases independently of their age in the galaxies analysed in this work. Naturally, the two

independent dust extinction tracers $A(\text{H}\alpha)$ and $A(\text{FUV})$ are mutually related (e.g. Calzetti 2001; Boselli et al. 2009). The trends observed between $A(\text{FUV})$ (and $A(\text{H}\alpha)$) and $\Sigma(\text{H}\alpha)$ or $\mu_e(H)$ suggest that all the stellar radiation participates to the heating of the dust emitting in the far infrared. The relations with the metallicity index $12 + \log(\text{O}/\text{H})$ are weak ($r \approx 0.4\text{--}0.5$), and suggest, as expected, that the dust extinction is more important in metal rich objects. They also show that the metallicity is coupled with the star formation history of galaxies (b), and is generally higher in high surface brightness evolved systems than in star forming, low surface brightness objects.

Appendix B: Colour–colour diagrams vs. physical parameters

To extend the analysis done in Sect. 4, we plot here for completeness all the colour–colour diagrams shown in Fig. 2 for galaxies coded according to different physical parameters. The codes used are the same than those used in Figs. 4–7, where the threshold in the physical parameters are: metal content $12 + \log(\text{O}/\text{H}) = 8.55$; $\text{H}\alpha$ surface brightness $\log \Sigma(\text{H}\alpha) = 38.6 \text{ erg s}^{-1} \text{ kpc}^{-2}$; birthrate parameter $\log b = 0.2$; H -band effective surface brightness $\mu_e(H) = 19 \text{ AB mag arcsec}^{-2}$; FUV attenuation $A(\text{FUV}) = 1.4 \text{ mag}$. The general behaviors described in Sect. 4 can be observed also in the Figs. B.1–B.5.

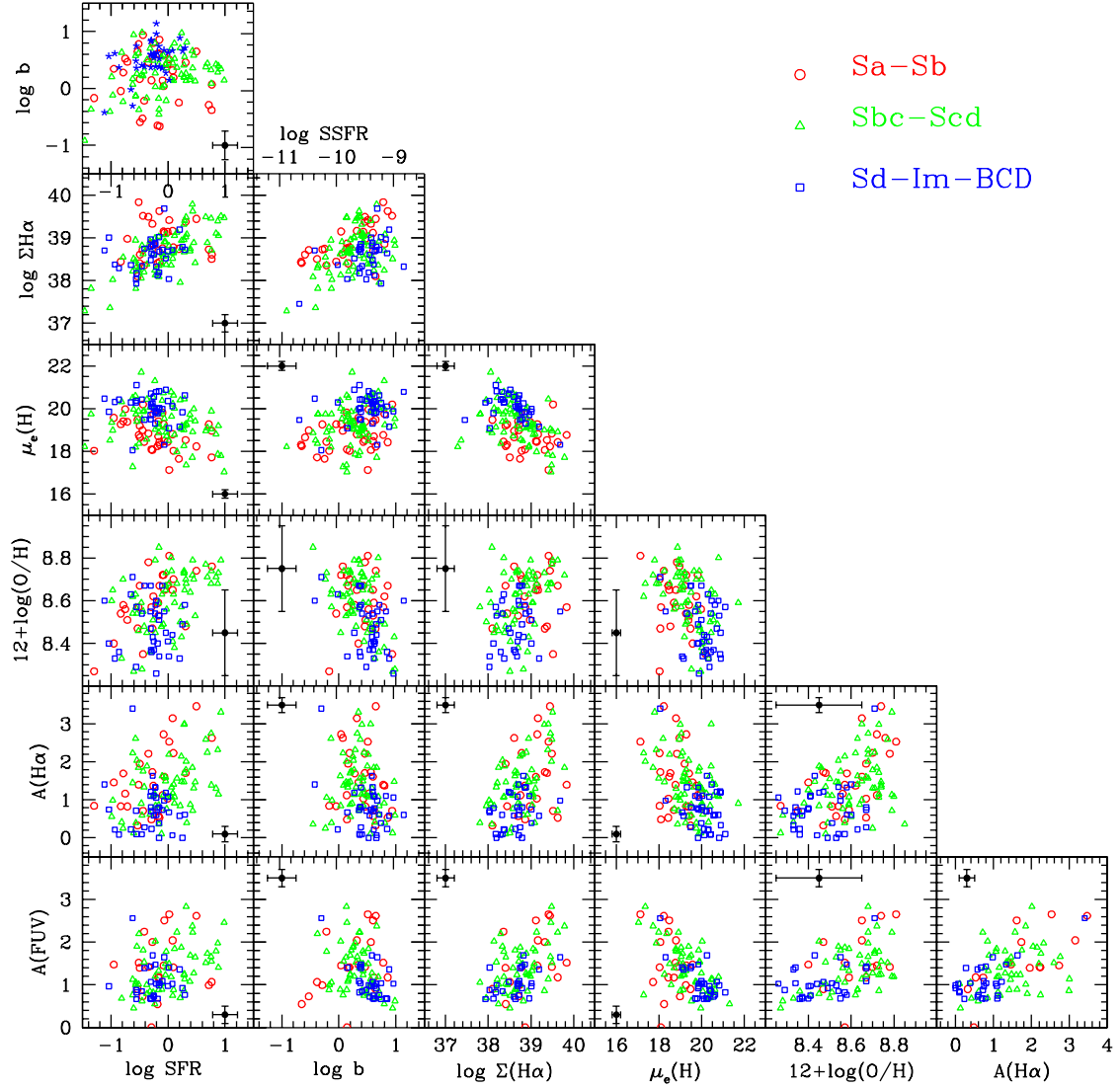


Fig. A.1. The relationships between the different parameters used to trace the physical properties characterizing the target galaxies: the logarithm of the star formation rate SFR (in $M_{\odot} \text{ yr}^{-1}$), the logarithm of the birthrate parameter b (or $SSFR$), the logarithm of the $H\alpha$ effective surface brightness (in $\text{erg s}^{-1} \text{ kpc}^{-2}$), the H -band effective surface brightness (in $\text{AB mag arcsec}^{-2}$), the metallicity index $12 + \log(O/H)$, the Balmer decrement $A(H\alpha)$ (in magnitudes) and the FUV attenuation $A(\text{FUV})$ (in magnitudes). Red open circles for Sa-Sb, green empty triangles for Sbc-Scd and blue open squares for Sd-Im-BCD. The typical error bar is indicated with a black cross.

Table A.1. Spearman correlation coefficients of the relations between the different physical parameters (Fig. A.1).

Y -colour	$\log SFR$	$\log b$	$\log \Sigma(H\alpha)$	$\mu_e(H)$	$12 + \log(O/H)$	$A(H\alpha)$
Units	$M_{\odot} \text{ yr}^{-1}$		$\text{erg s}^{-1} \text{ kpc}^{-2}$	$\text{AB mag arcsec}^{-2}$		mag
$\log b$	0.11					
$\log \Sigma(H\alpha)$	0.56	0.24				
$\mu_e(H)$	-0.25	0.35	-0.39			
$12 + \log(O/H)$	0.44	-0.55	0.43	-0.50		
$A(H\alpha)$	0.34	-0.28	0.33	-0.43	0.47	
$A(\text{FUV})$	0.36	-0.38	0.45	-0.57	0.54	0.56

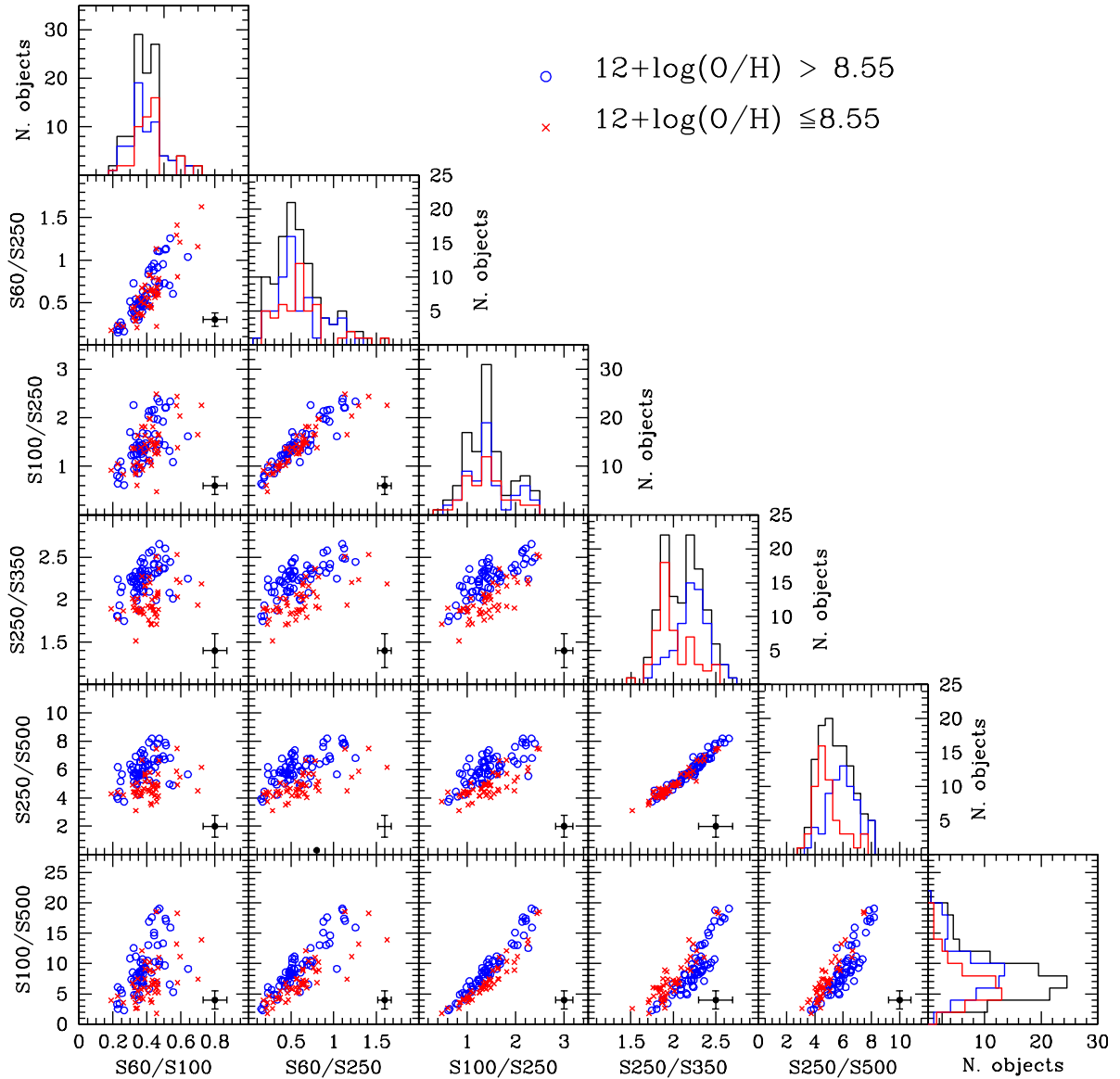


Fig. B.1. Far infrared colour–colour diagrams, equivalent to those shown in Fig. 2, with galaxies coded according to their mean metal content, with blue open circles for metal rich galaxies ($12 + \log(O/H) > 8.55$) and red crosses for metal poor objects ($12 + \log(O/H) \leq 8.55$). The black histogram gives the distribution of all galaxies along the X-axis, while the coloured histograms those of the two subsamples of galaxies selected according to their metallicity.

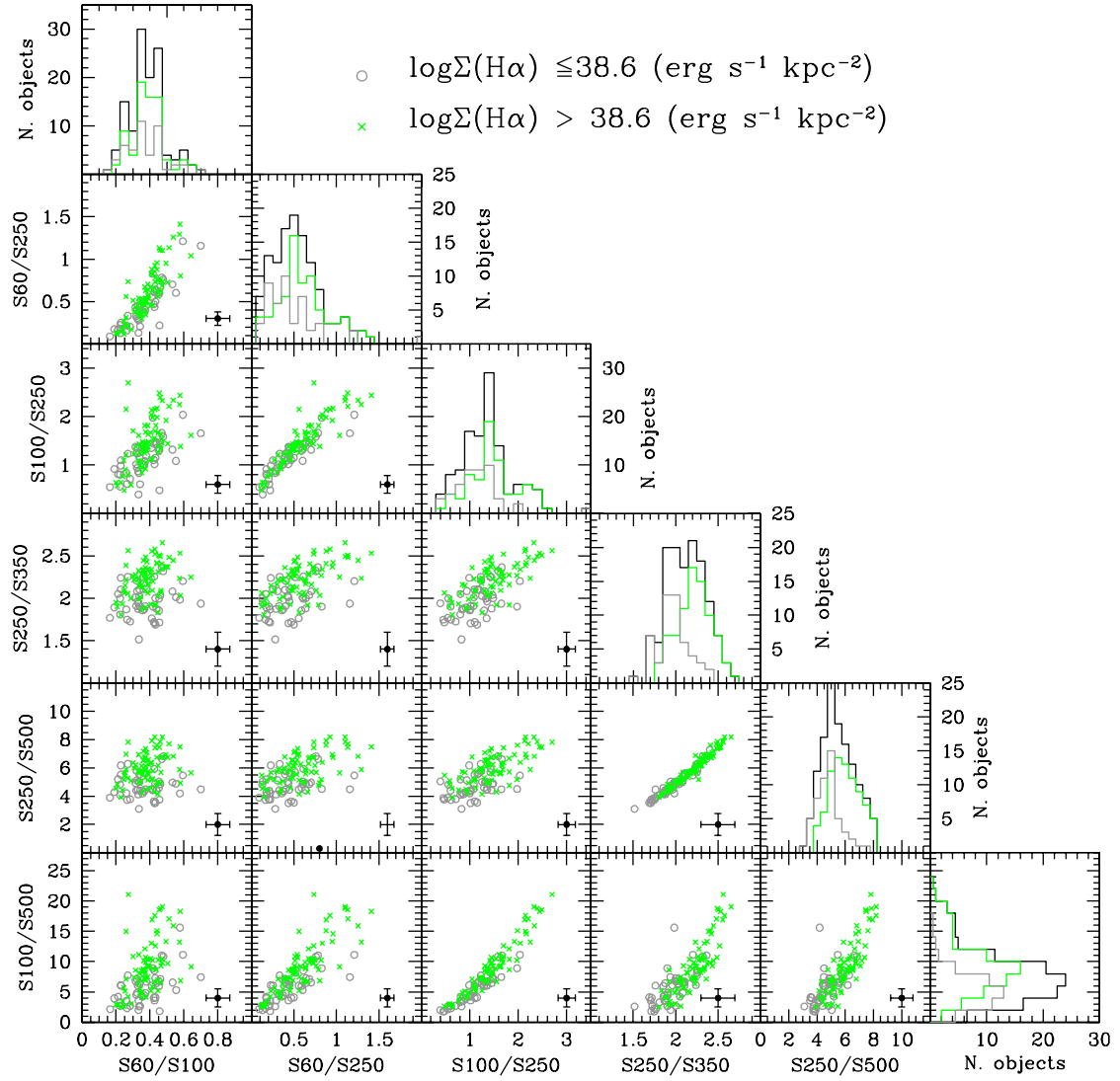


Fig. B.2. Far infrared colour-colour diagrams, equivalent to those shown in Fig. 2, with galaxies coded according to their $\text{H}\alpha$ surface brightness, with green crosses for objects with $\log \Sigma(\text{H}\alpha) > 38.6 \text{ erg s}^{-1} \text{ kpc}^{-2}$ and grey open circles for galaxies with $\log \Sigma(\text{H}\alpha) \leq 38.6 \text{ erg s}^{-1} \text{ kpc}^{-2}$. The black histogram gives the distribution of all galaxies along the X-axis, while the coloured histograms those of the two subsamples of galaxies selected according to their $\text{H}\alpha$ surface brightness.

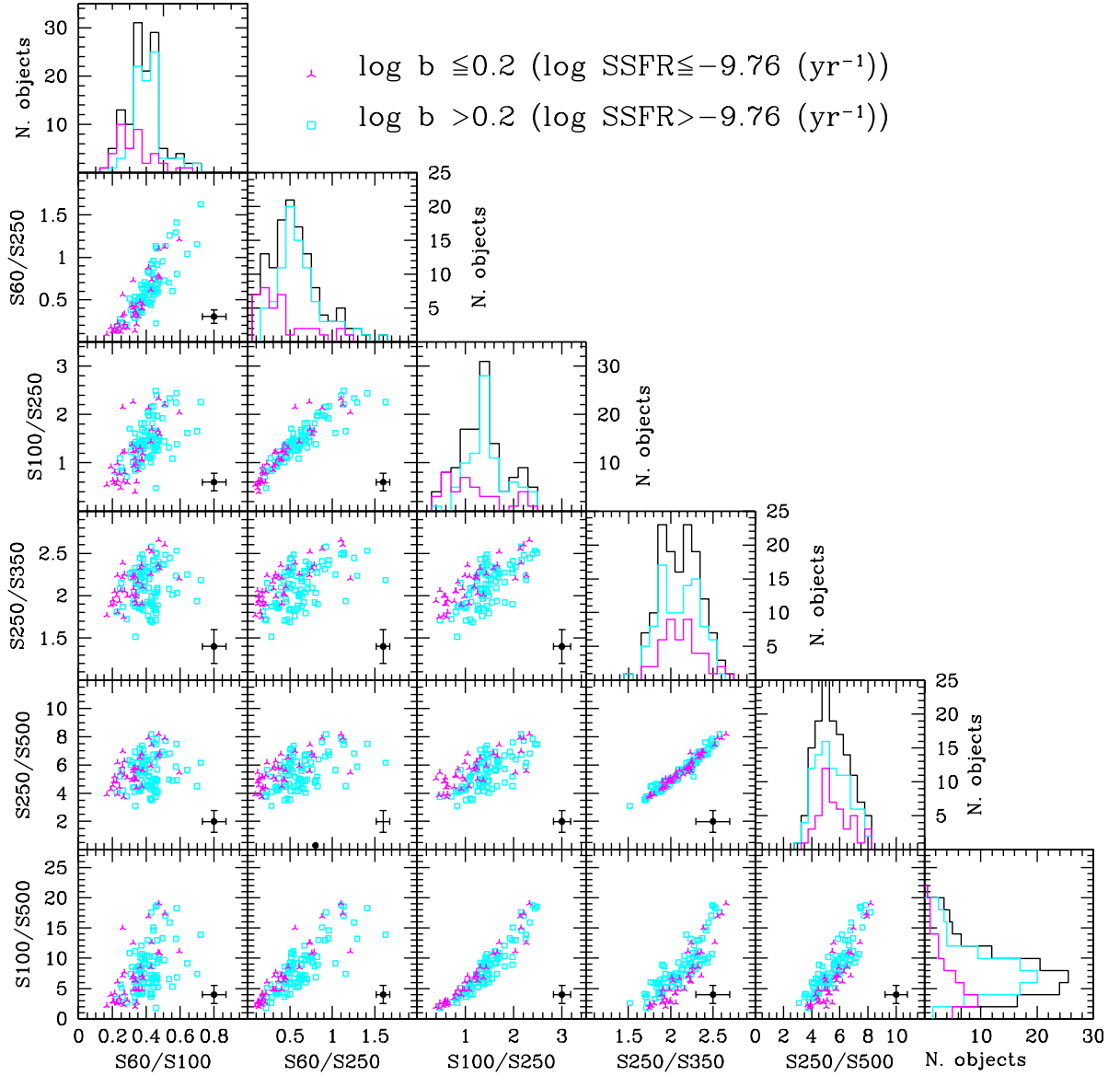


Fig. B.3. Far infrared colour–colour diagrams, equivalent to those shown in Fig. 2, with galaxies coded according to their birthrate parameter, with cyan open squares for objects with $\log b > 0.2$ and magenta three points stars for galaxies with $\log b \leq 0.2$. The black histogram gives the distribution of all galaxies along the X-axis, while the coloured histograms those of the two subsamples of galaxies selected according to their birthrate parameter.

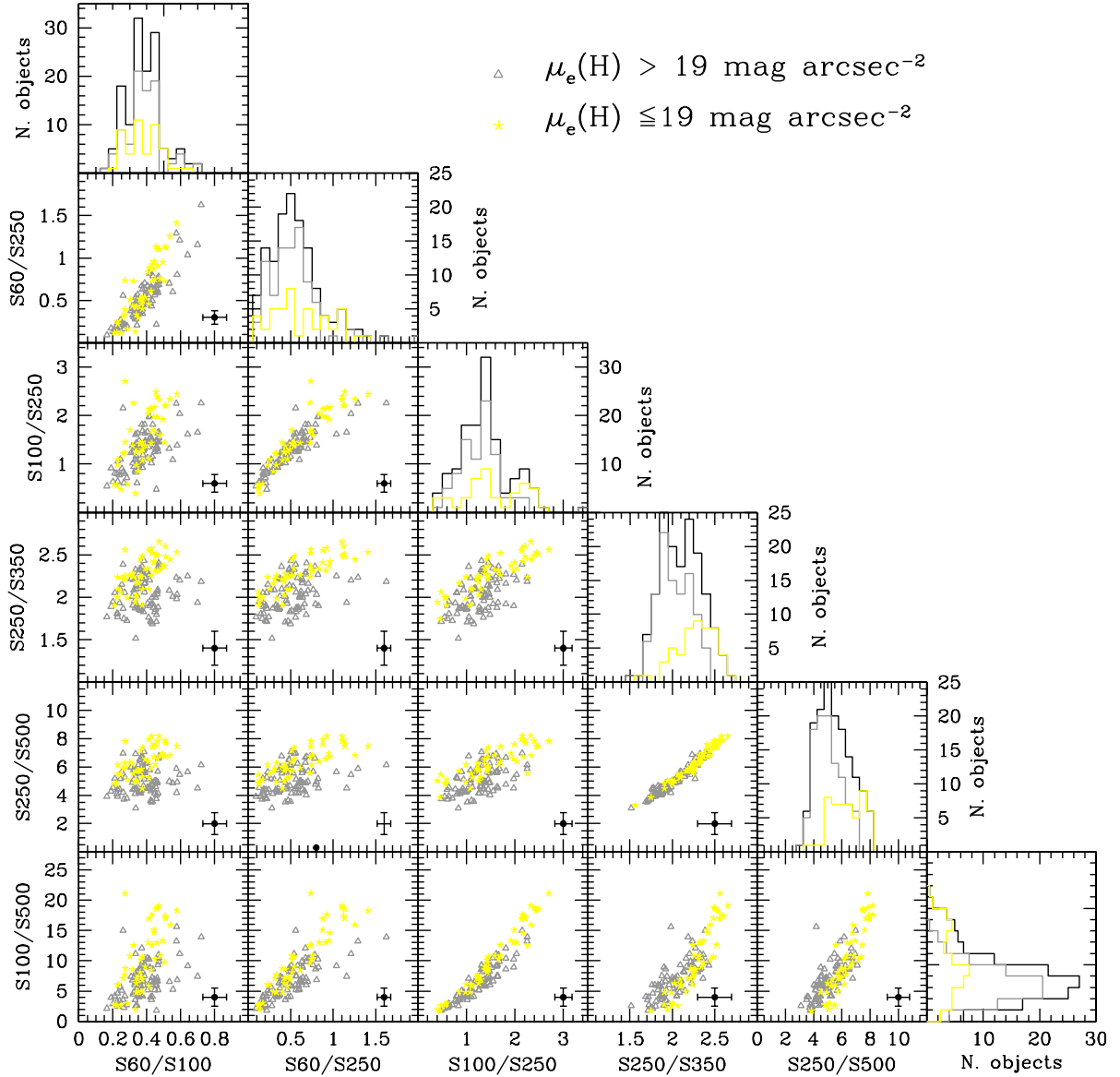


Fig. B.4. Far infrared colour-colour diagrams, equivalent to those shown in Fig. 2, with galaxies coded according to their effective surface brightness, with grey open triangles for objects with $\mu_e(H) > 19 \text{ mag arcsec}^{-2}$ and yellow asterisks for galaxies with $\mu_e(H) \leq 19 \text{ mag arcsec}^{-2}$. The black histogram gives the distribution of all galaxies along the X-axis, while the coloured histograms those of the two subsamples of galaxies selected according to their *H* band effective surface brightness.

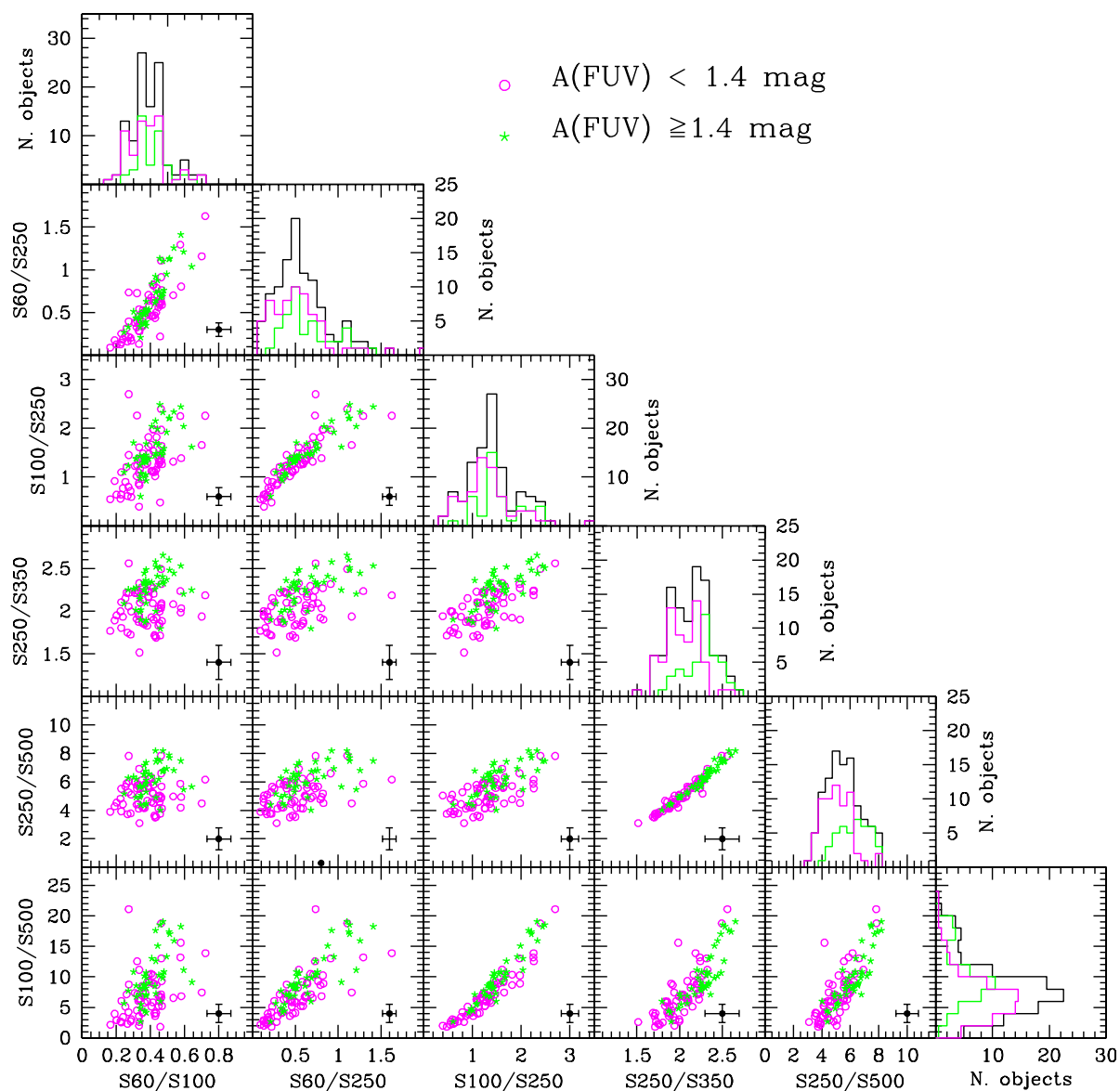


Fig. B.5. Far infrared colour–colour diagrams, equivalent to those shown in Fig. 2, with galaxies coded according to their FUV attenuation, magenta open circles for galaxies with a low attenuation ($A(\text{FUV}) < 1.4$) and green asterisks for objects with $A(\text{FUV}) \geq 1.4$. The black histogram gives the distribution of all galaxies along the X-axis, while the coloured histograms show those of the two subsamples of galaxies selected according to their FUV attenuation.